



1           **Impacts of Cloud Microphysics Parameterizations on Simulated Aerosol-Cloud-**  
2                           **Interactions for Deep Convective Clouds over Houston**

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13



14 **Abstract**

15 Aerosol-cloud interactions remain largely uncertain in predicting their impacts on weather and  
16 climate. Cloud microphysics parameterization is one of the factors leading to the large uncertainty.  
17 Here we investigate the impacts of anthropogenic aerosols on the convective intensity and  
18 precipitation of a thunderstorm occurring on 19 June 2013 over Houston with the Chemistry  
19 version of Weather Research and Forecast model (WRF-Chem) using the Morrison two-moment  
20 bulk scheme and spectral-bin microphysics (SBM) scheme. We find that the SBM predicts a deep  
21 convective cloud agreeing better with observations in terms of reflectivity and precipitation  
22 compared with the Morrison bulk scheme that has been used in many weather and climate models.  
23 With the SBM scheme, we see a significant invigoration effect on convective intensity and  
24 precipitation by anthropogenic aerosols mainly through enhanced condensation latent heating (i.e.,  
25 the warm-phase invigoration). Whereas such effect is absent with the Morrison two-moment bulk  
26 microphysics, mainly due to limitations of the saturation adjustment approach for droplet  
27 condensation and evaporation calculation.

28



## 29 **1 Introduction**

30 Deep convective clouds (DCCs) produce copious precipitation and play important roles in  
31 the hydrological and energy cycle as well as regional and global circulation (e.g., Arakawa, 2004;  
32 Houze, 2014). DCCs and associated precipitation are determined by water vapor, vertical motion  
33 of air, and cloud microphysics that could be affected by aerosols through aerosol-radiative  
34 interactions (ARI) or aerosol-cloud interactions (ACI) or both. The cloud-mediated aerosol effects  
35 are recognized by the Intergovernmental Panel on Climate Change (IPCC) as one of the key  
36 sources of uncertainty in our knowledge of Earth's energy budget and anthropogenic climate  
37 forcing (e.g., Arakawa, 2004; Andreae et al., 2005; Haywood and Boucher, 2000; Lohmann and  
38 Feichter, 2005).

39 Precipitation, latent heat, and cloud radiative forcing associated with DCCs are strongly  
40 associated with cloud microphysical processes, which can be modulated by aerosols through  
41 serving as cloud condensation nuclei (CCN) and ice nuclei (IN). For aerosol-DCC interactions, a  
42 well-known theory is that increasing aerosol concentrations can suppress warm rain as a result of  
43 increased droplet number but reduced droplet size. This allows more cloud water to be lifted to a  
44 higher altitude wherein the freezing of this larger amount of cloud water induces larger latent  
45 heating associated with stronger ice microphysical processes, thereby invigorating convective  
46 updrafts (referred to as “cold-phase invigoration,”; Khain et al. 2005; Rosenfeld et al., 2008). It is  
47 significant in the situations of warm-cloud bases ( $> 15^{\circ}\text{C}$ ; Fan et al., 2012b; Li et al., 2011;  
48 Rosenfeld et al., 2014) and weak wind shear (Fan et al., 2009, 2012b, 2013; Li et al., 2008; Lebo  
49 et al., 2012). Another theory is that increasing aerosols enhances droplet nucleation particularly  
50 secondary nucleation after warm rain initiates, which promotes condensation because of larger  
51 integrated droplet surface area associated with a higher number of small droplets (Fan et al., 2007,



52 2013, 2018; Koren et al., 2014; Lebo, 2018; Sheffield et al., 2015). This so-called “warm-phase  
53 invigoration” , which is manifested in a warm, humid, and clean environment under which the  
54 addition of a large number of ultrafine aerosol particles from urban pollution leads to stronger  
55 invigoration than the “cold-phase invigoration” (Fan et al., 2018). Many factors can affect whether  
56 aerosols invigorate or suppress convective intensity through ACI, such as environmental wind  
57 shear (Fan et al., 2009; Lebo et al., 2012), relative humidity (Fan et al., 2007; Khain et al., 2008),  
58 and Convective Available Potential Energy (Lebo et al., 2012; Morrison, 2012; Storer et al., 2010).  
59 For DCCs with complicated dynamics, thermodynamics, and microphysics, aerosol impacts are  
60 extremely complex and still remain poorly known.

61 Modeling of ACI is quite dependent on cloud microphysics parameterization schemes (e.g.,  
62 Fan et al., 2012a; Khain and Lynn, 2009; Khain et al., 2009, 2015; Lebo and Seinfeld, 2011; Lee  
63 et al., 2018; Loftus and Cotton, 2014; Wang et al., 2013). Two-moment bulk and bin schemes have  
64 been widely used in ACI studies (e.g., Chen et al., 2011; Fan et al., 2013; Khain et al., 2010). In  
65 two-moment bulk schemes, hydrometeor size distributions are diagnosed from the predicted  
66 number and mass with an assumed spectral shape (e.g., gamma function). Saturation adjustment  
67 approach is often used for calculating condensation and evaporation, meaning supersaturation and  
68 undersaturation with respect to water are removed in cloud within a timestep. In bin schemes, the  
69 size distributions of hydrometeors are discretized by a number of size bins and predicted, which  
70 represents some aerosol-cloud interaction processes more physically compared with bulk schemes  
71 (Fan et al., 2016; Khain et al., 2015).

72 Many studies have shown that bulk schemes are limited in representing certain important  
73 microphysical processes such as aerosol activation, condensation, deposition, sedimentation, and  
74 rain evaporation (Ekman et al., 2011; Khain et al., 2009; Lee et al. 2018; Li et al., 2009; Milbrandt



75 and Yau, 2005; Morrison, 2012; Wang et al., 2013). Though bin cloud microphysics can provide  
76 a more rigorous numerical solution and a more robust cloud microphysics representation than  
77 typical bulk microphysics, it is often applied in simulations for process understanding but rarely  
78 in operational applications due to the expensive computation cost. For not introducing further  
79 computation cost, bins schemes are also often run with a prescribed aerosol spectrum assuming a  
80 fixed composition and a simple aerosol budget treatment without coupling with chemistry/aerosol  
81 calculations. As a result, many aerosol life cycle processes such as aerosol nucleation, growth,  
82 aqueous chemistry, aerosol resuspension, and below-cloud wet removal are missing or crudely  
83 parameterized. Therefore, it is difficult to simulate spatial and temporal variabilities of aerosol  
84 chemical composition and size distribution. In Gao et al. (2016), we have coupled a spectral-bin  
85 microphysics scheme (SBM; Fan et al., 2012a; Khain et al., 2004) with the Chemistry version of  
86 Weather Research and Forecast model (WRF-Chem; Grell et al., 2005; Skamarock et al., 2008),  
87 called WRF-Chem-SBM, to address above-mentioned limitations. In this new model, the SBM  
88 was coupled with the Model for Simulating Aerosol Interactions and Chemistry (MOSAIC; Fast  
89 et al., 2006; Zaveri et al., 2008). The newly coupled system was initially evaluated for warm  
90 marine stratocumulus clouds and showed much improved simulation of cloud droplet number  
91 concentration and liquid water content compared with the default Morrison two-moment bulk  
92 scheme (Gao et al., 2016).

93 The Houston area in summer, where isolated convective clouds with very warm cloud-bases  
94 often occurred in the afternoon (Yuan et al., 2008), offers (a) a combination of polluted aerosols  
95 from the urban and industrial area of Houston with significantly low background aerosol  
96 concentrations surrounding Houston, (b) aerosol sources that are not correlated with meteorology,  
97 and (c) weak synoptic forcing along with strong local triggering in the form of land-sea contrasts



98 and sea breeze fronts. This combination allows the manifestation of potentially large aerosol  
99 effects. In this study, we choose a sea-breezed induced DCC case occurring 19-20 June 2013 near  
100 Houston to (1) evaluate the performances of WRF-Chem-SBM in simulating deep convective  
101 clouds and (2) gain a better understanding of the differences in aerosol effects predicted by SBM  
102 and the Morrison two-moment bulk scheme as well as the major factors/processes responsible for  
103 the differences. Considering that the convective clouds over the Houston area are mainly impacted  
104 by the aerosols produced from anthropogenic activities, we focus on the anthropogenic aerosol  
105 effect in this study. The simulated storm case is the same as the case for the Aerosol-Cloud-  
106 Precipitation-Cloud (ACPC) Model Intercomparison Project (Rosenfeld et al., 2014;  
107 [www.acpcinitiative.org](http://www.acpcinitiative.org)).

## 108 **2 Case Description and Observational Data**

109 A local convective event near Houston, Texas on 19-20 June 2013 is selected for the study  
110 owing to the most favorable conditions for simulating isolated convective cells. As above-  
111 mentioned, the case is also selected for the ACPC Model Intercomparison Project  
112 ([www.acpcinitiative.org](http://www.acpcinitiative.org)). The isolated relatively weak convective clouds started from the late  
113 morning because of a trailing front. With increased solar radiation in early afternoon and  
114 strengthening of a sea breeze circulation that transports warm and humid air from the Gulf of  
115 Mexico to Houston urban area, deep convective cells over Houston and Galveston bay areas  
116 developed (Fig. 1). The strong convective cell observed near the Houston city was initiated around  
117 2145 UTC (local time 16:45), and developed to its peak precipitation at 2217 UTC based on radar  
118 observation (Fig. 1). The maximum reflectivity was more than 55 dBZ. This storm cell lasted for  
119 about 1.5 hours.



120 We used the following observation data for model evaluation. Particulate matter (PM) 2.5  
121 data provided by Texas Commission for Environmental Quality (TCEQ) at  
122 <https://www.tceq.texas.gov/agency/data/pm25.html> are used to evaluate the simulated aerosols  
123 near the surface. The data for evaluating cloud base heights and CCN number concentration at  
124 cloud base are obtained from the Visible Infrared Imaging Radiometer Suite (VIIRS) retrievals  
125 based on the method of Rosenfeld et al., (2016). The 2-m temperature and 10-m winds are from  
126 the North American Land Data Assimilation System (NLDAS) with 0.125-deg resolution at  
127 <https://climatedataguide.ucar.edu/climate-data/nldas-north-american-land-data-assimilation->  
128 system. The observed radar reflectivity is used to evaluate the simulated convective system. The  
129 radar reflectivity is obtained from Next-Generation Weather Radar (NEXRAD) network at  
130 <https://www.ncdc.noaa.gov/data-access/radar-data/nexrad-products>, with a temporal frequency of  
131 every ~ 5 minutes and 1 km horizontal spatial resolution.

