I believe that the revised version has improved, but I found the responses and the subsequent modifications introduced in the revised version still a bit qualitative and not very elaborated. I believe the authors have a complete data set that it is interesting for the reader of ACP, but they do not fully exploit it with a more thorough analysis.

As mentioned in my first revision, it will be very interesting to quantify some of the terms, in particular the role of clouds and entrainment in the mixing and transport of O3 and CO2. I believe this can reinforce the measurements interpretation. I realize that this can be “rough” first-order estimations, in absence to a more complete model study, but I consider that these estimations are an adding value to the paper. Below, I include some specific suggestions.

Taking Equation 5 from Ouwersloot et al. (2013) or Eq.4 at Karl et al. (2007, Journal of Geophysical Research-Atmosphere) as a guidance, and applying it for CO2 and O3 budgets, the entrainment flux for ozone and carbon dioxide can be roughly estimated from the data by combing information of the boundary layer growth presented at Figure 6 with the specific characteristics of CO2 and O3 (figures 7 and 8). This vale of the entrainment flux for ozone can be then compared with an estimation of the deposition flux calculated using the surface measurements of ozone and a characteristic deposition velocity. Similarly, one can obtain a first guess of the contribution of the mass flux that mixes CO2 and O3 from the boundary layer into the free troposphere. In doing so, they could provide some quantification to the new statements regarding the role of clouds and the mass flux (4th paragraph page 17) and the magnitude of entrainment (6th paragraph page 18).

We thank the reviewer for his specific suggestions and his help to improve our study. We agree that the shallow cumulus clouds play a major role for the mass flux and associated tracer fluxes when analyzing vertical transport of tracers, since they might dominate vertical mass flux.

We followed the guidance by the reviewer according to the suggested literature. First, we used equation 6 from Ouwersloot et al. (2013) and rearranged-ordered it (EQ 1) to calculate the mass flux (expressed as velocity) for CO2 and O3. The entrainment velocity ($w_e$) is approximated by the change of the PBL height and the tracer gradient across the entrainment zone. Having estimated the mass flux velocity $M$, we can calculate an entrainment velocity based on eq.5 from Ouwersloot et al., 2013 (here EQ 2):

\[
\text{EQ 1: } \frac{\mathbf{M}}{\Delta \Phi} = \left( \frac{\Phi_{cc} - \Phi_{mix}}{\Phi_{mix}} \right)
\]

\[
\text{EQ 2: } \frac{\mathbf{M}}{\Delta \Phi} = \left( \frac{\Phi_{cc} + \Phi_{mix} - \Phi_{mix}}{\Phi_{mix}} \right)
\]

The first term on the right hand side denotes the surface flux, the second term denotes the entrainment flux and the third term in the numerator is the temporal change of the mixing ratio (or more general any scalar quantity) within the boundary layer. The denominator includes the differences between the scalar ratio in the cloud core ($\Phi_{cc}$) and the mean in the PBL ($\Phi_{mix}$, EQ 2) or the scalar at the top of the PBL layer ($\Phi (h^*)$, EQ 1). To solve both equations it is necessary to estimate the mixing ratio within the cloud core ($\Phi_{cc}$). As mentioned before, no data inside the cloud were measured, therefore we defined the layer where we observed the mixing line as cloud layer and averaged the CO2 and O3 profiles. The ozone
surface flux was calculated with typical ozone deposition velocity of \( v_d = 0.1 - 0.5 \text{ cm/s} \) (Zhu, 2015), which leads to an ozone surface flux of \( \frac{w' \bar{O}_3}{\partial t} \big|_0 = -0.04 \) to -0.2 ppb m/s, where the average between both values is used. The CO2 surface flux of \( \frac{w' \bar{CO}_2}{\partial t} \big|_0 = -0.2 \) to -0.4 ppm m/s was measured by Vilà-Guerau de Arellano et al (2004), with fair weather conditions and clouds development after 14 UTC, again the average between both values is used. The temporal change of each scalar \((\frac{d\Phi}{dt})\) is calculated from the averaged scalar between 1000 to 1600 m at 7:30 UTC and 11:45 UTC. All values are summarized in the table below:

<table>
<thead>
<tr>
<th></th>
<th>O3</th>
<th>CO2</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>EQ 1</td>
<td>EQ 2</td>
</tr>
<tr>
<td>(w'\Phi' \big</td>
<td>_0)</td>
<td>-0.12 ppb m/s</td>
</tr>
<tr>
<td>(\Delta\Phi)</td>
<td>0.04</td>
<td>2.8 ppm</td>
</tr>
<tr>
<td>(\frac{dh}{dt} \text{ [m/s]})</td>
<td>....</td>
<td>0.04</td>
</tr>
<tr>
<td>(w_i \text{ [m/s]})</td>
<td>-0.08</td>
<td></td>
</tr>
<tr>
<td>(\frac{d\Phi}{dt})</td>
<td>0.00019 ppb m/s</td>
<td>0.00016 ppm m/s</td>
</tr>
<tr>
<td>(h \text{ [m]})</td>
<td>800</td>
<td></td>
</tr>
<tr>
<td>(\Phi_{cc})</td>
<td>44.5 ppb</td>
<td>384.3 ppm</td>
</tr>
<tr>
<td>(&lt;\Phi&gt;)</td>
<td>41.2 ppb</td>
<td>382.9 ppm</td>
</tr>
<tr>
<td>(\Phi(h^+\big</td>
<td>))</td>
<td>38.9 ppb</td>
</tr>
<tr>
<td>(M \text{ [m/s]} \text{ (EQ2)})</td>
<td>....</td>
<td>-0.0172</td>
</tr>
</tbody>
</table>

Calculation of the entrainment velocity based on EQ 4 in Ouwersloot et al (2013).

\[ w_e = \frac{dh}{dt} - w_s + M \]

<table>
<thead>
<tr>
<th></th>
<th>O3</th>
<th>CO2</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>EQ 4</td>
<td>EQ 1</td>
</tr>
<tr>
<td>(w_e)</td>
<td>0.2506</td>
<td>0.23</td>
</tr>
<tr>
<td>(M \text{ [m/s]} \text{ (EQ1)})</td>
<td>....</td>
<td>0.28</td>
</tr>
</tbody>
</table>

The mass flux velocity obtained from the ozone measurements is -0.02 m/s. The negative mass flux most likely reflects the missing ozone production in our budget estimate. Fig. 13 shows, that the surface ozone increased at the surface during the morning from 35 ppb to 45 ppb. This is most likely due to moderate ozone production during the morning, which partly may balance the ozone deposition flux. Accounting for a net surface ozone increase by a positive flux, the ozone budget becomes positive in accordance with the CO2 budget. The mass flux velocity estimate from the CO2 measurements is from \( M = 0.13 \text{ m/s} \), which fits to the mass flux velocity of 0.08 m/s reported in Karl et al. 2013, and 0.06 m/s by Ouwersloot et al. (2013). Since this value is better constrained from our observations than the ozone derived values, which are affected by chemistry, we use \( M(\text{CO2}) \) in the following.

