Replies and revised text in response to two reviewer comments are identical to those published online in ACPD, with the exception of one update (highlighted in red).

Reply to comments of Referee No.1
We thank Referee No.1 for insightful comments which helped to further improve the manuscript. Referee comments (in italics) are addressed below. Revised text, where necessary, is shown in blue, and is included in the final manuscript we will submit to ACP.

General: I enjoyed the paper and appreciate the challenges in linking NO\textsubscript{x} chemistry, the diurnal cycle, and boundary layer effects. The discussion of limitations in the evaluation of the effects of interferences was quite complete. It is clear that much has been accomplished but more remains to be done. A peeve of mine that the various papers in this collection do not appear to have been coordinated in the time Periods highlighted (particularly in the figures): For example, Gallee et al provides details on model versus observed boundary layer depth for 26-28 December but the really interesting chemistry is earlier in Periods II and III in Frey. Similarly earlier published work (Argentini et al.) looked carefully at 10 January boundary layer behaviour but Frey et al. did detailed profile measurements a day earlier. It would useful in collections such as these to identify specific Periods of common interest prior to extensive analysis.

Reply: Thank you. The comparison of modelled boundary layer depth with observations based on sodar measurements presented in Gallée et al. (2015) was possible on a few days only during the OPALE campaign (see also reply to a similar reviewer comment to Gallée et al. (2015)). NO\textsubscript{x} concentrations were measured from 23 November 2011 to 12 January 2012 (this work), HONO from 4 December 2011 to 11 January 2012 (Legrand et al., 2014), HO\textsubscript{2} radicals from 19 December 2011 to 9 January 2012 (Kukui et al., 2014), and CH\textsubscript{2}O from 14 December 2011 to 11 January 2012 (Preunkert et al., 2014). However, sodar measurements were only available on 12, 13, 18, 21 26, 27, 28 December 2011 and on 3, 4 and 10 January 2012. The best period to compare the MAR model with sodar measurements was 26 to 28 December 2011, because it was the longest Period and also included all chemical trace gases targeted during OPALE.

Regarding NO\textsubscript{x} observations, additional modelled and observed BL heights were available only during Period III. (9-22 December 2011), which we’ll note in the manuscript. Unfortunately, the NO\textsubscript{x} vertical profiles measured on 9 January 2012 can only be compared to modelled BL heights due to the lack of sodar data.

Revised text 31298, Lines 8-13: The strong increase of NO\textsubscript{x} around 11 December 2011 falls into a Period when F\textsubscript{NO\textsubscript{x}} almost tripled, while wind speeds slightly decreased and shallow boundary layer heights prevailed with daily h\textsubscript{z} maxima below 100-200 m (Fig.1, Table 2). On 12 December and 13 December the modelled diurnal ranges of h\textsubscript{z} were 3.4-224 m and 3.6-251 m, respectively, while sodar observations yielded 10-150 m and 5-125 m, respectively (Gallée et al., 2015). After 13 December 2011 F\textsubscript{NO\textsubscript{x}} remained at high values, thus, the decrease of NO\textsubscript{x} mixing ratios appears to be primarily caused by stronger upward mixing into a larger volume, i.e. wind speeds increased and daily h\textsubscript{z} maxima grew, exceeding 600 m on 18 December (Fig. 1).

Specific:
Abstract: Should be more specific about the difference between the South Pole and Concordia: It is not just the diurnal cycle but the sudden collapse of the boundary layer in the evening that is unique to Concordia (when the surface flux of NO\textsubscript{x} is suddenly confined to a shallow layer).

Revised text 31283, Lines 5-7: Profiles of NO\textsubscript{x} mixing ratios of the lower 100 m of the atmosphere confirm that, in contrast to South Pole, air chemistry at Dome C is strongly influenced by large diurnal cycles in solar irradiance and a sudden collapse of the atmospheric boundary layer in the early evening.