### 132 **3. Model description and experiments**

133 We conducted model simulations using the version of WRF-Chem based on Gao et al.  
134 (2016) coupling with the Morrison two-moment scheme (Morrison et al., 2005; Morrison et al.,  
135 2009; Morrison and Milbrandt, 2011) and SBM (Khain et al., 2004; Fan et al., 2012). The version  
136 of SBM employed in this study is a fast version of the Hebrew University Cloud Model (HUCM)  
137 described by Khain et al. (2004) with improvements from Fan et al. (2012a) and (2017). The  
138 considered hydrometer size distributions are droplets/raindrops, cloud ice/snow and graupel. The  
139 graupel version is used because it is more appropriate for simulating the convective storm over the  
140 Houston area than the hail version. SBM is currently coupled with the four-sector version of  
141 MOSAIC (0.039-0.156, 0.156-0.624, 0.624-2.5 and 2.5-10.0  $\mu\text{m}$ ). As detailed in Gao et al. (2016),  
142 the aerosol processes including aerosol activation, resuspension, and in-cloud wet-removal are also



143 improved. Theoretically, both aerosol and cloud processes can be more realistically simulated  
144 particularly under the conditions of complicated aerosol compositions and aerosol spatial  
145 heterogeneity compared with original WRF-Chem. The dynamic core of WRF-Chem-SBM is the  
146 Advanced Research WRF model that is fully compressible and nonhydrostatic with a terrain-  
147 following hydrostatic pressure vertical coordinate (Skamarock et al., 2008). The grid staggering is  
148 the Arakawa C-grid. The model uses the Runge-Kutta 3rd order time integration schemes, and the  
149 3rd and 5th order advection schemes are selected for the vertical and horizontal directions,  
150 respectively. The positive definite option is employed for advection of moist and scalar variables.

151 Two nested domains with horizontal grid spacings of 2 and 0.5 km and horizontal grid points  
152 of  $450 \times 350$  and  $500 \times 400$  for Domain 1 and Domain 2, respectively, are used (Fig. 2a), with 51  
153 vertical levels up to 50 hPa. The chemical and aerosol lateral boundary and initial conditions for  
154 Domain 1 simulations were from a quasi-global WRF-Chem simulations at 1-degree grid spacing,  
155 and meteorological lateral boundary and initial conditions were created from MERRA-2 (Gelaro  
156 et al., 2017). Two simulations were run over Domain 1 with anthropogenic emissions turned on  
157 and off, respectively, to provide two different aerosol scenarios for the initial and boundary  
158 chemical and aerosol conditions for Domain 2 simulations: (1) a polluted aerosol scenario with  
159 anthropogenic aerosols accounted which is for the real situation; (2) an assumptive clean scenario  
160 without anthropogenic aerosols. Domain 2 is run with initial and lateral boundary chemical and  
161 aerosols fields from Domain 1 outputs and initial and lateral boundary meteorological conditions  
162 from MERRA-2. Note that we use the meteorology from MERRA-2 as the initial and lateral  
163 boundary conditions for Domain 2 instead of Domain 1 outputs, because we want to keep the  
164 initial and lateral boundary meteorological conditions the same for all the sensitivity tests with



165 different microphysics and aerosol setups (meteorology is different between the two simulations  
166 over Domain 1).

167 The simulations in Domain 1 were initiated at 0000 UTC on 14 Jun and ended at 1200 UTC  
168 20 June with about 5 days for the chemistry spin-up. The meteorological field was reinitialized  
169 every 36 hours to prevent the model drifting. The dynamic time step was 6 s for Domain 1 and 3  
170 s for Domain 2. The anthropogenic emission was from NEI-2011 emissions. The biogenic  
171 emission came from the Model of Emissions of Gases and Aerosols from Nature (MEGAN)  
172 product (Guenther et al., 2006). The biomass burning emission was from the Fire Inventory from  
173 NCAR (FINN) model (Wiedinmyer et al., 2011). We used the Carbon Bond Mechanism Z  
174 (CBMZ) gas-phase chemistry (Zaveri and Peters, 1999) and MOSAIC aerosol model with four  
175 bins (Zaveri et al., 2008). The physics schemes other than microphysics applied in the simulation  
176 are the Unified Noah land surface scheme (Chen and Dudhia, 2001), Mellor-Yamada-Janjic  
177 planetary boundary layer scheme (Janjic et al., 1994), Multi-layer, Building Environment  
178 Parameterization (BEP) urban physics scheme (Salamanca and Martilli, 2010), the RRTMG  
179 longwave and shortwave radiation schemes (Iacono et al., 2008).

180 The main purpose of the simulations in Domain 1 is to provide initial and boundary chemical  
181 and aerosol conditions for the simulations in Domain 2. To save computational cost, WRF-Chem  
182 coupled with Morrison two-moment bulk microphysics scheme (Morrison et al., 2005) is used for  
183 the simulations in Domain 1. Two simulations run for Domain 1 are referred to as D1\_MOR\_anth  
184 in which the anthropogenic emissions are turned on and D1\_MOR\_noanth where the  
185 anthropogenic emissions are turned off. Then four major experiments are carried out to simulate  
186 the convective event near the Houston over Domain 2 with two cloud microphysics schemes and  
187 two aerosol scenarios, respectively. We refer to the simulation in which SBM is used and the



188 anthropogenic emissions are included using the initial and boundary chemicals and aerosols from  
189 D1\_MOR\_anth, as our baseline simulation (referred to as “SBM\_anth”). SBM\_noanth is based on  
190 SBM\_anth but uses initial and boundary chemicals and aerosols from D1\_MOR\_noanth and turns  
191 off the anthropogenic emissions, meaning that anthropogenic aerosols are not taken into account.  
192 MOR\_anth and MOR\_noanth are the two corresponding simulations to SBM\_anth and  
193 SBM\_noanth, respectively, using the Morrison two-moment bulk microphysics scheme. To  
194 examine the contribution of the saturation adjustment approach for condensation and evaporation  
195 to the simulated aerosol effects with the Morrison scheme, we further conducted two sensitivity  
196 tests, based on MOR\_anth and MOR\_noanth, by replacing the saturation adjustment approach in  
197 the Morrison scheme with the condensation and evaporation calculation based on an explicit  
198 representation of supersaturation over a time step as described in Lebo et al. (2012). Note in  
199 both SBM and this modified Morrison schemes, the supersaturation for condensation and  
200 evaporation are calculated after the advection. These two simulations are referred to as  
201 MOR\_SS\_anth and MOR\_SS\_noanth, respectively. To present more robust results, we carry out  
202 a small number of ensembles (three) for each case over Domain 2 (we do not have computer time  
203 to do more ensemble runs). The three ensemble runs are only different in the initialization time:  
204 0000 UTC, 0600 UTC, and 1200 UTC on 19 June. All the simulations end at 1200 UTC 20 June.  
205 All analysis results for Domain 2 simulations in this study are the mean values of three ensemble  
206 runs.

207 We evaluate the aerosol and CCN properties simulated by D1\_MOR\_anth to ensure realistic  
208 aerosol fields, which are used for the Domain 2 simulations with anthropogenic aerosols  
209 considered. These evaluations are included in the section 4.1.



210 From D1\_MOR\_anth simulation, we see a very large spatial variability of aerosol number  
211 concentrations (Fig. 2b). There are three regions with significantly different aerosol loadings over  
212 the domain as shown by the black boxes in Fig. 2b: (a) the Houston urban area, (b) the rural area  
213 about 100 km northeast to Houston, and (c) Gulf of Mexico. Aerosols over the Houston urban area  
214 are mainly contributed by organic aerosols, which are highly related with industrial and ship  
215 channel emissions. The rural area aerosols are mainly from sulfate and sea salt aerosol is the major  
216 contributor over the Gulf of Mexico. This suggests that aerosol properties are extremely  
217 heterogenous in this region. The aerosols over Houston urban area are generally about 5 and 10  
218 times higher than the rural and Gulf area, respectively (Fig. 2c). The size distributions show a  
219 three-mode distribution with the largest differences from the Aitken mode (peaks at 50 nm; Fig.  
220 2c). These ultrafine aerosol particles are mainly contributed by anthropogenic activities (Fig. 2b,  
221 d). With the anthropogenic emissions turned off, the simulated aerosols are much lower and have  
222 much less spatial variability (Fig. 2d).

## 223 **4 Result**

### 224 **4.1 Model Evaluation**

225 We first show the evaluation of the aerosol and CCN properties simulated by  
226 D1\_MOR\_anth, which runs over Domain 1, much larger than Domain 2. As described in Table 1,  
227 there are eight PM monitoring sites from TCEQ around the Houston area. Surface PM<sub>2.5</sub> shows  
228 high concentrations at Houston and its downwind regions (Fig. 3). The values from  
229 D1\_MOR\_anth show a very good agreement with the observations in terms of the surface PM<sub>2.5</sub>  
230 averaged over 24 hours (the day before the convection near Houston). The hourly variations of  
231 ground-level PM<sub>2.5</sub> concentrations from both observation and D1\_MOR\_anth for these sites in



232 the day before the convective initiation is depicted in Fig. 4. Generally, the simulated hourly  
233 pattern agrees with the observation for eight stations. D1\_MOR\_anth reproduces the diurnal  
234 variations, especially the increasing trend from 1200 UTC to 1800 UTC 19 Jun prior to the  
235 initiation of deep convective cells over Houston and Galveston bay areas.

236 The evaluation of the cloud base heights and CCN at cloud bases at the warm cloud stage  
237 before transitioning to deep clouds (2000 UTC) are shown in Fig.5. Over the Houston and its  
238 surrounding area (black box in Fig. 5), the simulated cloud base heights are about 1.5-2 km, in an  
239 agreement with the retrieved values from VIIRS satellite, which are around 1.2-1.8 km (Fig. 5a-  
240 b). The retrieved CCN concentrations at cloud bases vary significantly over the domain and this  
241 spatial variability is generally captured by the model (Fig. 5c-d). For example, D1\_MOR\_anth  
242 simulates some high CCN concentrations ( $400\text{-}800\text{ cm}^{-3}$  with some above  $1000\text{ cm}^{-3}$ ) over the  
243 Houston and around the Bay area, relatively low CCN values at the rural areas (about  $200\text{-}600\text{ cm}^{-3}$ )  
244 and very low values over the Gulf of Mexico (less than  $200\text{ cm}^{-3}$ ), as shown in Fig. 5d. This is  
245 consistent with the spatial variability from the retrievals (Fig. 5c). The evaluation of aerosol  
246 properties before the initiation of Houston convective cells and CCN at the warm cloud stage  
247 before transitioning to deep clouds provides us confidence in using the chemical and aerosol fields  
248 from Domain 1 outputs to feed Domain 2 simulations.

249 Now we are evaluating near-surface temperature and winds, reflectivity and precipitation  
250 simulated by SBM\_anth and MOR\_anth. Fig. 6 shows the comparisons in 2-m temperature and  
251 10-m winds at 1800 UTC (before the convective initiation). Both SBM\_anth and MOR\_anth  
252 capture the general temperature pattern with a little overestimation at the northeast part of the  
253 domain (mainly rural area). SBM\_anth predicts a slightly higher temperature than MOR\_anth in  
254 the northern part of the Houston region (purple box in Fig. 6), which agrees with NLDAS better.



255 SBM\_anth gets the similar southerly winds from Gulf of Mexico to Houston as shown in NLDAS,  
256 while the southerly winds from Gulf of Mexico become very weak or disappear prior to reaching  
257 Houston in MOR\_anth.