We calculate an entrainment velocity \((w_e)\) with EQ 4 from Ouwersloot et al. (2013) by accounting for the mass flux \((M)\) from EQ2 (modified entrainment velocity in the following to distinguish from an entrainment velocity, which does not account for the mass flux velocity). This modified entrainment velocity is about 0.25 m/s for CO2, and it is larger than the mass flux velocity by an amount of 50% of the
entrainment velocity, which depends on the mass flux (based on the CO2 budget). We used this entrainment velocity with Eq 5 from Ouwersloot, et al (here Eq 1) to recalculate the mass flux velocity. For this CO2 – based velocity the recalculated mass fluxes for both, ozone and CO2 yield 0.23 to 0.28 m/s. The difference between the mass fluxes by Eq1 and Eq2 can be explained by the large uncertainty of the unknown mixing ratio in the cloud core (\(\Phi_{cc}\)), which is in the nominator of equation 5 and 6. However, the ratio between the CO2 entrainment flux (calculated with the modified entrainment velocity) and the surface flux (\(\overline{w'CO_2}\)) ranges from \(\beta_{CO2} = -0.2\) to \(-0.4\) ppm m/s) ranges from \(\beta_{CO2} = 1.8\) to \(3.5\), which is close to the ratio of \(\beta_{CO2} = 2.9\) by Vilà-Guerau de Arellano et al. (2004). Therefore, the cloud induced mass flux, which acts as a sink for the PBL leads to an increased modified entrainment velocity, which is about twice of the value without accounting for the additional cloud induced mass flux.

This illustrates the importance of the mass flux for the tracer budget of the PBL and agrees with the hypothesis that the mass flux might play an important impact on the vertical transport of atmospheric compounds from the PBL into the FT according to Ouwersloot et al. (2013), and that the descending air in between the clouds can be affected by air, which has been uplifted by convective motions.

We added to the text in section 4:

As shown by Ouwersloot et al. (2013) shallow cumulus clouds play a major role for the mass flux and associated tracer fluxes when analyzing vertical transport of tracers, since they might dominate vertical mass flux. Over the course of the day the cloud fraction increased from clear sky to 100% and constituted shallow cumulus clouds until noon. According to the ceilometer data and the soundings the aerosol/boundary layer height changed from 1000 m asl to 1640 m asl between 7:30 UTC and the second flight at 11:45 UTC, along with the cloud cover, which is reflected by the decrease of the global radiation in Fig.5. If we assume absence of chemical loss and production for the tracers, an upward mass flux would lead to an enhancement of CO2 from the nocturnal accumulation close to the surface accompanied by a decrease of ozone in the ascending branch of the subcloud layer. Such an upward mass flux would lead to tracer contrasts between the in-cloud and out cloud region as well as between the in-cloud and downward mixing region. As indicated by structured CO2 and profiles where the CO2 maxima are accompanied by high specific humidity, a cloud induced upward mass flux fits the observations of these peaks. In contrast the minima can be related to both the free tropospheric tracer composition of ozone, humidity and O3. Based on these gradients we estimated the entrainment velocity and the mass flux (expressed as velocity) according to equation 5 and 6 by Ouwersloot et al (2013). We find that the entrainment velocity (0.25ms\(^{-1}\)) is strongly increased by a cloud induced upward mass flux. Not accounting for the mass flux the entrainment velocity would be about 50% lower in the case of the CO2 budget. The largest uncertainties arise from the missing measurements of the surface fluxes and the estimates within the cloud layer, which was not penetrated from the aircraft.

I also found that the authors need to elaborate further on the role of subsidence or other large-scale forcing. By looking to the new figure 6, there is a large decrease in the mixed-layer at 6 UTC (almost 21000 as shown at Figure 7) to less than 1500 m at 10 UTC. I believe that this dynamic behaviour can only be explained by large scale forcing, either of synoptic or mesoscale character. Could the authors describe in a more elaborate manner this behaviour and placed in context for the O3 ad CO2 transport? Could they relate it to section 2.5?
The decrease of the residual layer height from 6 UTC to 10 UTC is also evident from the analysis data of the COSMO-7 model. This is a clear indication, that this behavior is related to large scale forcing and the adiabatic warming by large scale descending air. Trajectory analysis showed that the probed air masses above 2000 m (at 7:30 UTC) had no surface contact and descended from previously 6 km altitude to 2 km in the last 5 days. Thus the observation of descending air masses is consistent with higher ozone and low water mixing ratio, as observed above the inversion layer in the free troposphere.
Airborne observation of mixing across the entrainment zone during PARADE 2011

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Abstract

This study presents the analysis of the structure and air mass characteristics of the lower atmosphere during the field campaign PARADE (PArticles and RAdicals: Diel observations of the impact of urban and biogenic Emissions) on Mount Kleiner Feldberg in southwestern Germany during late summer 2011. We analysed measurements of meteorological variables (temperature, moisture, pressure, wind speed and direction) from radio soundings and of chemical tracers (carbon dioxide, ozone) from aircraft measurements. We focus on the thermodynamic and dynamic properties, that control the chemical distribution of atmospheric constituents in the boundary layer. We show that the evolution of tracer profiles of CO\textsubscript{2} and O\textsubscript{3} indicate mixing across the inversion layer (or entrainment zone). This finding is supported by the analysis of tracer–tracer correlations which are indicative for mixing and the relation of tracer profiles in relation to the evolution of the boundary layer height deduced from radio soundings. The study shows the relevance of entrainment processes for the lower troposphere in general and specifically that the tracer–tracer correlation method can be used to identify mixing and irreversible exchange processes across the inversion layer.

1 Introduction

The evolution of the planetary boundary layer (PBL) from stably stratified in the morning to well mixed in the afternoon is a multifaceted process and potentially leads to enhanced mixing between the PBL and the free troposphere (FT) (Stull, 1988). At sunset, a stable boundary layer with low turbulence intensity grows near the surface (Mahrt, 2014). Between this layer and the FT, a neutrally stratified layer remains resulting from the decay of turbulence of the convective PBL during daytime. This layer, called the residual layer, maintains the chemical composition of the original convective PBL (Stull, 1988). The turbulent-convective PBL and the overlaying FT are separated by the capping inversion, which often acts as a transport barrier. This barrier, which is indicated by changes of the vertical gradients
of temperature, humidity, aerosol content and chemical constituents, can be overcome by frontal activity (Sinclair et al., 2010). Frontal systems ahead of the respective cold and warm sectors of a cyclonic depression lead to vertical exchange between the PBL and the FT and potentially mix both air masses (Donnell et al., 2001). This is also the case for shallow and deep convection (e.g. Flaty et al., 1995; Hauf et al., 1995). The so called entrainment zone (EZ) is a layer of intermittent turbulence at the top of the PBL where air masses from the FT are entrained into the capping inversion and thus can interact with convective thermals from the PBL (e.g. Angevine, 2007; Bange et al., 2007). Based on LIDAR (Träumner et al., 2011), wind profiler (Cohn and Angevine, 2000) and ceilometer measurements (Eresmaa et al., 2006), the extent of the EZ was estimated to 20–40 % of the PBL depth (Martin et al., 2014).