31284, Lines 15-19: You list four factors leading to high NO\textsubscript{x} at the South Pole (Davis et al. 2008) in the introduction. Your conclusion should come back and summarize which of these are relevant to Concordia. In particular, the statement low temperatures leading to low primary production rates of
Hox radicals should be addressed insofar as Davis et al argue that this is what contributes to the non-linear increase in the lifetime of NO₂ and high accumulation levels at the South Pole. Is there any relevance to the chemistry at Concordia?

Reply: Note reply and revised text have been updated compared to the version posted online in ACPD. The flux values observed at Dome C are typically large enough to explain the average increase in mixing ratios in the early evening (1700-1800 LT). For example in 2009-10 estimated net production rates of NO₂ at night are on the order of 100 pptv/h and are consistent with the average increase in NO₂ from 110 to 300 pptv over about 2 h from 1700 to 1900 LT (Frey et al., 2013). Similarly in 2011-12 observed increases in NO₂ mixing ratios are consistent with NO₂ production rates estimated with a simple box model. It is therefore in general not necessary to invoke non-linear HOₓ-NO₂ chemistry and associated increase in NO₂ lifetime as suggested in the case of South Pole (Davis et al., 2008, and references therein). We updated the text accordingly.

Revised text 31299, after Line 2: Previously, non-linear HOₓ-NO₂ chemistry and the associated increase in NO₂ lifetime were suggested to be an additional factor needed to explain large increases in NO mixing ratios observed at South Pole (Davis et al., 2008, and references therein). In order to assess the relevance of this factor at Dome C we apply a simple box model to estimate net NO₂ production rates as done previously (Frey et al., 2013). It is assumed that that mixing is uniform and instantaneous, that the snow emission flux F_{NO₂} is the main NO₂ source and the reaction with the OH radical is the dominant NO₂ sink.

\[
\frac{d[NO₂]}{dt} \sim \frac{F_{NO₂}}{h_z} - k[NO₂][OH]
\]  

(1)

where \( k \) is the respective reaction rate coefficient. In 2009-10 no OH observations were available at Dome C and average values from South Pole were used instead. In 2009-10 estimated net production rates of NO₂ at night were on the order of 100 pptv/h and therefore explained the observed average increase in NO₂ from 110 to 300 pptv over about 2 h from 1700 to 1900 LT (Frey et al., 2013). In 2011-12 the same analysis is repeated using OH measurements available for most of Period IV. (Kukui et al., 2014) as well as \( h_z \) calculated with the MAR model (Gallée et al., 2015). Resulting night time values of net NO₂ production rates are with about 40 pptv/hr smaller than in 2009-10 but again to a first order consistent with a smaller observed increase in NO₂ mixing ratios in the evening hours; i.e. during Period IV. median NO₂ increased between 1630 and 1930 LT from 114 to 242 pptv (Fig.6a,f). The above model is oversimplified as the very likely presence of HO₂NO₂ will modulate the diurnal variability of NO₂ sinks and sources with an impact on NO₂ lifetime as suggested by Davis et al. (2008). However without any information on the diurnal cycle of HO₂NO₂ at Dome C further modelling is not warranted.

Revised text 31301, after Line 16 (Conclusions): Large mixing ratios of NO₂ at Dome C arise from a combination of several factors: continuous sunlight, large NO₂ emissions from surface snow and shallow mixing depths after the evening collapse of the convective boundary layer. Unlike at South Pole it is not necessary to invoke non-linear HOₓ-NO₂ chemistry to explain increases in NO₂ mixing ratios. However, uncertainties remain regarding atmospheric levels of HO₂NO₂ and its impact on NO₂ life time being a temporary NO₂ reservoir.


Reply: Agreed, we replaced the Cohen et al. (2007) reference with the one above.

31290, Lines 11-12: Can you argue that the NO₂ flux is constant with time through the collapse of the boundary layer. Eliminating 22% of the data when the boundary layer depth is < 10 m may be problematic if this 22% occurs during the evening transition when NOX levels get large.