258 For the Houston convective cell that we focused (red box in Fig. 7a), SBM\_anth simulates  
259 it well in both location and high reflectivity value (greater than 50 dBZ) in comparison with the  
260 NEXRAD observation (Fig. 7a-b). The simulated composite reflectivities (i.e., the column  
261 maximum) are up to 55-60 dBZ, consistent with NEXRAD. With the Morrison scheme,  
262 MOR\_anth simulates several small convective cells near Houston with maximum reflectivity of  
263 55 dBZ or less (Fig. 7c). The contoured frequency by altitude diagram (CFAD) plots for the entire  
264 storm period show that SBM\_anth is in a better agreement with observation compared with  
265 MOR\_anth, especially for vertical structure of the high reflectivity range (greater than 48 dBZ,  
266 black dashed lines in Fig. 8) and echo top heights, which can reach up to 14-15 km (Fig. 8a-b).  
267 MOR\_anth overestimates the occurrence frequencies of the 35-45 dBZ range and underestimates  
268 those of the low and high reflectivity ranges (less than 15 dBZ or larger than 50 dBZ) as well as  
269 the echo top heights (1-2 km lower than SBM\_anth; Fig. 8c).

270 For the precipitation rates averaged over the study area (red box in Fig. 7), the observation  
271 shows two peaks, which are captured by both SBM\_anth and MOR\_anth (Fig. 9a). However, the  
272 timing for the first peak is about 30 and 60 min earlier in SBM\_anth and MOR\_anth than the  
273 observation, respectively. Also, SBM\_anth predicts the rain rate intensities at the two peak times  
274 more consistent with the observations whereas MOR\_anth underestimates the rain rate intensity at  
275 the second peak time (Fig. 9a). The large precipitation rates (greater than 15 mm h<sup>-1</sup>) in SBM\_anth  
276 has a ~1.5 times larger occurrence probability than those in MOR\_anth, showing a better  
277 agreement with the observation (Fig. 9b). Overall, the performance of SBM\_anth is superior to



278 MOR\_anth in simulating the location and intensity of the convective storm and associated  
279 precipitation.

#### 280 4.2 Simulated Aerosol Effects on Cloud and Precipitation

281 Now we look at the effects of anthropogenic aerosols on the deep convective storm  
282 simulated with SBM and Morrison microphysics schemes. Fig. 9a shows that with the SBM  
283 scheme, anthropogenic aerosols remarkably increase the mean surface rain rates (by ~30%; from  
284 SBM\_noanth to SBM\_anth), mainly because of the increased occurrence frequency (nearly  
285 doubled) for relatively large rain rates (i.e., 10-15 mm h<sup>-1</sup> and >15 mm h<sup>-1</sup>) in Fig. 9b. With the  
286 Morrison scheme, the changes in mean precipitation and the PDF from MOR\_noanth to  
287 MOR\_anth are relatively small, showing a very limited aerosol effect on precipitation. With the  
288 SBM scheme, the increase in the updraft speeds by the anthropogenic aerosols is even more notable  
289 than the precipitation (Fig. 10a-b). Above 5-km altitude, the occurrence frequencies of updraft  
290 speeds greater than 0.4% extend to a much larger values, with 36 m s<sup>-1</sup> at the upper levels in  
291 SBM\_anth while only ~ 20 m s<sup>-1</sup> in SBM\_noanth. With the Morrison scheme, the changes are not  
292 significant by the anthropogenic aerosols (MOR\_noanth vs MOR\_anth in Fig. 10c-d). From  
293 MOR\_noanth to MOR\_anth, there is a slight increase in updraft speed at around 9-11 km altitudes  
294 but a slight decrease at 6-8 km altitudes. The significant invigoration of convective intensity by  
295 anthropogenic aerosols with the SBM scheme explains the much larger occurrences of relatively  
296 large rain rates and overall more surface precipitation due to the anthropogenic aerosol effect (Fig.  
297 9).

298 Now the question is why the anthropogenic aerosols enhance convective intensity of the  
299 storm with the SBM scheme while the effect is very small with the Morrison scheme. Fig. 11  
300 shows the vertical profiles of mean updraft velocity, thermal buoyancy, and total latent heating



301 rate of the top 25<sup>th</sup> percentile updrafts with value greater than  $2 \text{ m s}^{-1}$  during the deep convective  
302 cloud stage. With the SBM microphysics scheme, the increased convective intensity due to  
303 anthropogenic aerosol effect corresponds to the increased thermal buoyancy which is particularly  
304 notable at upper levels ( $\sim 20\%$ ) from SBM\_noanth to SBM\_anth (Fig. 11a, c). The increased  
305 thermal buoyancy can be explained by the increased total latent heating (Fig. 11e), which is mainly  
306 from the larger condensation latent heating (Fig. 12a). From SBM\_noanth to SBM\_anth, the latent  
307 heating from ice-related microphysical processes (including deposition, drop freezing, and riming)  
308 has a relatively smaller increase than that from condensation (about half of the increase in  
309 condensation latent heating as shown in Fig. 12a). As shown in Fan et al., (2018), the increase in  
310 lower-level condensation latent heating has a much larger effect on intensifying updraft intensity  
311 compared with the same amount of increase in high-level latent heating from ice-related  
312 microphysical processes. This suggests that the convective invigoration by the anthropogenic  
313 aerosols with the SBM scheme should be mainly through the “warm-phase invigoration”  
314 mechanism. Compared with the Morrison scheme, the increase of total latent heating by the  
315 anthropogenic aerosols is almost doubled with the SBM scheme, explaining more remarkable  
316 enhancement of thermal buoyancy and thus the convective intensity (red lines vs blue lines in Fig.  
317 11). From MOR\_noanth to MOR\_anth, there is a small increase in both the condensation latent  
318 heating and high-level latent heating associated with ice-related processes (blue lines in Fig. 12b).  
319 The major difference in the increase of latent heating by the anthropogenic aerosols between SBM  
320 and Morrison microphysics schemes comes from the condensation latent heating, with a  $\sim 20\%$   
321 increase with SBM but only  $\sim 8\%$  with Morrison (Fig. 12). The lack of significant increase in  
322 condensation latent heating limits the “warm-phase invigoration”, mainly responsible for the



323 limited aerosol impacts on convective intensity and associated precipitation with the Morrison  
324 scheme.

325 To understand why the responses of condensation to the anthropogenic aerosols are  
326 different between the SBM and Morrison schemes, we look into the process rates of drop  
327 nucleation and condensation (Fig. 13). The calculations of aerosol activation and  
328 condensation/evaporation in the SBM scheme are based on the Köhler theory and diffusional  
329 growth equations in light of particle size and supersaturation, respectively. Whereas in the Morrison  
330 scheme, the Abdul-Razzak and Ghan (2002) parameterization is used for aerosol activation and  
331 the saturation adjustment method is applied for condensation and evaporation calculation. With  
332 the SBM scheme, the anthropogenic aerosols increase the drop nucleation rates by a few times  
333 over the profile (red lines in Fig. 13a), and the condensation rates are also drastically increased  
334 (doubled between 4–6 km altitudes as shown in Fig. 13c). The enhanced condensation rate by the  
335 anthropogenic aerosols is because much more aerosols are activated to form a larger number of  
336 small droplets, increasing the integrated droplet surface area for condensation, as documented in  
337 Fan et al., (2018). As a result, supersaturation is drastically lower in SBM\_anth than SBM\_noanth  
338 (green lines in Fig. 13a). With the Morrison scheme, we still see a large increase in droplet  
339 nucleation rate (Fig. 13b). However, the condensation rates are barely increased (blue solid vs.  
340 dashed lines in Fig. 13d). We hypothesize that the lack of response of condensation to the increased  
341 aerosol activation with the Morrison scheme is mainly because of the saturation adjustment  
342 calculation of the condensation and evaporation process. The approach does not allow  
343 supersaturation in cloud and the calculation does not depend on supersaturation, thus removes the  
344 sensitivity to the anthropogenic aerosols.



345 To verify our hypothesis and examine how much the saturation adjustment method is  
346 responsible for the weak responses of condensation latent heating and convection to the added  
347 anthropogenic aerosols, we conducted two additional sensitivity tests by replacing the  
348 saturation adjustment approach in Morrison scheme with the condensation and evaporation  
349 calculation based on an explicit representation of supersaturation over a time step, as described  
350 in Section 3. The result shows the Morrison scheme with the simple calculation of supersaturation  
351 for condensational growth significantly changes the condensation rate (orange vs. blue lines in Fig.  
352 13d) and a similarly large enhancement (from MOR\_SS\_noanth to MOR\_SS\_anth in Fig. 13d) is  
353 seen as the SBM scheme (Fig. 13c). This leads to a larger increase in condensation latent heating  
354 (orange lines in Figure 12b) compared with the original Morrison scheme, resulting a similarly  
355 large increase in thermal buoyancy by the anthropogenic aerosols as with the SBM scheme (orange  
356 lines in Fig. 11d), thus a similarly large increase in the convective intensity (orange lines in Fig.  
357 11b). The increase of precipitation from MOR\_SS\_noanth to MOR\_SS\_anth is also similar to that  
358 with the SBM scheme (not shown). These results verify that the saturation adjustment approach  
359 for parameterizing condensation and evaporation is the major reason responsible for limited  
360 aerosol effects on convective intensity and precipitation with the Morrison scheme. Past studies  
361 also showed the limitations of the saturation adjustment approach in simulating aerosol impacts  
362 on deep convective clouds (e.g., Fan et al., 2016; Lebo et al., 2012; Lee et al., 2018; Wang et al.,  
363 2013).

364 Fig. 14 shows the responses of hydrometeor mass to anthropogenic aerosol effects. With  
365 the SBM scheme, the increases of cloud mass, rain mass, and total ice mass (ice, snow, and graupel)  
366 by the anthropogenic aerosols are very significant (Fig. 14, left), corresponding to convective  
367 invigoration. The increase of the total ice mass is particularly significant (from 3.5 to 5.5 g kg<sup>-1</sup>



368 around 10-km altitude), suggesting a large effect of enhanced convective intensity on ice  
369 hydrometeors. However, with the Morrison scheme, little change is seen (Fig. 14, right, blue lines).  
370 By replacing the saturation adjustment with a simple calculation based on supersaturation for  
371 condensation and evaporation in the Morrison scheme, the increases in those hydrometeor masses  
372 become as evident as those with the SBM scheme (Fig. 14, right, orange lines).

### 373 **5 Conclusions and Discussion**

374 We have conducted model simulations of a deep convective cloud case occurring on 19 June  
375 2013 over the Houston area with WRF-Chem coupled with the SBM and Morrison microphysics  
376 schemes to (1) evaluate the performance of WRF-Chem-SBM in simulating the deep convective  
377 clouds, and (2) explore the differences in aerosol effects on the deep convective clouds produced  
378 by the SBM and Morrison schemes and the major factors responsible for the differences.

379 We have evaluated the simulated aerosols, CCN, cloud base heights, reflectivity, and  
380 precipitation. The model simulates the large spatial variability of aerosols and CCN from Gulf of  
381 Mexico, rural area, to Houston city. On the bulk magnitudes, the model captures the surface PM<sub>2.5</sub>,  
382 cloud base height, and CCN at cloud base near the Houston reasonably well. These realistically  
383 simulated aerosol fields were fed to higher resolution simulations (0.5 km) using the SBM and  
384 Morrison schemes. With the SBM scheme, the model simulates a deep convective cloud over the  
385 Houston in a better agreement with the observed radar reflectivity and precipitation, compared  
386 with using the Morrison scheme.