Airborne flux measurements of heat and moisture were used to determine the impact of entrainment on their budgets (e.g. Driedonks and Tennekes, 1984; Canut et al., 2010), and to understand the influence of entrainment on the development of boundary layer clouds (Stevens et al., 2003) and the decay of stratocumulus clouds (e.g. Lilly, 1968). Malinowski et al. (2013) showed that turbulent mixing between the cloud top and the FT depends on the thermodynamic properties (static stability) and wind shear (dynamic instability). Moreover, radiative effects can change the entrainment flux at the cloud tops and lead to exchange and mixing as well between the FT and PBL (Moeng et al., 1999). Modelling studies revealed that changes of the (static) stability, low-level wind shear, and turbulent motions affect not only the growth of the boundary layer but also the entrainment at the PBL top (e.g. Driedonks and Tennekes, 1984; Sullivan et al., 1998; Cohen, 2000). Lock et al. (2000) showed that mixing and entrainment at the top of the PBL depend on the surface buoyancy flux, the inversion strength and on the influence of wind shear (Conzemius and Fedorovich, 2006).

In recent years substantial efforts have been made to understand the driving physical processes of temperature and humidity transport during the growth or decay of the PBL as summarized by Lothon et al. (2014). It is still difficult to quantify the role of the EZ for the mixing of trace gases, because of its high spatial, vertical and temporal variability (Martin 2014).
et al., 2014). Vila-Guerau de Arellano (2004) was able to show that the downward transport of CO$_2$ from the FT into the PBL can enhance or reduce CO$_2$ mixing ratios within the PBL which then can lead to a misinterpretation of possible CO$_2$ sources. Vertical gradients of ozone exist as well between the PBL and FT, and they are strongly affected by transport processes and chemical reactions. Measurements and modelling studies of O$_3$ revealed that convective transport of ozone precursors lead to chemical ozone formation in the FT and enhances the background concentrations (Hov and Flatøy, 1997; Henne et al., 2005).

Cooper et al. (2015) highlighted that surface air quality regulations in the US strongly depend on accurate knowledge of baseline ozone, which is defined as ozone transported downwind from anthropogenic and natural sources, while forecast of baseline ozone needs large knowledge of all different transport processes. This includes intercontinental transport from Asia, which is often associated with episodes of enhanced ozone concentrations over large parts of the western US that is located at relatively high altitudes (Lin et al., 2012). Moreover, downward transport of ozone rich air from the FT can enhance ozone mixing ratio in the PBL (Beck et al., 1997; McKendry and Lundgren, 2000; Jaffe, 2011).

Neuman et al. (2012) used tracer–tracer (CO–O$_3$) correlations to identify the in-mixing of ozone from the FT into the PBL. This approach is commonly used to identify stratosphere–troposphere exchange (e.g., Fischer et al., 2000; Hoor, 2002). Generally the approach makes the assumption of two well-mixed reservoirs, which give rise to linear tracer–tracer relationships in the case of mixing. However, the boundary layer acts as a barrier for mixing and often exhibits strong trace gas gradients. Therefore the distribution of data points in tracer–tracer scatter plots subdivide into distinct groups depending on the mixing ratios of the respective well mixed reservoirs. Mixing between these air masses leads to linear connections between these groups (mixing lines).

Handisides (2001) presented ozone and peroxy radical observations from the local valley and from Mount Kleiner Feldberg and linked his findings to transport within different boundary layer structures in summertime. Moreover during winter, Wetter (1998) observed on the top of Mount Kleiner Feldberg CO plumes transported out of the stable boundary layer below the summit. In both studies the PBL structure was not investigated in detail. However,
this is necessary to understand and interpret the temporal evolution of the spatial distribution of atmospheric trace species measured within and above the PBL, especially over complex terrain where handover processes between the PBL and the FT are additionally triggered by thermally induced flow systems (Kossmann et al., 1999; Henne et al., 2004).

In this study, we will use CO₂ and O₃ to study mixing across the EZ with the tracer–tracer correlation method. Both trace species have been measured by aircraft during the field campaign PARADE at Mount Kleiner Feldberg in late summer 2011, and are used for the first time together to identify mixing and transport within this atmospheric region. Additionally we discuss possible processes which lead to mixing across the entrainment zone. In this context we focus on the boundary layer development and the weather conditions.

We will first present the synoptic condition during the day of measurements and put it into perspective with the entire situation of the measurement campaign. We focus on the diurnal evolution of the boundary layer which we determined with balloon and ceilometer measurements. We then present the results of the aircraft trace species measurements. Finally, we discuss possible processes which may be responsible for the measured distribution of CO₂ and O₃ and its influence from local mixing and long-range transport.

2 Measurement and methods

2.1 PARADE 2011

The PARADE (PArticles and RAdicals: Diel observations of the impact of urban and biogenic Emissions) field campaign took place in August and September 2011 located at the Taunus Observatory at the summit of Mount Kleiner Feldberg (825 m a.s.l.), about 20 km northwest of Frankfurt am Main in southwestern Germany. The goal of PARADE was to study the impact of summertime biogenic and urban emissions in a semi-remote environment. A comprehensive set of measurements of several reactive species and meteorological variables were conducted during PARADE, also including in situ measurement of ozone, and carbon dioxide (Phillips et al., 2012; Bonn et al., 2014; Li et al., 2015).
Mount Kleiner Feldberg is mostly covered with coniferous forest, except on the top where the observations took place during PARADE. Meteorological measurements for temperature, wind, humidity, global radiation are performed by a permanently operating observatory of the German Weather Service (DWD) and of the State of Hessen environmental agency (HLUG) \cite{Crowley2010}. The station is known for its quite remote character under prevailing northwesterly winds \cite{Klein2010}. However, this rural region is affected by pollution from the Rhein-Main area (including the cities of Frankfurt/Main, Wiesbaden, and Mainz) when wind directions have a southerly component. The area north of Mount Mount Kleiner Feldberg (50–100 km) is less populated and without major industry. In the direct vicinity of Mount Kleiner Feldberg are few main roads, some small towns within 5 km, and two similar sized mountain peaks are located: Großer Feldberg (878 m) and Altkönig (798 m).

During this field campaign the synoptic conditions involved predominately south westerly winds, mainly related to troughs over the northeastern Atlantic and to high pressure systems over Central Europe. The south westerly winds advected enhanced levels of anthropogenic pollution to Kleiner Feldberg and were interrupted by northern to north easterly winds several times. In these cases the probed air masses showed rather low levels of anthropogenic pollution \cite{Phillips2012}. We will present additional details on the synoptic conditions below.

### 2.2 Measurements

#### 2.2.1 Aircraft instrumentation

A P68D Partenavia turbo prop aircraft operated by Enviscope GmbH (www.enviscope.de, last access: 12 November 2014) was equipped with several instruments to perform measurements in and above the planetary boundary layer on three days (31 August, 2 September, 6 September 2011) during the campaign. The aircraft flew 5 research flights from Mainz to Mount Kleiner Feldberg. Ascending profiles were flown up and downstream of the surface station, while descending profiles were performed directly over Mount Kleiner
Feldberg. The ascending flight legs were conducted in different regions and are not included in the averaged profiles of the following analysis. The descending flight legs were flown in spirals with a diameter of about 10 to 15 km. The maximum flight altitude was 3000 m asl and the flight time was approximately two to three hours. A typical flight track is shown in Fig. 1 and illustrates the sampling strategy. In total, 13 descent profiles were obtained at different locations and times for each variable.