387 By excluding the anthropogenic aerosols in the simulations, the effects of anthropogenic  
388 aerosols on the deep convective clouds and differences in aerosol effects using the two  
389 microphysics schemes were examined. With the SBM scheme, anthropogenic aerosols notably  
390 increase convective intensity, enhance the peak precipitation rate over the Houston area (by ~30%),



391 and double the frequencies of relatively large rain rates ( $> 10 \text{ mm h}^{-1}$ ). The enhanced convective  
392 intensity by anthropogenic aerosols makes the simulated storm agree better with the observed,  
393 mainly attributed to the increased condensation latent heating, indicating the “warm-phase  
394 invigoration”. In contrast, with the Morrison scheme, there is no significant anthropogenic aerosol  
395 effect on the convective intensity and precipitation.

396 Sensitivity tests by replacing the saturation adjustment with the condensation and evaporation  
397 calculation based on an explicit representation of supersaturation over a time step show the  
398 similar aerosol effects on condensation, convective intensity, hydrometeor mass mixing ratios, and  
399 precipitation as with the SBM scheme. Therefore, the saturation adjustment method for the  
400 condensation and evaporation calculation is mainly responsible for the limited aerosol effects with  
401 the Morrison scheme. This is because the saturation adjustment method does not allow for the  
402 “warm-phase invigoration”, which is different from Lebo et al. (2012) showing that the saturation  
403 adjustment artificially enhanced condensation latent heating at low levels and limited the potential  
404 for aerosols to invigorate convection through the “cold-phase invigoration” mechanism in their  
405 idealized simulations of a supercell storm with the thermal bubble initiation. In this study of the  
406 thunderstorm with WRF real-case simulations for both chemistry/aerosols and clouds, the  
407 saturation adjustment method actually leads to a smaller condensation latent heating than the  
408 explicit calculation with supersaturation (solid bold blue vs. solid bold orange line in Fig. 12b).  
409 Thus, when the computational resource is not sufficient or in other situations such as the  
410 application of SBM is not available, the Morrison scheme modified with the condensation and  
411 evaporation calculation based on a simple representation of supersaturation can be applied to study  
412 aerosol effects on convective clouds, especially for warm and humid cloud cases in which the  
413 response of condensation to aerosols is particularly important.



414           Following Fan et al., (2018), which showed that the “warm-phase invigoration” mechanism  
415 was manifested by ultrafine aerosol particles in the Amazon warm and humid environment with  
416 extremely low background aerosol particles. Here we showed that in summer anthropogenic  
417 aerosols over the Houston area may also enhance the thunderstorm intensity and precipitation  
418 through the same mechanism by secondary nucleation of numerous ultrafine aerosol particles from  
419 the anthropogenic sources. But the magnitude of the effect is not as substantial as in the Amazon  
420 environment. Possible reasons include that background aerosols are much higher over the Houston  
421 area and air is not as humid as Amazon.

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432

433



434 **Reference**

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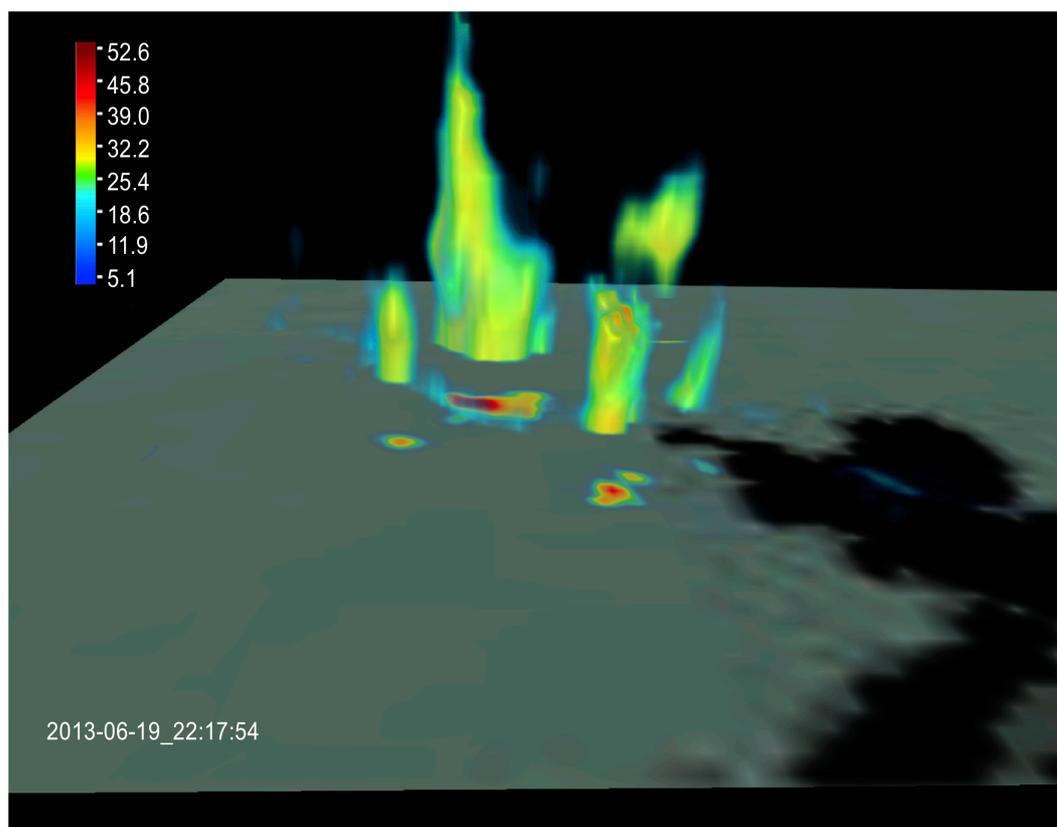
623 **Table 1** Descriptions of the PM<sub>2.5</sub> Monitoring Sites over the Houston area from TCEQ

Abbreviation	Site Descriptions	Latitude	Longitude
HA	Houston Aldine	29.901	-95.326
HDP	Houston Deer Park 2	29.670	-95.129
SFP	Seabrook Friendship Park	29.583	-95.016
CR	Conroe Relocated	30.350	-95.425
KW	Kingwood	30.058	-95.190
CT	Clinton	29.734	-95.258
PP	Park Place	29.686	-95.294
GS	Galveston 99th Street	29.254	-94.861

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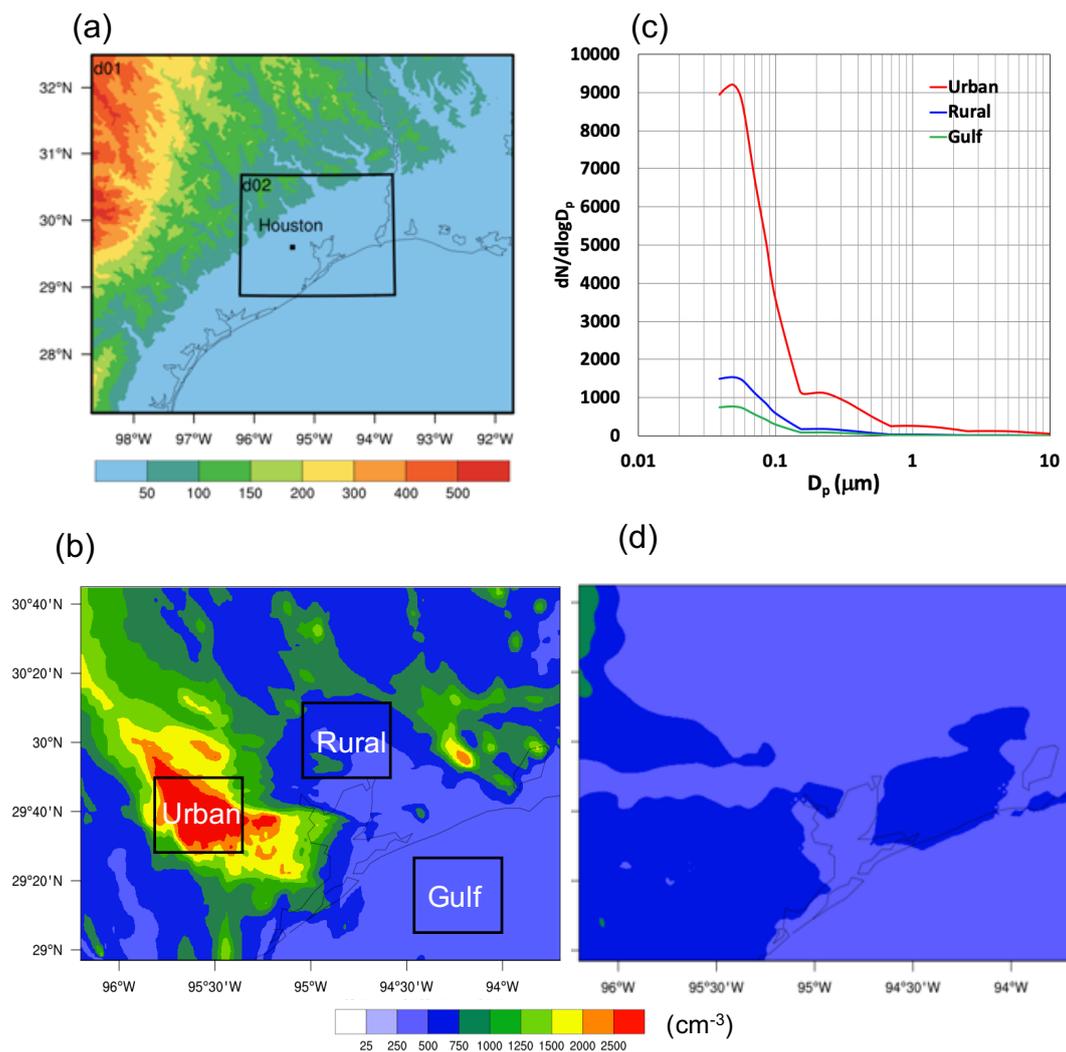
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628 **Figure 1** 3D structure snapshot of radar reflectivity (unit: dBZ) from NEXRAD, overlaid with the  
629 composite reflectivity shown on the surface at the time when the maximum reflectivity is observed  
630 (2217 UTC). The dark shade shows the water body and the largest cell is in the Houston.

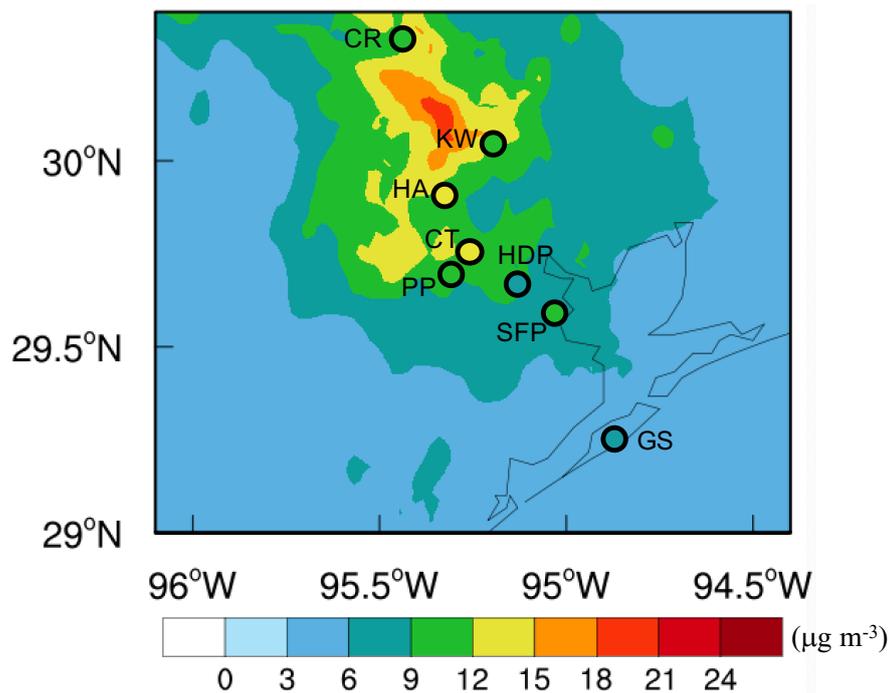
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632

633 **Figure 2** (a) Simulation domains with the terrain heights (unit: m), (b) aerosol number  
634 concentration (unit: cm<sup>-3</sup>) from D1\_MOR\_anth, (c) aerosol size distributions over the urban, rural,  
635 and Gulf of Mexico as marked by three black boxes in Fig. 2b at 1200 UTC, 19 Jun 2013 (6-hr  
636 before the convection initiation), and (d) the same as Fig. 2b, but for D1\_MOR\_noanth in which  
637 the anthropogenic aerosols are excluded.

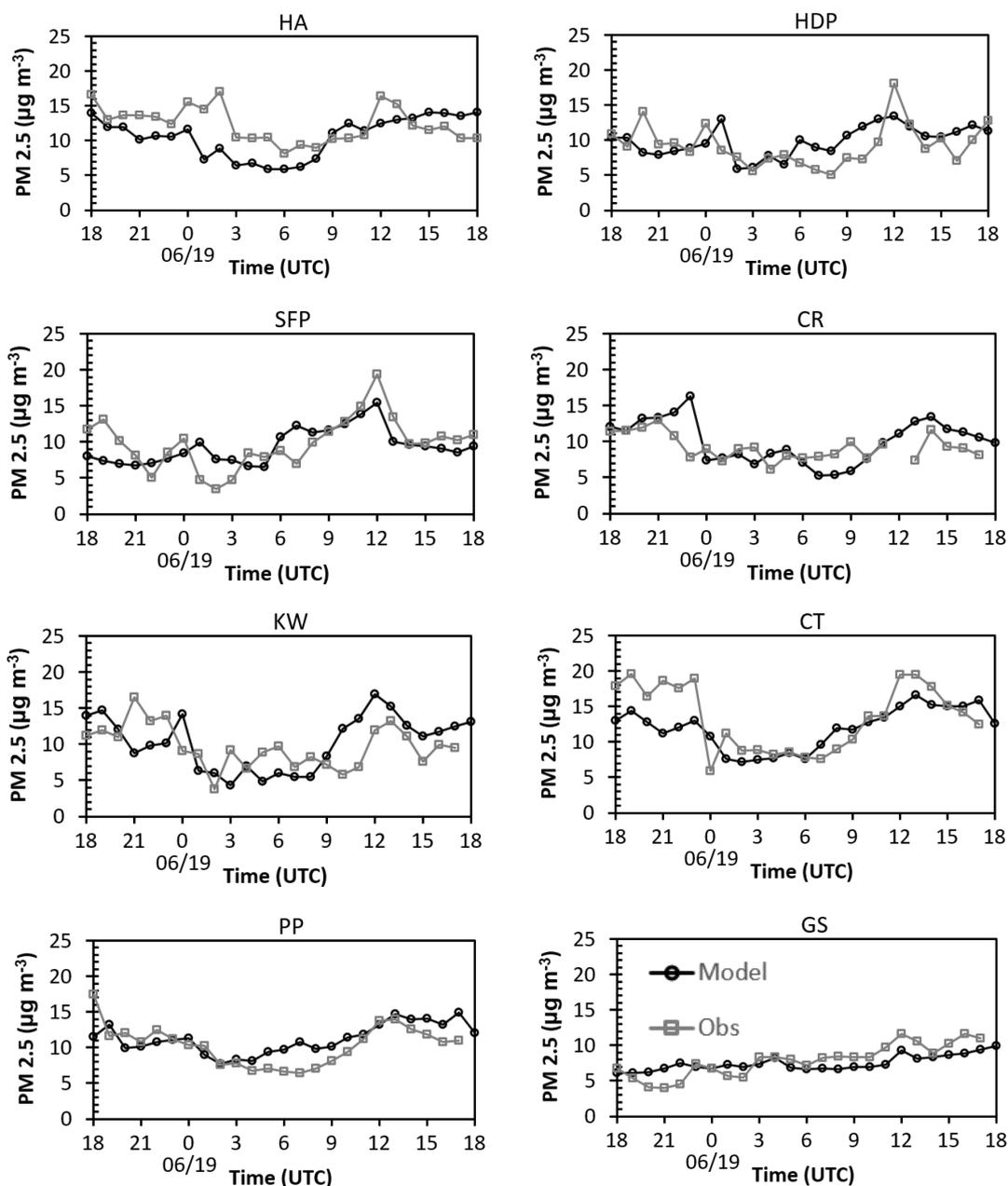
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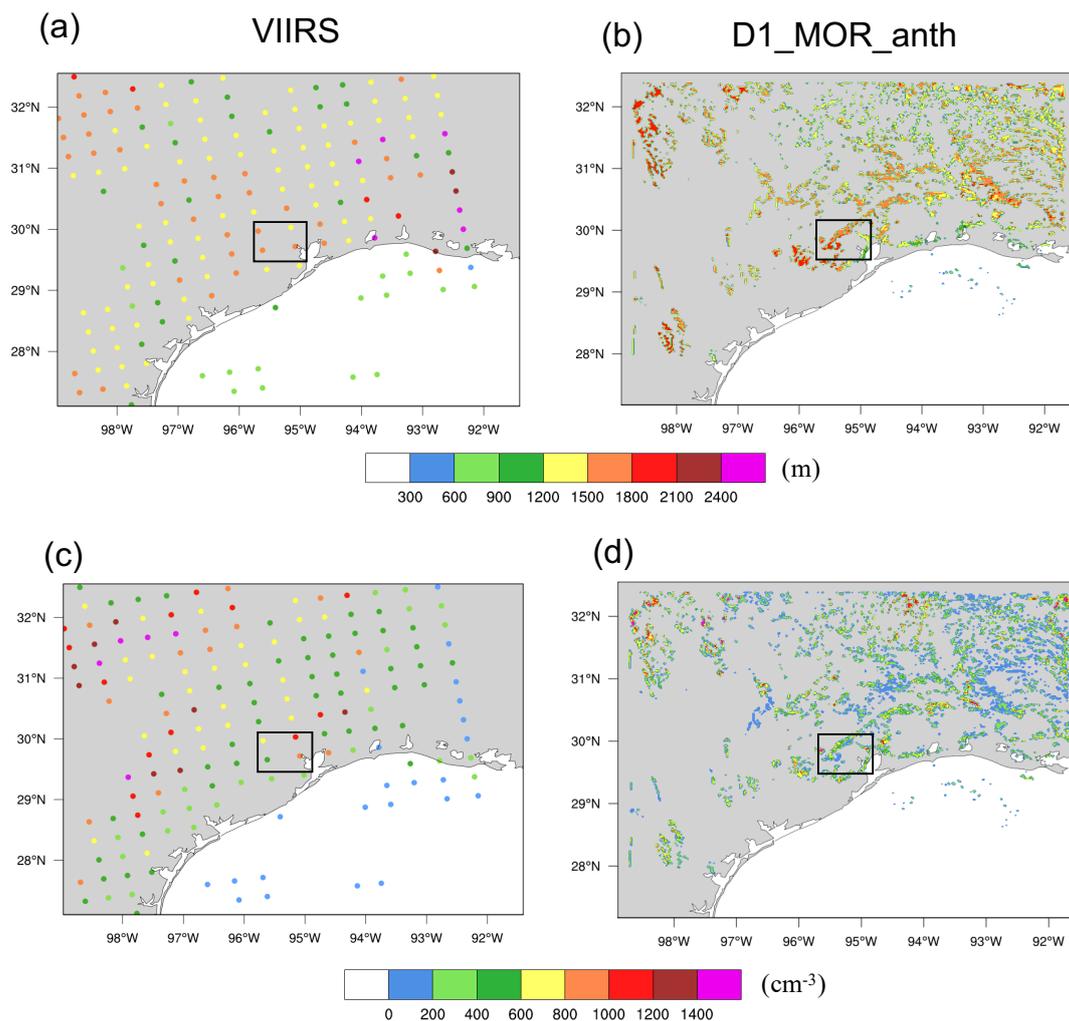
640 **Figure 3** Comparisons of 24-hr averaged PM<sub>2.5</sub> mass concentrations (unit:  $\mu\text{g m}^{-3}$ ) between model  
641 simulation D1\_MOR\_anth (contoured) and site observation from TCEQ (colored circles) from  
642 1800 UTC, 18 June 2013 to 1800 UTC, 19 June 2013 (1 day before the convection initiation). The  
643 site names and other information are shown in Table 1.

644



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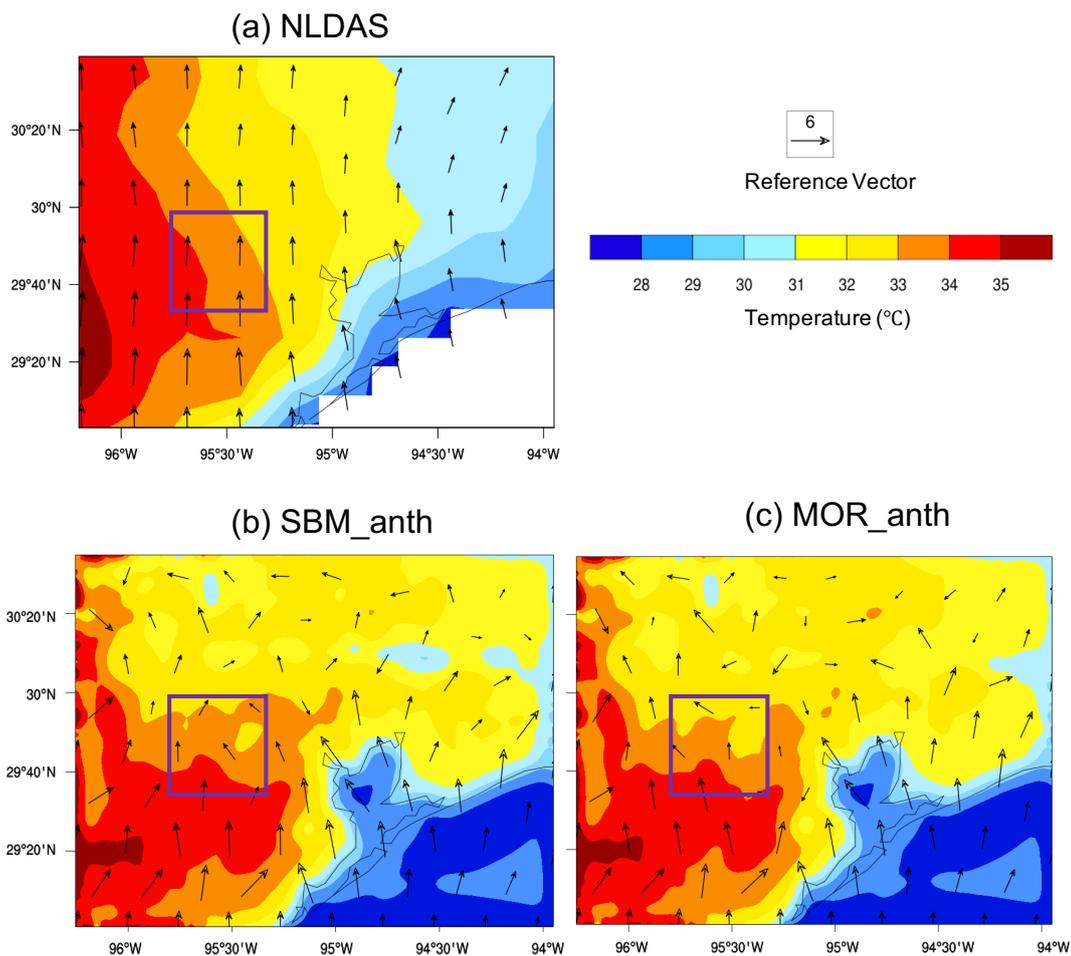
646 **Figure 4** Site-by-site comparisons of hourly PM<sub>2.5</sub> mass concentrations (unit:  $\mu\text{g m}^{-3}$ ) from  
647 D1\_MOR\_anth and TCEQ site observation over 24 hours from 1800 UTC, 18 June 2013 to 1800  
648 UTC, 19 June 2013 (1 day before the convection initiation).