CO₂ was measured with a modified commercial CO₂/H₂O analyser (LI-6262, LI-COR Inc.) with a total uncertainty of 0.55 ppmv (= one standard deviation σ) determined from regular in-flight calibrations and a time resolution of 1 Hz [Gurk et al., 2008]. Ozone (O₃) was measured by UV absorption every 10 seconds with the commercial gas analyser O342M (Environnement S.A.) with a total uncertainty of 1.3 ppbv for a moist air mass with less than 10 g kg⁻¹ specific humidity [Köllner, 2013]. The response time of the instruments (5%/95%) was 3 s (CO₂) and 9 s (O₃) and the final data correction accounted for different residence times in the inlet lines. The ambient temperature and moisture were measured every second with a commercial HMP230 (Vaisala) with an accuracy of two to three percent for relative humidity and 0.5 °C for temperature. The ambient pressure (SETRA Model 270) was measured with accuracy of 0.55 hPa (pers. com., Enviscope GmbH). The particle number concentration was measured by the Optical Particle Counter (Grimm SKY-OPC 1.129) for a size range of 0.25–32 µm (32 channels) and a time resolution of 1 Hz . Additionally, particle number concentrations of particles with diameters between 10 nm and 5 µm were measured with a modified condensation particle counter (CPC-Modell 3007, TSI) every second. These measurements are reported for standard temperature (273.15 K) and pressure (1013.25 hPa) conditions (STP).

Three-second average measurements are analysed corresponding to a spatial resolution of about 150 m at an average aircraft velocity of 50 m s⁻¹. The aircraft descent took about 10 min, so that the 3 s data yielded a vertical resolution of three to four meters. For further analysis of the vertical structure of the trace gases measurements between 1000 and 3000 m a.s.l. Kleiner Feldberg (approx. 150 m over the mountain summit), each vertical profile is linearly interpolated to a uniform vertical grid with 10 m resolution.
2.2.2 Surface-based and radiosonde measurements

Vertical profiles of temperature, relative humidity and position were obtained from 174 radiosondes (GRAW DFM-06, [www.graw.de](http://www.graw.de), last access: 12 November 2014). Four to ten radiosondes were launched each day, starting usually one hour before sunrise and ending one hour after sunset. The sampling rate of 1 Hz results in a vertical resolution of about 1–2 m. The data has been interpolated onto an equidistant vertical grid with 20 m resolution. The pressure is obtained from these observations. Wind speed and wind direction are calculated from temporal changes of the current position of the balloon relative to the surface station. The near surface temperature measurements were compared to the temperature measurement from the local observatory. The comparison revealed a high positive correlation with an $r^2 = 0.99$ and a mean absolute deviation of $|T_{RS} - T_{DWD}| = 0.56 \pm 0.50^\circ C$.

The Ceilometer CHM 15 k-x from Jenoptik ([www.jenoptik.com](http://www.jenoptik.com), last access: 12.11.2014) is primarily designed to determine the cloudbase from 200 m to 15 km above ground. In the absence of clouds, the backscatter signal can also be used for detecting aerosol layers ([Münkel et al., 2007](http://www.jenoptik.com)). The laser emits a light pulse with a repetition rate of 5–7 kHz, a pulse time of 1 ns and a wavelength of $\lambda = 1064$ nm, with a vertical resolution of 15 m. Thus, every 30 s a vertical backscatter profile is obtained.

Surface based $CO_2$ measurements were performed with a commercial $CO_2/H_2O$ analyser (LI-6262, LI-COR Inc.) with a total uncertainty of 2 ppmv determined from regular calibrations and a time resolution of 1 Hz. Additionally $CO_2$ was measured from the mobile laboratory (MoLa) ([Drewnick et al., 2012](http://www.jenoptik.com)) with a commercial $CO_2/H_2O$ analyser (LI-840, LI-COR Inc.) without regular calibrations and a separate inlet line. This data was provided as 10 min mean values and is only used here to illustrate the diurnal cycle of $CO_2$. Surface based ozone measurements were performed by various instruments with different techniques at different locations during this campaign. At Mount Kleiner Feldberg $O_3$ was measured by UV absorption with a commercial $O_3$ analyser (Model 49, Thermo Environmental Instruments Inc.). The total uncertainty was 5% (two $\sigma$) and the data were provided as 10
minute averages (Li et al., 2015). HLUG provides 30 minute averaged O$_3$ measurements using a commercial UV absorption instrument (API 400) (Crowley et al., 2010).

### 2.3 Determination of the planetary boundary layer height

The planetary boundary layer height (PBLH) was calculated from the radio soundings using the surface based bulk Richardson number $Ri_B$ (Vogelezang and Holtslag, 1996; Seibert et al., 2000). The dimensionless bulk Richardson number $Ri_B$ is defined as the ratio of the vertical stability and the vertical shear of the horizontal wind:

$$ Ri_B = \frac{g}{\theta_v} \frac{\Delta \theta_v \Delta z}{(\Delta u)^2 + (\Delta v)^2} $$

with $g$ the gravitational acceleration, $\theta_v$ the virtual potential temperature, $z$ the altitude, and $u$, $v$ the horizontal wind components (Vogelezang and Holtslag, 1996). The differences $\Delta$ are obtained from the respective values at the surface and the layers above. $Ri_B$ provides an appropriate measure to characterize the stratification of the atmosphere and to determine the PBLH. The thermal stratification is determined from the sign of the bulk Richardson number. The flow is statically stable if $Ri_B$ is greater than zero, else statically unstable or neutral. Furthermore, if $Ri_B$ is below a critical value, the flow will become dynamically unstable and marks the altitude between the turbulent boundary layer and the stable stratified free troposphere. We used a threshold value of $Ri_B = 0.38$ which distinguish between a turbulent (low values of $Ri_B$) and a stable (large values of $Ri_B$) boundary layer. The same value is also used in the mesoscale model that we will introduce later (Fay, 1998). For clarity, the threshold value has been varied between 0.2 and 0.5 following suggestions from Zilitinkevich and Baklanov (2002) and Zhang et al. (2014), however, with a negligible effect on the results.

The PBLH has further been determined with the ceilometer, following the approach from Schneider and Eixmann (2002), to observe the height of the aerosol layer more frequently than with the radio soundings. The method is based on iterative averaging of the backscatter profile. Once the mean value has decreased to half the maximum average value inside the
PBL, this altitude is designated to be the PBLH. Additionally, the height of the residual layer can be determined with this method at night. For our analysis we excluded all backscatter profiles which were observed during cloudy conditions. The determined boundary layers of the remaining backscatter profiles are averaged on an hourly basis.

2.4 **Lagrangian air mass analysis**

For the surface observations 24 h backward trajectories were calculated using LAGRANTO \cite{Wernli1997, Sprenger2015} from the position of Mount Kleiner Feldberg at 200 m above surface level (approximately between the modeled and actual mountain height) at 12:00 UTC for each day of the campaign. For the aircraft observations 24 h backward trajectories were calculated every 10 s along the flight path for each individual flight using LAGRANTO. In both cases, the trajectories were driven by model wind fields from the COSMO-7 analysis \cite{Doms2011, Doms2011}, operated by Swiss National Weather Service (MeteoSwiss). The meteorological fields from COSMO-7 are available every hour, the horizontal resolution is about 7 km, and 61 vertical levels are used from the surface up to 20 hPa. The trajectories are calculated from the three dimensional, kinematic wind field, not taking into account additional parametrisations for convection or turbulence.

2.5 **Synoptic conditions during PARADE 2011**

The synoptic conditions were alternating mainly between two kinds of weather patterns. Prior to the start of the measurement period (15–18 August 2011) the situation was dominated by a low pressure system over northwestern Europe and a high pressure system over central Europe. On 19 August 2011 a cold front related to a low pressure system over the north sea passed over the measurement area. This was then followed by high pressure conditions over central Europe (20–26 August 2011) with very warm and humid temperatures. This period was again succeeded by the passage of a low pressure system late in the day on 26 August 2011 which formed over the eastern Atlantic. This led to relatively low temperatures with \( T_{\text{max}} < 14^\circ \text{C} \) and continued until the end of August (31 August 2011).
Similar conditions were found during the last days of the campaign, only interrupted by fair weather between the end of August and 4 September 2011.

This synoptic pattern is well reflected by the daily mean carbon dioxide (CO₂) mixing ratio observed at the surface for the entire campaign (Fig. 2). CO₂ showed larger values during high pressure conditions (gray area), while lower CO₂ values were observed during low pressure conditions (white area). Figure 3 shows that the near surface based LAGRANTO backward trajectories match well with these synoptic conditions. Clean air masses from the north-west to westerly direction over the North Atlantic are linked to low CO₂ mixing ratios, while enhanced CO₂ values are related to the potentially polluted continental regions in southerly wind directions or related to local fair weather conditions. Note that biogenic uptake of CO₂ through photosynthesis was of minor importance at the end of the European growing season and that the CO₂ flux was probably dominated by anthropogenic emissions (Rivier et al., 2010).

3 Results

3.1 Synoptic situation on 6 September 2011

In the following we focus on one particular day (6 September 2011) when two research flights were completed. The two flights were flown in a prefrontal high pressure zone with a cyclonic system in northern Europe and high pressure in the south (Fig. 4). The approaching frontal system was associated with high wind speeds over Central Europe, carrying air masses rapidly from the Atlantic Ocean to Central Europe. A first trough reached western Europe on 5 September 2011, and marked the beginning of a series of low pressure systems. The next low was already present over the British Isles. Its cold front was semi-occluded with the prefrontal warm front. Part of this warm front was situated over northwestern France (Fig. 4). The strong westerly wind transported colder air masses from the Atlantic Ocean and fostered the evolution of small-scale convective cells.
3.2 Diurnal evolution of the boundary layer on 6 September 2011

In this section, we briefly discuss the diurnal evolution of the PBL on 6 September 2011. We show the temporal evolution of $R_i_B$ over Mount Kleiner Feldberg (Fig. 5). The vertical profiles of $R_i_B$ have been calculated solely from the radio soundings and were linearly interpolated between each profile in this figure. In the morning hours low values of $R_i_B$ are confined to the lowest layers (so called stable boundary layer). The extent of these low values increases during the day up to 2000 m a.s.l. This reflects the growth of the PBL. In the morning hours a second inversion layer was indicated by the radiosounding profiles at around 2000–2500 m a.s.l.. This inversion further coincides with an increased backscatter signal detected by the ceilometer. The backscatter signal is related to an aerosol layer as well as the bottom of shallow cumulus clouds and marks the altitude of the residual layer.

The stable and residual layers merged between 10:00 and 11:00 UTC, afterwards only a single layered structure (now referred to as planetary boundary layer (PBL)) is detectable. Figure 6 shows profiles of the virtual potential temperature, wind speed and wind direction from the radio soundings just before and after the flights during that day. The PBLH at 11:00 UTC is located at 1640 m a.s.l. with an inversion strength of about 3 K. The maximum PBLH of 1981 m a.s.l. occurred at 14:00 UTC, while the strengthening of the temperature inversion increased to about 10 K. The wind direction within the PBL and FT was mostly constant during that day, whereas the wind speed increased within the FT and decreased within the PBL during the day. Later on, the PBLH decreased to about 500 m before the onset of precipitation at around 18:00 UTC (Fig. 5).

The air mass composition was probed twice during this day. The first flight took place between 07:00 and 09:30 UTC, while the second flight was between 11:00 and 14:00 UTC. Thus, a composition change of the PBL was expected from flight one to flight two, since the first flight took place before the two layers merged, while the second occurred after this event.
3.3 Vertical profiles of trace gases during the boundary layer growth

Figure 7 also shows CO₂ profiles for the morning (07:27 UTC), at noon (11:42 UTC) and afternoon (13:12 UTC). The profiles are color-coded with specific humidity along the flight leg. The PBLH (black dashed line) and residual layer height (gray line), derived from the gradient of the humidity profiles, correspond well to the derived boundary layer heights from the radiosonde measurements.

In the morning the vertical profiles of CO₂ clearly show the three-layer structure of the PBL (Fig. 7a). Within the stable layer we observed a large variability of CO₂ mixing ratios ranging between 380.0 ppmv and 394.3 ppmv below 1000 m altitude. Between 1000 and 2300 m a.s.l. the vertical profile of CO₂ in the residual layer was almost constant with mean values (\(\overline{\text{CO}_2}\)) of 380.5 (±0.16) ppmv, indicating that the residual layer was well mixed. A strong CO₂ gradient at 2300 m a.s.l. reveals a sharp separation between the PBL and the FT which has a mean CO₂ of about 383.0 (±0.58) ppmv. In the afternoon (13:12 UTC) CO₂ is again well mixed within the PBL with \(\overline{\text{CO}_2}\) = 382.1 (±0.35) ppmv (see Fig. 5). A clear separation to the FT is evident again, with mean FT-CO₂ mixing ratios of 383.7 (±0.38) ppmv.

The transition period between the measurements of these two profiles is characterized by a much more variable CO₂ profile, especially above the PBLH (Fig. 7b). CO₂ decreases within the PBL with a mean value of 382.9 (±0.49) ppmv. Above 1500 m a.s.l. in the FT mean CO₂ mixing ratios show a double peaked structure with maximum values of almost 386 ppmv and minimum values of 382 ppmv. Thus, a large variability with altitude is evident. Above 2300 m a.s.l. CO₂ mixing ratios are more comparable again to mixing ratios observed in the morning and the afternoon with a mean value of 383.9 (±0.61) ppmv.

The vertical profiles of O₃ show a similar temporal evolution development (Figs. 7 and 8). Largest O₃ values are found in the FT above the boundary layer throughout the day and exceed local PBL values by typically 10 ppbv. The only exception is the transition period at 11:42 UTC, when O₃ shows a large variability. As for CO₂ a double peak structure is evident, however, the maxima and minima are anti-correlated between CO₂ and O₃. Furthermore, O₃-minima and CO₂-maxima show large values of water vapor mixing ratio \(q_v\). These values
exceed the commonly measured humidity values (less than 2 g kg\(^{-1}\)) in the FT. Since the aircraft was allowed only to operate in cloud-free conditions, the large humidity values were only observed within the PBL. We will now analyse in detail how this complex structure is related to the observed development of the boundary layer structure.