649

650 **Figure 5** Evaluation of (a,b) cloud base heights (unit: m) and (c,d) CCN number concentration at  
651 cloud base (unit: cm<sup>-3</sup>) from VIIRS satellite (left) retrieved at 1943 UTC (Rosenfeld et al. 2016)  
652 and model simulation D1\_MOR\_anth (right) at 2000 UTC, 19 June 2013. The Houston area is  
653 marked as the black box. Satellite-retrieved cloud base height was calculated from the difference  
654 between reanalysis surface air temperature (from reanalysis data) and VIIRS-measured cloud base  
655 temperature (warmest cloudy pixel) divided by the dry adiabatic lapse rate, while modeled cloud  
656 base height was determined by the lowest cloud layer with cloud mass mixing ratio greater than  
657 10<sup>-5</sup> kg kg<sup>-1</sup>.

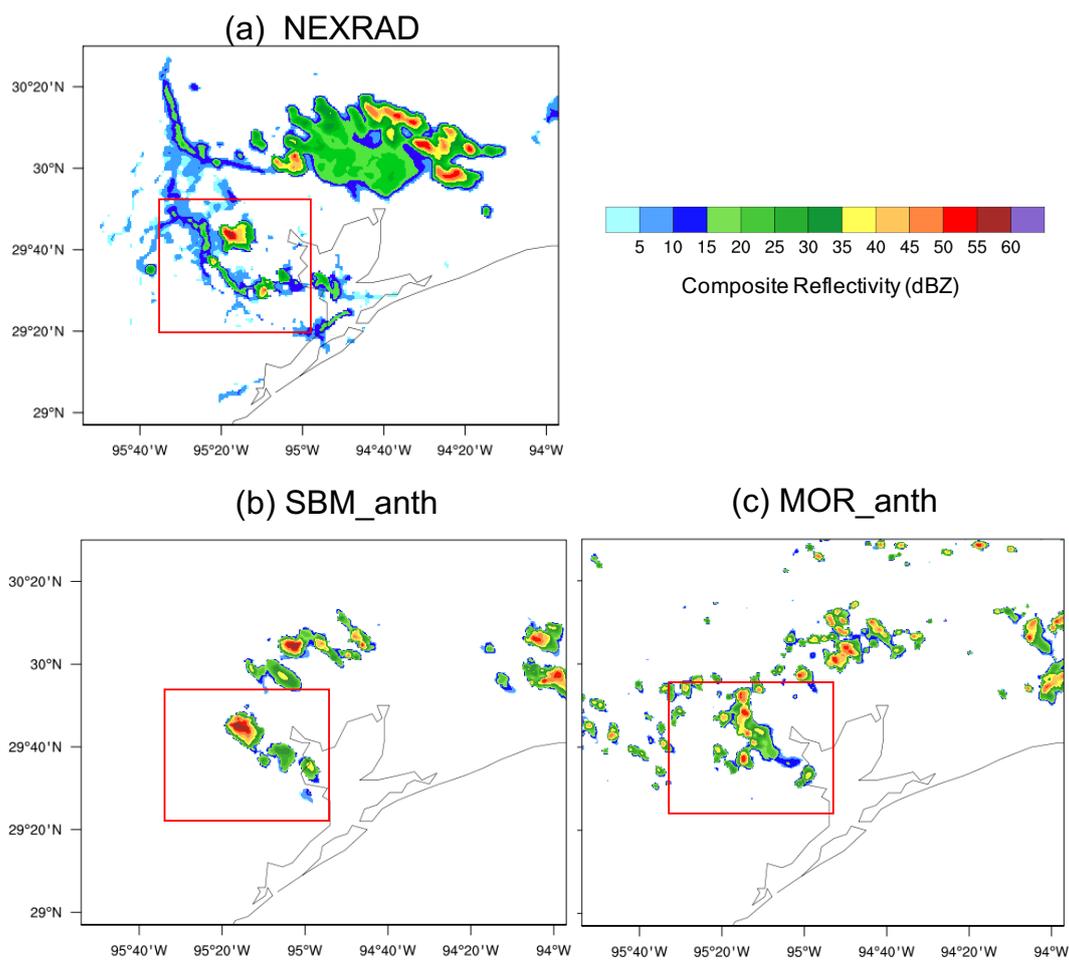
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659

660 **Figure 6** 2-m Temperature (shaded; unit: °C) and 10-m winds (vectors; unit:  $\text{m s}^{-1}$ ) from (a)  
661 NLDAS, (b) SBM\_anth and (c) MOR\_anth at 1800 UTC, 19 Jun 2013. The purple box denotes  
662 the Houston area.

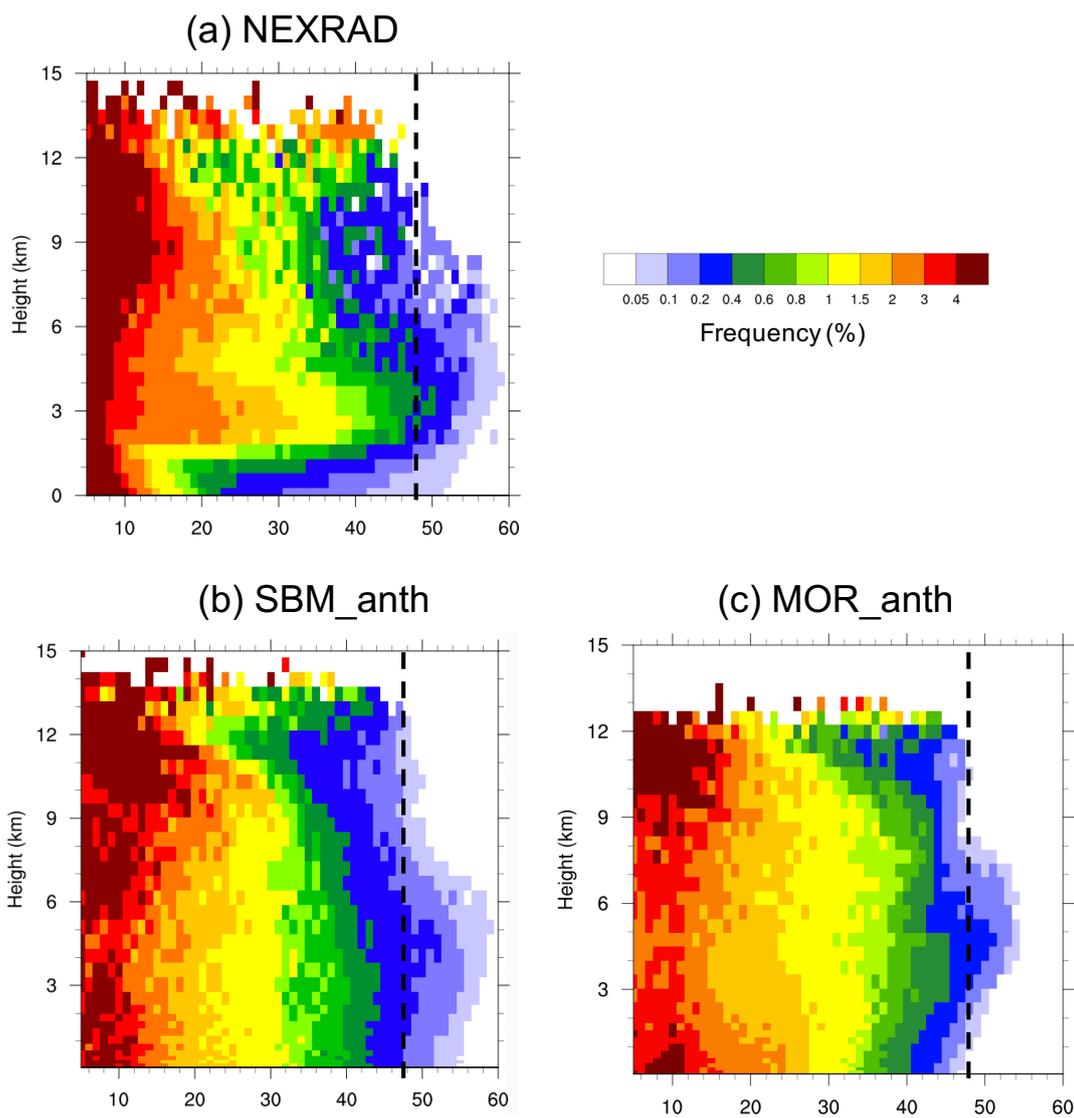
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664

665 **Figure 7** Composite reflectivity (unit: dBZ) from (a) NEXRAD (2217 UTC), (b) SBM\_anth  
666 (2140 UTC) and (c) MOR\_anth (2125 UTC) when maximum reflectivity in Houston is observed  
667 on 19 June 2013. The red box is the study area for convection cells near Houston.