### 3.4 Mixing deduced from correlations

The vertical profiles of O\(_3\) and CO\(_2\) and their temporal evolution over the course of the day, are used to study a potential mixing event on 6 September 2011 above Mount Kleiner Feldberg. For this we now apply tracer–tracer correlations (e.g., Fischer et al., 2000; Hoor, 2002). Figure 9 shows the scatter plot between O\(_3\) and CO\(_2\) in the morning (07:27 UTC). Using the moisture information and mean values of CO\(_2\) and O\(_3\) it is possible to identify two separated air masses. The air mass in the residual and stable layers has low O\(_3\) values and variable values of CO\(_2\) and \(q_v\). The air mass in the FT is rather characterized by larger values of O\(_3\), low values of \(q_v\) and less variability in CO\(_2\).

At 11:42 UTC the appearance of the scatter plot changes significantly (Fig. 10). The chemical characteristics of two different air masses still can be identified, of which one has the characteristics of a boundary layer air mass (larger values for \(q_v\)) and one of a FT air mass (lower values of \(q_v\)). However, the air mass in the FT shows now also a larger variability in CO\(_2\) which we see as an indication of a PBL air mass that has entered the FT. Moreover, the O\(_3\) mixing ratios in the PBL are also slightly increased which could be related to downward transport of an FT air mass.

Using now also the temporal and spatial information of the measured data points, a linear relationship between the FT and PBL air masses in the tracer–tracer scatter plot becomes evident (Fig. 10). This line is highlighted with the black boxes. These data points have been measured consecutively in an altitude range between 1640–2200 m, which is just above the PBLH at 11:42 UTC (Fig. 5). The EZ was situated at this location and time. Thus we regard these data points as the result of mixing, and the observed linear tracer–tracer relationship as mixing line, since it relates to the evolution of the boundary layer structure during this day. The specific humidity of this mixing line shows a slightly larger variability than the variability...
of the FT background (Fig. 8). This further supports our hypothesis of measuring a rather fresh mixing event. Note further that during this time clouds were present in the observation area with the cloud tops at around 2400 m a.s.l., according to the on board observations and a post flight analysis of photos taken during the flight. However, the strong variations of ozone, moisture and CO$_2$ for the profile at 11:42 (Fig. 7) have been observed exclusively in between clouds.

In short, we identify local mixing between air masses from the PBL and the FT from the results of the tracer tracer correlation technique for the trace gases CO$_2$ and O$_3$. Additionally, this mixing is supported by a larger variability of the specific humidity within the mixing region. In the following subsections, we will discuss processes which might have contributed to this mixing event as well as if this mixing event was influenced by long-range transport.

4 Discussion

4.1 Impact of long-range transport?

The origin of the probed air masses is investigated with 24 h kinematic backward trajectories using LAGRANTO (see Sect. 2.4) for this one-day period. Figure 11 shows the zonal position of the backward trajectories along the flight path at release time (0 h), −12 and −24 h prior to the measurements. Additionally, the trajectories are color-coded with the ozone mixing ratio as measured onboard the aircraft. Ozone was chosen, as tracer in this figure, because of its strong vertical gradient. The ozone mixing ratio is averaged for the trajectory positions backward in time for each hour in bins of 0.5° longitude and 250 m altitude. The binning compromises all calculated trajectories and reflects the transport of air for a time period 24 h before the observations for the second flight around noon. It is evident that the measurement region has been influenced by air masses from two different origins (Fig. 11). The trajectory analysis reveals that the air mass probed above 2000 m with large O$_3$ mixing ratios, has been transported rapidly from westerly directions towards the measurement site. Furthermore, only these trajectories reside within the FT with the high O$_3$
mixing ratios. The trajectories at lower levels indicate that the air mass at this altitude range with low $O_3$ mixing ratios, has been advected more slowly. Importantly, air masses initialized below 2000 m, thus mainly within the turbulent boundary layer, show a larger spread of the trajectories and tend to stay in the boundary layer corresponding to low ozone mixing ratios. A similar behaviour is evident from CO$_2$ (not shown) confirming a clear separation between free tropospheric and boundary layer transport.

We conclude that the entrainment zone separates ozone rich air masses, which were transported in the free troposphere with higher wind speed than from air masses, which resided in the boundary layer with lower wind speed, as indicated by the low ozone mixing ratios. The clear vertical separation of air masses between the FT and PBL is therefore reflected by the different chemical signature of the air masses and is an indication, that the two respective air masses were spatially separated. It is thus an indication that the air masses must have been mixed shortly before or at the measurement location.

### 4.2 Impact of the local boundary layer conditions and cloud development

We will now show, if local mixing between boundary layer air and the air in the free troposphere can explain the observed signatures of mixing. The PBLH reaches 1640 m a.s.l. with an inversion strength of 3 K at 11:00 UTC. The residual layer height in the morning, as identified from the ceilometer and secondary temperature inversions from the radio soundings, reaches 2100 m. From 11:42 to 13:12 UTC the PBLH grew to 2000 m a.s.l. and reaches a temperature inversion strength of 8 to 10 K. On the one hand, the growth of the boundary layer since 11:42 UTC is an indication for further entrainment of the overlying air masses into the PBL, which lead to further mixing. On the other hand, the increase of the inversion strength limits the entrainment from air above and also the upward transport to the FT and reduced mixing (Angevine, 2007).

As shown in Fig. 5 the boundary layer top during the measurements was affected by the occurrence of clouds. These clouds were aligned along straight lines (Fig. 12), which is also evident from satellite images (not shown). The presence of clouds can affect the strength of the inversion layer as well as of the mixing between PBL and FT (Stevens, 2007).
and Malinowski et al. (2011) showed that radiative cooling at the upper edge of the clouds promotes a downward mixing of air and thus influences the inversion strength and might also contribute to the observed mixing event by entrainment of dry air from the FT. This can be seen by the decrease of the mean specific humidity between 1640 and 2200 m a.s.l. from the morning ($q_v = 2.8 (\pm 1.2) \text{g kg}^{-1}$) to noon time ($q_v = 2.1 (\pm 0.7) \text{g kg}^{-1}$).

An important feature of these cloud structures is the formation of roll vortices, which are aligned along the axis of the cloud bands. The cloud bands can be regarded as an indicator of shallow convection within the boundary layer. Moist air is lifted from near the surface up to the top of the PBL where the clouds form and the air descends in-between the clouds. The indicated air motion contributes especially to a further mixing of the PBL air mass, but can also increase the mixing potential between the PBL and the FT. This can then result in the vertical profiles of CO$_2$ as observed between the morning and the afternoon on 6 September 2011 (see Fig. 6). Hägeli et al. (2000) observed similar relationships for temperature and for moisture profiles within the cloud free area between the cloud streets. It should be noted, that many experimental studies are restricted to cloud-free conditions as summarised by Kalthoff et al. (2013).

So we suggest that the mixing event is linked to the evolution of shallow convection and consecutive cloud formation at the top of the PBL. Therefore, we will discuss the three CO$_2$ profiles during this day (Fig. 7), which reflects the suspected dynamic processes from theoretical approaches (e.g. Etling and Brown 1993; Lock et al. 2000; Stevens et al. 2003). The descending air in-

As shown by Ouwersloot et al. (2013) shallow cumulus clouds play a major role for the mass flux and associated tracer fluxes when analyzing vertical transport of tracers, since they might dominate vertical mass flux. Over the course of the day the cloud fraction increased from clear sky to 100% and constituted shallow cumulus clouds until noon. According to the ceilometer data and the soundings the aerosol/boundary layer height changed from 1000 masl to 1640 masl between 7:30 UTC and the second flight at 11:45 UTC, along with the cloud cover, which is reflected by the decrease of the global radiation in Fig. 5. If we assume absence of chemical loss and production for the
tracers, an upward mass flux would lead to an enhancement of CO\textsubscript{2} from the nocturnal accumulation close to the surface accompanied by a decrease of ozone in the ascending branch of the subcloud layer. Such an upward mass flux would lead to tracer contrasts between the in-cloud and out cloud region as well as between the in-cloud and downward mixing region. As indicated by structured CO\textsubscript{2} and profiles where the CO\textsubscript{2} maxima are accompanied by high specific humidity, a cloud induced upward mass flux fits the observations of these peaks. In contrast the minima can be related to both the free tropospheric tracer composition of ozone, CO\textsubscript{2} and humidity. Based on these gradients we estimated the entrainment velocity and the mass flux (expressed as velocity) according to equation 5 and 6 by Ouwersloot et al. (2013). We find that the entrainment velocity (0.25 m s\textsuperscript{-1}) is strongly increased by a cloud induced upward mass flux. Not accounting for the mass flux the entrainment velocity would be about 50\% lower in the case of the clouds can be affected by air, which has been uplifted by convective motions. The mass flux dominates mostly the vertical transport of atmospheric compounds from the PBL into the FT (Vila-Guerau de Arellano, 2005; Ouwersloot et al., 2013). CO\textsubscript{2} budget. The largest uncertainties arise from the missing measurements of the surface fluxes and the estimates within the cloud layer, which was not penetrated from the aircraft.

As discussed before (Figs. 5 and 7), in the afternoon (11:42 UTC) the CO\textsubscript{2} profiles between 1600 to 2300 m show double peak CO\textsubscript{2}-mixing ratios around \(\sim 386\) ppmv, which were only present near the surface in the morning. Since there are no significant sources or sinks of CO\textsubscript{2} in the FT, these enhanced values indicate that surface air has been lifted locally and remains at this altitude when the cloud dissolved (Vila-Guerau de Arellano, 2005). This is also the case for specific humidity and ozone at this height, which exhibit exactly the same air mass composition as previously observed at the surface.

In principle, high ozone values can also be the result of photochemical ozone production, but ozone production would not add more than a few ppb per day (Morris et al., 2010). Ouwersloot et al. (2013) showed that ozone can be influenced by UV radiation extinction by clouds and atmospheric compounds like NO and NO\textsubscript{2}. However, the high ozone values are accompanied by low humidity, which indicates potential downward transport from higher al-
This is confirmed by 5 days backward trajectories (calculated with ECMWF analysis (0.5 °C horizontal resolution, 90 vertical layer, 3 h time resolution)), which show that these air masses reached a maximum height of about 6 km over the Atlantic ocean, without any surface contact (not shown). Finally, during the entire day the specific humidity in the free troposphere showed very low mean values of 1.4 to 2.0 g kg⁻¹. This indicates that the free troposphere was dominated by an air mass of similar chemical characteristics during the measurements on the 6 September 2011. Between 2200–2500 m a.s.l. the ozone rich air mass is also observed with higher humidity. The variations of O₃, CO₂ and qᵥ above 2000 m clearly indicate layers of different origin in the region between 1600 and 2300 m a.s.l. in the cloud interstitial regions between the roll clouds.

Therefore, the growth of the PBL is an indication of entrainment from the FT into the PBL. On the other hand the increase of the temperature inversion (cause from large-scale subsidence) provides strengthening of the inversion layer and limited the further boundary layer growth, but entrainment still occurs (Ouwersloot et al., 2013).

4.3 Impact of mixing on surface measurements

It has previously been shown that boundary layer air trapped within the stable layer can suddenly be mixed in the newly growing PBL as soon as the PBL becomes unstable (Stull, 1988). In addition, large scale synoptic situations can lift these air masses out of the PBL into the FT (Sinclair et al., 2010). Vice versa, polluted plumes from the FT can enter the PBL by downward mixing and potentially lead to an increase of the pollution levels near the surface in initially clean regions (McKendry and Lundgren, 2000; Neuman et al., 2012).

Figure 13 shows the ozone mixing ratios onboard the aircraft and the time series at the surface. The airborne measurements below 800 m a.s.l. agree well with the surface observations of about 30 ppb. The higher values around 07:30 UTC observed onboard the aircraft can be explained by the fact, that the measurements directly at the surface are affected by deposition and ozone loss processes within the stable nocturnal boundary layer, which was not accessible for the aircraft. As shown in Fig. 5 around 09:00 UTC at the end of the first flight the stable layer of the lowest level disappears and is then mixed with the newly
growing boundary layer. At this time, between the two flights, there occurs an increase of O₃ of 6 ppbv per hour up to a maximum mixing ratio of 43 ppbv at noon. The increase of ozone can be attributed to photochemical ozone production and downward mixing from the free troposphere. McKendry and Lundgren (2000) showed that an increase of the ozone mixing ratio larger than 5 ppbv h⁻¹ has been related to downward transport from the FT and not to photochemical production alone. As mentioned earlier, clouds formed initially between 09:00–10:00 UTC. For that reason the global radiation decreased and hence the photochemical ozone production efficiency has been reduced (Fig. 5).

Figure 14 shows the CO₂ mixing ratio onboard the aircraft and near the surface. The airborne CO₂ measurements show low mixing ratios of 381 ppmv (8:00 UTC) compared to the high CO₂ surface measurements. This large difference is due to accumulated CO₂ in the stable boundary layer during night. This is consistent with slightly enhanced ozone measured onboard the aircraft for this period (see Fig. 13). The aircraft was always flying at least 100 m above the surface where the CO₂ values are already much lower compared to the stable surface layer below. Note, that the CO₂ profile in Fig. 5 also indicates much higher CO₂ values at the lowest levels, which are in the same range as observed at the surface at Mount Kleiner Feldberg. Thus during night, CO₂ at the surface is dominated by plant respiration and anthropogenic CO₂ emissions. During day a potential effect from downward mixing at the surface is probably masked by some local uptake of CO₂ by vegetation.