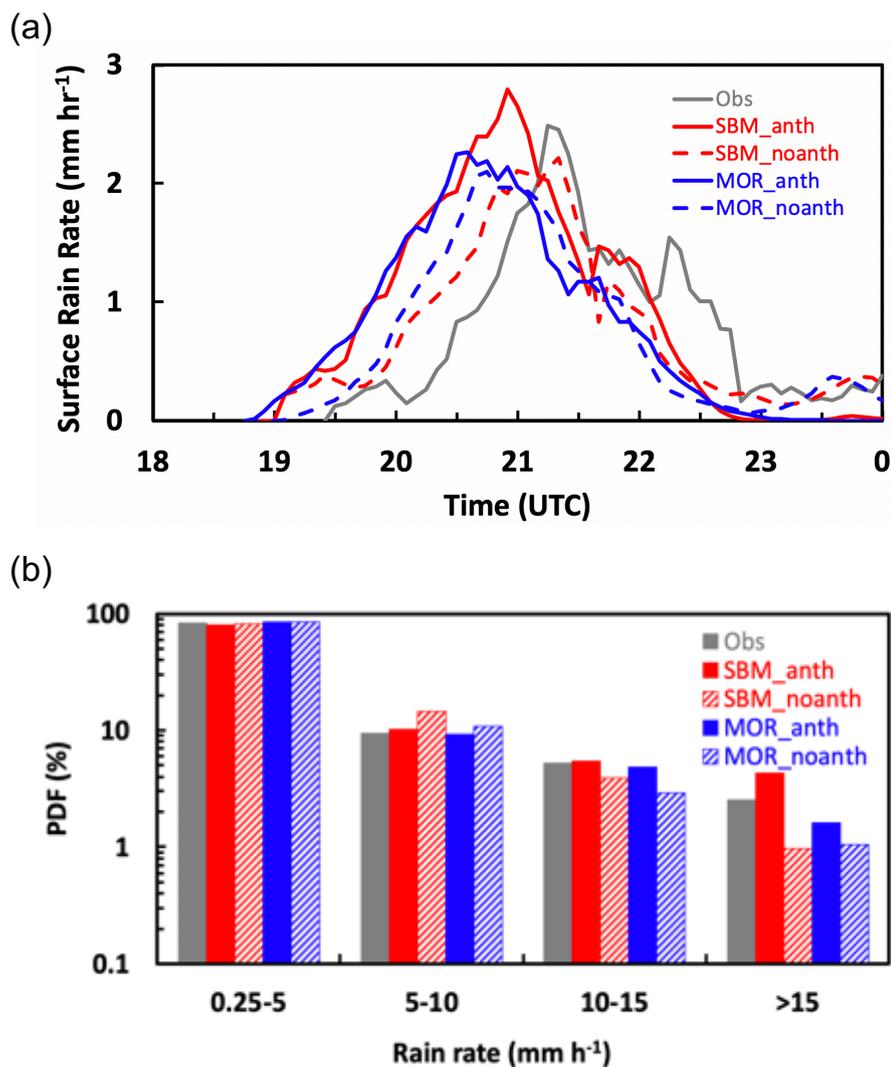
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669

670 **Figure 8** The CFAD of reflectivity (unit: dBZ) for the values larger than 0 dBZ from (a) NEXRAD,  
671 (b) SBM\_anth and (c) MOR\_anth over the study area (red box in Fig. 7) from 1800 UTC, 19 Jun  
672 to 0000 UTC, 20 Jun 2013. The black solid lines denote the reflectivity with the value of 48 dBZ.

673



674

675 **Figure 9** (a) Time series of averaged surface rain rate (unit: mm h<sup>-1</sup>) and (b) PDFs of rain rate for

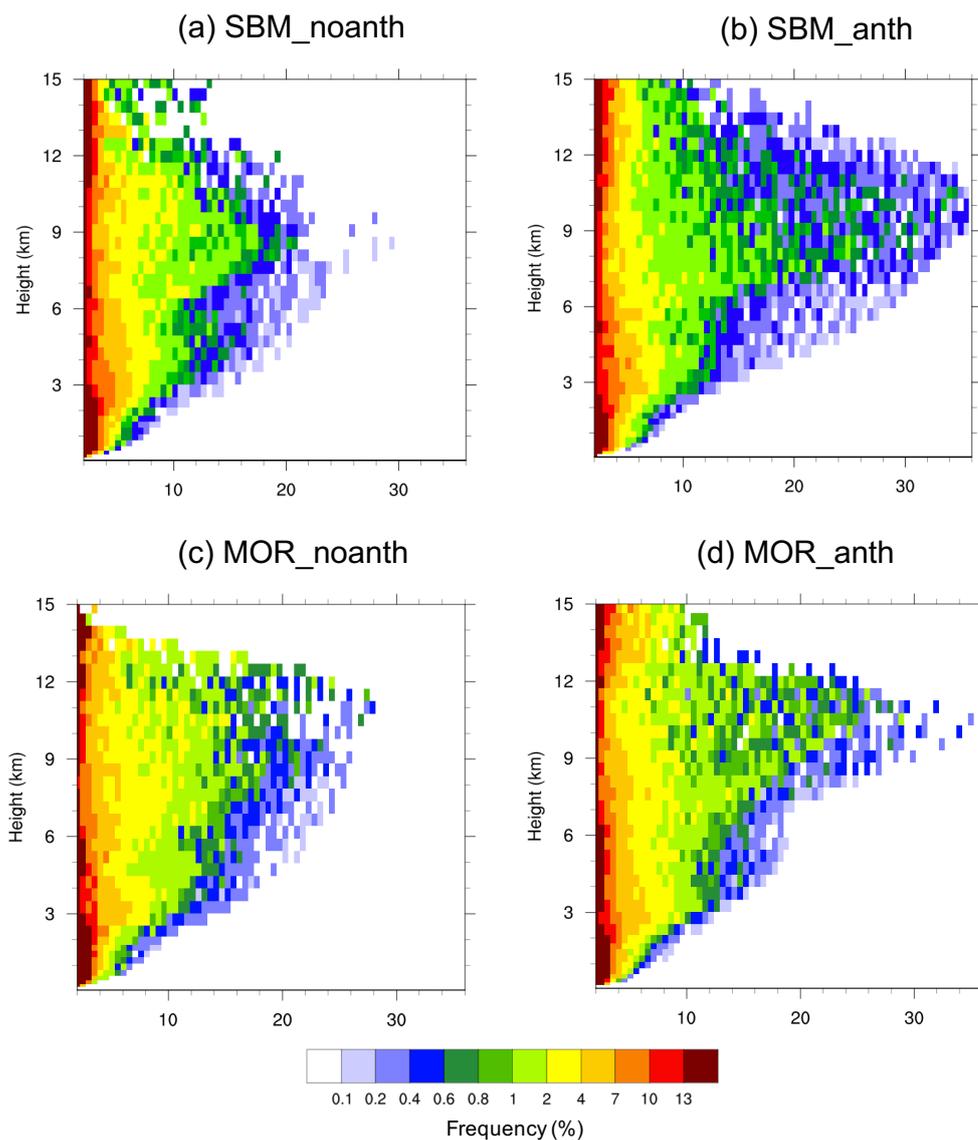
676 the values larger than 0.25 mm h<sup>-1</sup> over the study area (red box in Fig. 7) from observation (grey),

677 SBM\_anth and SBM\_noanth (red), MOR\_anth and MOR\_noanth (blue) from 1800UTC, 19 Jun

678 2013 to 0000 UTC, 20 Jun 2013. The observed precipitation rate is obtained by NEXRAD

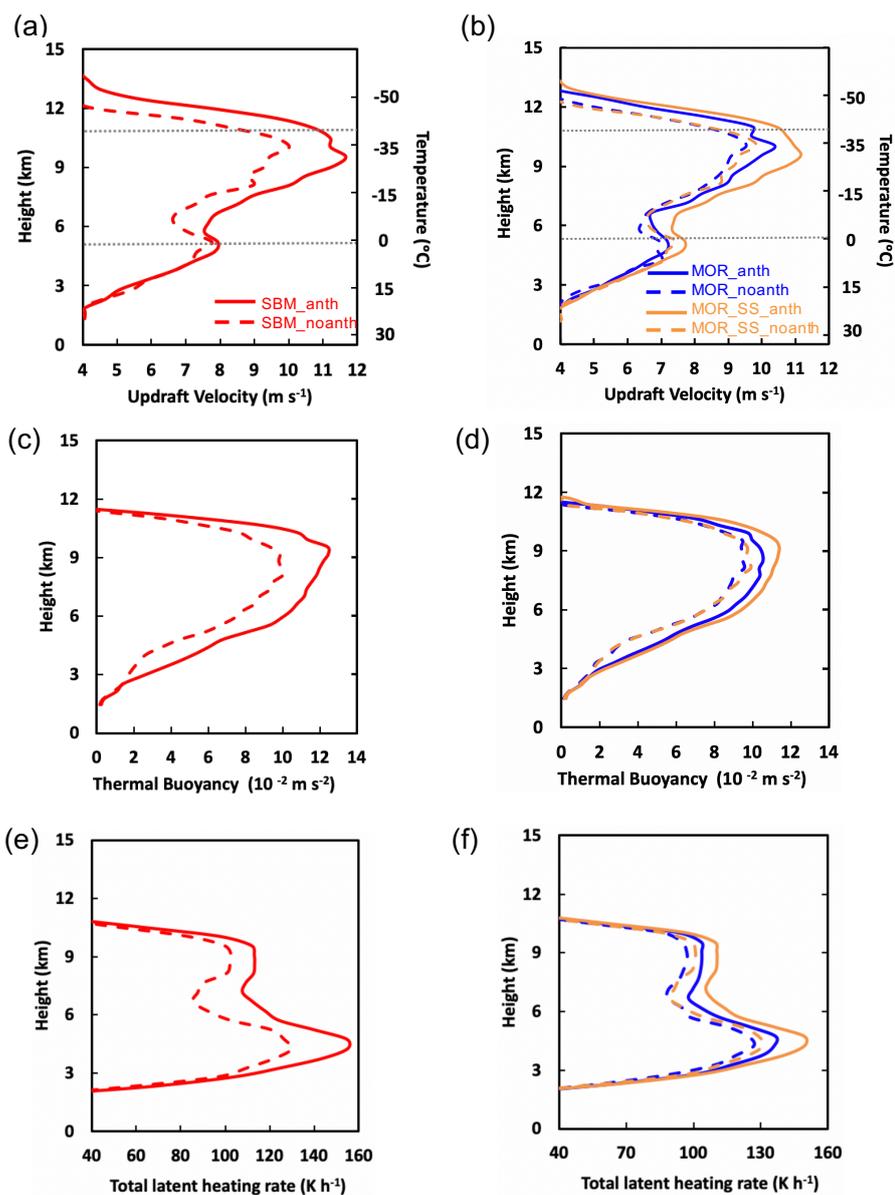
679 retrieved rain rate. Both observation and model data are in every 5-min frequency.