The particle number concentration measurements show very low values of 1 to 2 particles ccm⁻¹ within the FT. Within the PBL we observe larger particle number concentration from the CPC (210 particles ccm⁻¹) than from the OPC (40 particles ccm⁻¹). Since the measurements site is covered by cumulus clouds (see fig. 4), it is obvious that particles larger than 100 nm were activated to form cloud droplets. These measurements could not be used to further support the mixing process. Furthermore, the backscatter signal from the ceilometer was mostly affected by the clouds (after 10:00 UTC) and therefore we observed only the cloud base height. The increase of O₃ near the surface is most likely the result of mixing of O₃ rich air masses from the FT. This downward transport process would match well with the time of occurrence of the mixing line.
5 Conclusions

Extensive surface based, balloon based and airborne measurements were collected as part of the measurement campaign PARADE on Mount Kleiner Feldberg (825 m) in summer 2011. We used this data to investigate the diurnal evolution of the boundary layer height over this complex terrain as well as the potential impact of local and long-range transport of trace gases. Our focus was on two research flights on 6 September 2011 where vertical tracer properties of CO₂ and O₃ indicated the occurrence of mixing across the entrainment zone. The planetary boundary layer height was determined using the Bulk-Richardson-number (\( R_i \)) from measurements taken from radio soundings. This approach can be used both under stable and unstable conditions using the same criteria. The boundary layer top grew from 860 m a.s.l. in the morning up to 2000 m a.s.l. in the afternoon. During this growth it merged with the residual layer (2100 m a.s.l. at 06:00 UTC). Negative \( R_i \) values indicated that the PBL is convectively driven and turbulent during that day, and thus well mixed. Additionally, airborne measurements of various trace gases (CO₂ and O₃) and particle number concentrations were performed in cloud-free areas during that day.

The temporal evolution of the measured vertical trace gas profiles showed the evolution of the PBL as observed with the radio soundings. The merging of the stable boundary layer, residual layer and free troposphere with the convective driven boundary layer is involved with mixing of the trace gases from these different atmospheric layers, and lead to complex vertical profiles of the trace gases. The tracer–tracer correlations have been applied to identify this mixing event and the effect on the trace gas distribution change across the transport barrier at the boundary layer top. This mixing line is probably formed by vertical uplift of CO₂ rich and O₃ poor air masses from the morning stable boundary layer. The strong western wind transported cold air masses from the Atlantic Ocean and caused the development of cloud streets. The structure of the trace gas profiles resembles the vertical cloud structure. We found the mixed air masses in the cloud interstitial space. This is consistent with convective uplifting and mixing with FT air at the cloud top within the clouds and subsidence of outside the clouds. It is suggested that these conditions forced mixing between the FT and
the PBL. It cannot be excluded that the mixing event was affected by long-range transport. The analysis of the local meteorological situation and the trace gas observations, however provide strong evidence, that the local boundary layer growth lead to mixing at the PBL top. This is supported by the specific humidity, which shows on the mixing line a slightly larger variability than the variability of the FT background humidity. The additional formation of shallow convective clouds at the PBL top (and roll convection) probably led to an efficient exchange between the boundary layer and the free troposphere.

We conclude, that rapid boundary layer growth and its interaction with convective clouds at the PBL top can efficiently mix both air masses and lead to variability that can be characterized by mixing lines. Moreover, the analysis of mixing lines provides a powerful tool to identify the occurrence and the effect of processes, which lead to a constituent exchange and mixing across the inversion layer.

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Table 1. Values to estimate the mass flux from the boundary layer growth budget, the entrainment and surface flux from aircraft, surface and radio sounding over Mount Kleiner Feldberg at 11:45 UTC.
Figure 1. Topographic map (colorbar in m) of the Rhein-Main Area in southwest Germany near Frankfurt am Main including the flight pattern on 6 September 2011 (black line). Vertical profiles are indicated by the flight track spirals. All descending profiles were flown near the summit of Mount Kleiner Feldberg, while ascending profiles have been taken up- and downwind.
Figure 2. Daily averaged CO$_2$ observations and the associated standard deviation during the measurement campaign parade on the Mount Kleiner Feldberg. The shaded areas indicate periods of high pressure influence (white) and low pressure influence (gray).
Figure 3. Backward trajectories, starting from the Mount Kleiner Feldberg at 200 m altitude over the surface for each day at 12:00 UTC up to 24 h based by wind fields from COSMO-7. The colors represent the daily averages of CO$_2$ measurements from Mount Kleiner Feldberg (Fig. 2).
Figure 4. Surface pressure chart provided from the German Weather Service (DWD) over Europe at the 6 September 2011, 12:00 UTC. The red box indicates the measurement area.
Figure 5. Diurnal variation of the residual layer height (RLH) and planetary boundary layer height (PBLH) derived from radiosondes (black) and ceilometer (red, hourly averaged). The distinction between the turbulent boundary layer and the stable FT is shown by the bulk Richardson number $Ri_B$ (gray scale). RLH is derived from the radiosonde temperature profiles (black dashed) and from one hour averaged aerosol layer heights (AHL) calculated from the ceilometer backscatter profiles (red dashed). Blue circles indicate the occurrence of clouds during that day. The global radiation (GR, blue line) is provided from HLUG in W m$^{-2}$. Vertical profiles of CO$_2$ are shown for different aircraft flights over Mount Kleiner Feldberg during this day in Rainbow colors, for details see Figs. 7 and 8.
Figure 6. Profiles of a) virtual potential temperature, b) wind speed, and c) wind direction derived from radiosondes during the day (black: 6 UTC, red: 10 UTC, blue: 11 UTC, orange: 14 UTC) on 6 September 2011.
Figure 7. Vertical profiles of a) CO$_2$ and of b) O$_3$ in the morning (07:27 UTC) on 6 September 2011. The color shows the specific humidity, and the horizontal lines show the boundary layer height (black dashed line) and the residual layer height (dashed gray line).
Figure 8. Same as Fig. 7 at noon. The measurement points between 1640 m and 2200 m a.s.l. are influenced by mixing from the PBL and the FT.
Figure 9. Correlations between CO₂ and O₃ during the observations at 07:27 UTC color-coded according to specific humidity measurements. The aircraft descended from the FT into the PBL, indicated by a clear separation from the chemical signature.
**Figure 10.** Same as Fig. 9 at noon. The measurement points between 1640 and 2200 m a.s.l. (black boxes) are related due to the mixed air masses from the stable Layer and the FT.
Figure 11. Temporal evolution of air mass origins between 11:00 and 13:00 UTC on 6 September 2011. The trajectory positions are colour coded by mixing ratios of ozone as measured along the flight leg. The data are binned by 0.5° longitude and 250 m altitude (details see text). The black line marks the maximum PBLH during that time.
Figure 12. Cloud formation below the flight track at about 11:32 UTC. Just before measuring a profile by descending above Mount Kleiner Feldberg. The arrows show the mean wind direction.
Figure 13. Surface ozone measurements of different instruments at different inlets (HLUG (blue, 30 min-mean) and UV absorption (red, 10 min-mean)) on Mount Kleiner Feldberg (825 m a.s.l.), including the aircraft observations (mean and standard deviation between 1000 to 1600 m a.s.l.) on 6 September 2011.
Figure 14. CO$_2$ surface measurements (on Mount Kleiner Feldberg (825 m a.s.l.), see Sect. 2.2), including the aircraft observations (mean and standard deviation between 1000 to 1600 m a.s.l.) on 6 September 2011. The high CO$_2$ fluctuations at 08:00 UTC is demonstrated by the mean (red star), standard deviation (black error bars), median (red bar) and maximum/minimum values (blue bar).