680



681

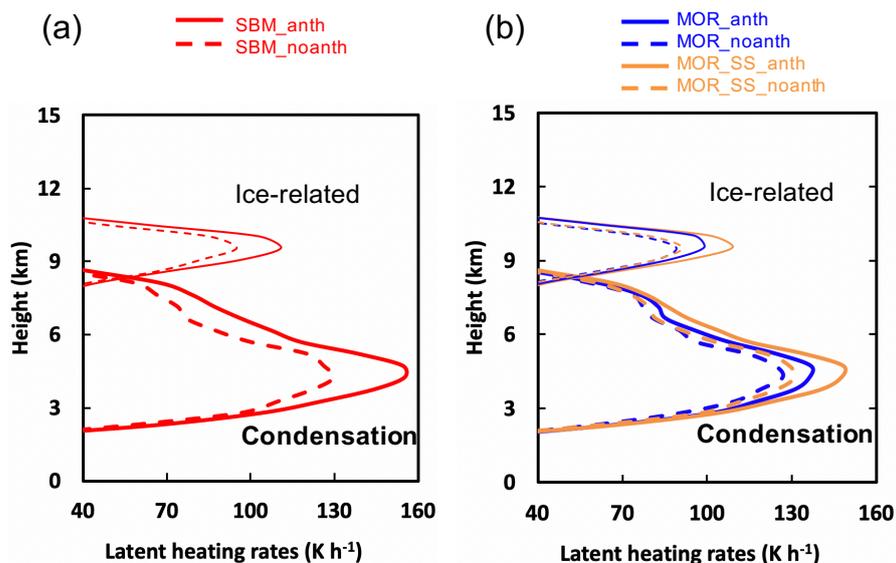
682 **Figure 10** CFADs of updraft velocity (unit:  $\text{m s}^{-1}$ ) for values larger than  $2 \text{ m s}^{-1}$  from (a)  
683 SBM\_noanth, (b) SBM\_anth, (c) MOR\_noanth, and (d) MOR\_anth over the study area (red box  
684 in Fig. 7) during the strong convection period (2000 – 2300 UTC, 19 Jun 2013).



685  
 686 **Figure 11** Vertical profiles of (a,b) updraft velocity (unit:  $\text{m s}^{-1}$ ), (c,d) thermal buoyancy (unit:  
 687  $\text{m s}^{-2}$ ) and (e,f) total latent heating rate (unit:  $\text{K h}^{-1}$ ) averaged over the top 25 percentiles (i.e.,  
 688 from 75<sup>th</sup> to 100<sup>th</sup>) of the updrafts with velocity greater than  $2 \text{ m s}^{-1}$  from the simulations  
 689 SBM\_anth and SBM\_noanth (red), MOR\_anth and MOR\_noanth (blue), and MOR\_SS\_anth  
 690 and MOR\_SS\_noanth (orange) over the study area (red box in Fig. 7) during the strong



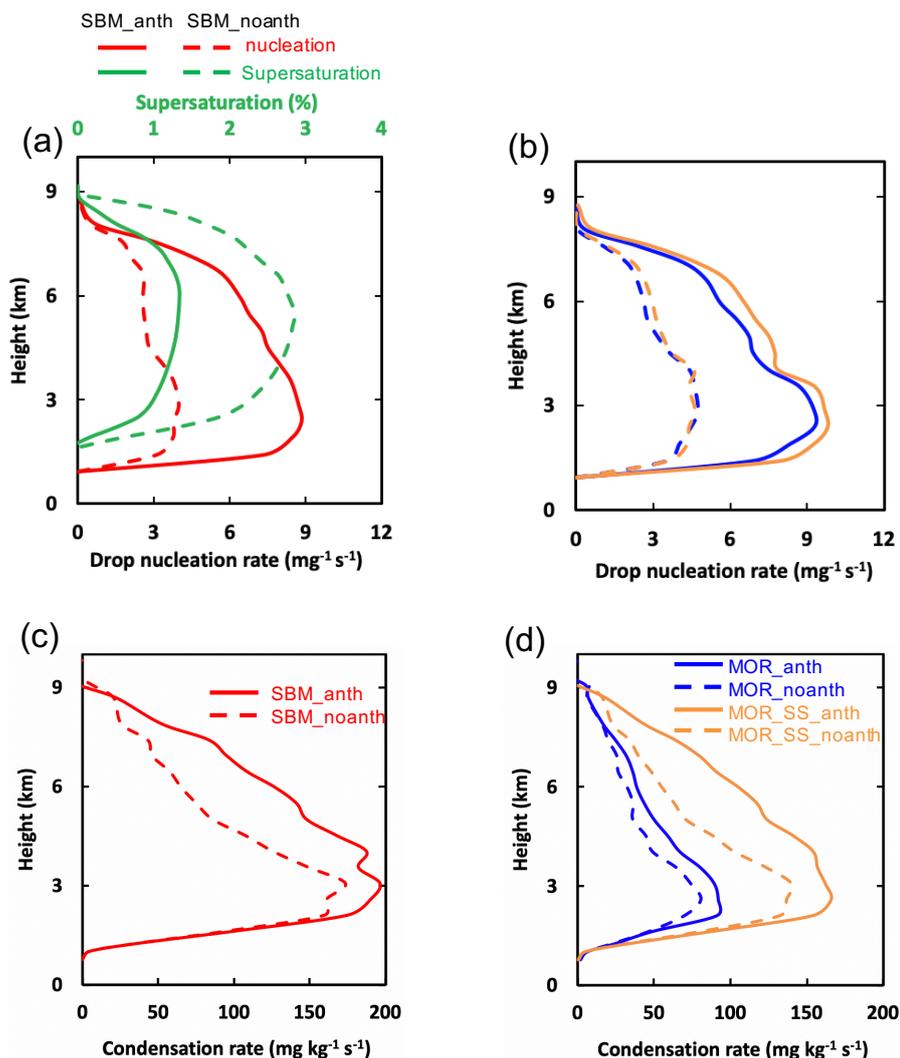
691 convection period (2000 – 2300 UTC, 19 Jun 2013). The dotted lines in (a) and (b) denote the  
692 freezing level (0 °C) and homogeneous freezing level (-40 °C).  
693



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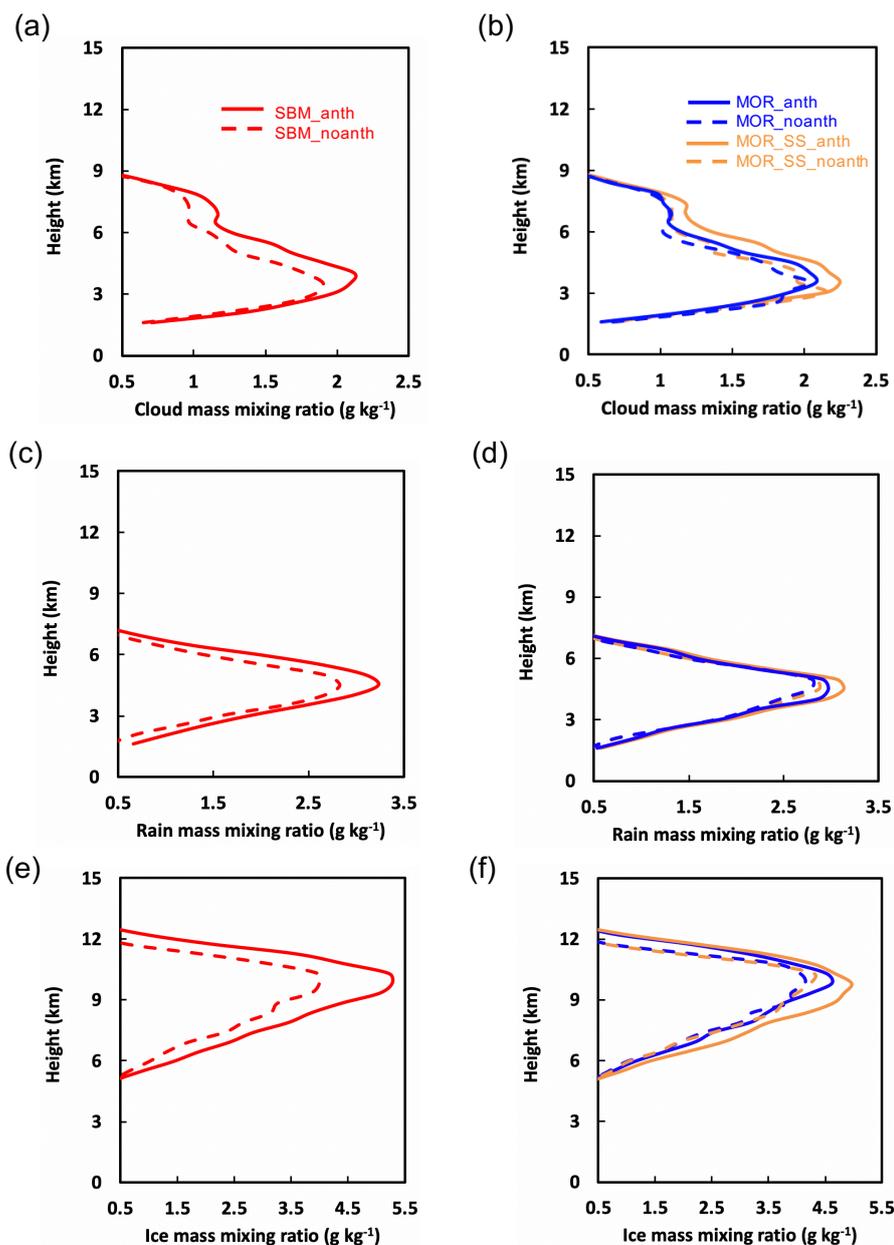
695 **Figure 12** Vertical profiles of condensation heating rate (thick lines below 9 km; unit: K h<sup>-1</sup>) and  
696 ice-related latent heating rate (thin lines above 9 km; unit: K h<sup>-1</sup>) averaged over the top 25  
697 percentiles (i.e., 75th to 100th) of the updrafts with velocity greater than 2 m s<sup>-1</sup> from the  
698 simulations (a) SBM\_anth and SBM\_noanth (red), and (b) MOR\_anth and MOR\_noanth (blue),  
699 and MOR\_SS\_anth and MOR\_SS\_noanth (orange) over the study area (red box in Fig. 7) during  
700 the strong convection period (2000 – 2300 UTC, 19 Jun 2013).

701



702

703 **Figure 13** Vertical profiles of (a) drop nucleation rate (red; unit:  $\text{mg}^{-1} \text{s}^{-1}$ ) and supersaturation  
704 with respect to water (green; unit: %) from SBM\_anth and SBM\_noanth, (b) drop nucleation rate  
705 (unit:  $\text{mg}^{-1} \text{s}^{-1}$ ) from MOR\_anth and MOR\_noanth (blue), and MOR\_SS\_anth and  
706 MOR\_SS\_noanth (orange), (c) condensation rate (unit:  $\text{mg kg}^{-1} \text{s}^{-1}$ ) from SBM\_anth and  
707 SBM\_noanth (red), and (d) the same as (c) but from MOR\_anth and MOR\_noanth (blue), and  
708 MOR\_SS\_anth and MOR\_SS\_noanth (orange), averaged over the top 25 percentiles (i.e., from  
709 75<sup>th</sup> to 100<sup>th</sup>) of the updrafts with velocity greater than  $2 \text{ m s}^{-1}$  over the study area (red box in  
710 Fig. 7) during the strong convection period (2000 – 2300 UTC, 19 Jun 2013).



711

712 **Figure 14** Vertical profiles of (a, b) cloud droplet, (c, d) rain drop and (e, f) ice particle (including  
713 ice, snow, and graupel) mass mixing ratios (unit:  $\text{g kg}^{-1}$ ) averaged over the top 25 percentiles (i.e.,  
714 75th to 100th) of the updrafts with value greater than  $2 \text{ m s}^{-1}$  from the simulations SBM\_anth and  
715 SBM\_noanth (red), MOR\_anth and MOR\_noanth (blue), and MOR\_SS\_anth and



716 MOR\_SS\_noanth (orange) over the study area (red box in Fig. 7) during the strong convection  
717 period (2000 – 2300 UTC, 19 Jun 2013).  
718