Reply: The point raised is unclear. We do not argue that flux is constant with time. To the contrary, we argue that the application of MOST requires that at a given time flux is constant between the two measurement heights (condition (a)). Constant flux can be assumed if the chemical lifetime (\( \tau_{chem} \)) of NO_x is much longer than the turbulent transport time scale (\( \tau_{trans} \)). This condition is met during
However, modelled mixing heights indicate that the upper inlet is frequently above the surface layer (condition (c)) during and after the collapse of the convective BL, as pointed out by the reviewer. Thus, the removal of flux estimates (22% of the total available) affects mostly the evening and night, when the BL is shallow. Hence, fluxes during night time are less well constrained, but nevertheless support a significant diurnal cycle (see Figure 6b,g and Figure 9 in Frey et al. (2013)). We clarify this point in the revised text.

Revised text 31290, Lines 8-17: Condition (a) is met in the surface layer if the chemical lifetime \( \tau_{chem} \) of \( \text{NO}_x \) is much longer than the turbulent transport time scale \( \tau_{trans} \). Based on observed \( \text{OH} \) and \( \text{HO}_2 \) \( \tau_{chem} \) for \( \text{NO}_x \) is estimated to be 3 h at 1200 LT and 7 h at 0000 LT during OPALE (Legrand et al., 2014). Estimating \( \tau_{trans} \) following the approach described previously (Eq. 6 and 7 in Frey et al., 2013) yields 0.6, 1.7 and 2.5 min during the day (0900-1700 LT), the typical time of BL collapse (1700-1900 LT) and during the night (1900-0900 LT), respectively. Thus, \( \tau_{chem} \) exceeds \( \tau_{trans} \) by at least a factor 100, confirming that vertical mixing always dominates over the gas phase photochemical sink and flux can be assumed constant between the two inlets. Condition (b) is met as discussed in Frey et al. (2013). For (c) the upper inlet height of 1 m is compared to estimates of mixing height \( h_z \) from the MAR model (Gallée et al., 2015). Calculated flux values of \( \text{NO}_x \) were removed when \( h_z < 10 \) m resulting in the removal of 22% (773 values) of all available 10 min flux averages. Flux estimates are removed specifically during the evening and night, when the BL is shallow. Hence, fluxes during night time are less well constrained, but nevertheless support a significant diurnal cycle (see Figure 6b,g and Figure 9 in Frey et al. (2013)).

Revised text 31293, Lines 10-19: This description of changes in NOx levels could use a bit more work. The intraseasonal trend should be characterized as intraseasonal variability.

Reply: As suggested we now characterize in the text the ‘trend’ as intra-seasonal variability.

Revised text 31293, Lines 10-19 (continued): Also, there is a gap in wind data Dec 3-7. I looked at the AWS data (ftp://amrc.ssec.wisc.edu/pub/aws/10min/dph/2011/089891211.r) for this Period and it looks like the wind speed was greater in Period II compared to Period III. The AWS anemometer data shows frequent stalling in Period III. However, a simple average yields Period II: 2.4 m/s whereas Period III: 1.3 m/s. This suggests a closer look at the depth of mixing between the two Periods.

Reply: Indeed, the statistics for Periods I-IV listed already in Table 2 indicate that wind speeds at 3.3 m were greater during Period II (median 3.6 m/s) than during Period III (median 2.5 m/s). We closed the gap in wind data Dec 3-7 using measurements at 2.0m height from the AWS managed by Univ. of Wisconsin, which reduces the median wind speeds during Period II to 3.0 m/s. Recalculating the average wind speeds from the Wisconsin AWS yields 2.98 and 2.34 m/s for Periods II and III, respectively, higher values but with a smaller difference than suggested by the reviewer. Thus, wind speeds in the two Periods are different, but not by much. The primary change driving the increase of \( \text{NO}_x \) mixing ratios is the increase in flux combined with lower wind speeds and a relatively shallow BL (see above for revised text 31298, Lines 8-13).

Revised Table 2 and Fig. 1: Median wind speed for Period II updated to 3.0 m/s and Fig. 1 includes now wind speed observations during 3-7 December.

Revised text 31293, Lines 10-19 (continued): With respect to the correlation between wind speed and \( \text{NO}_x \) levels, another factor to look at is the response of \( \text{NO}_x \) concentrations to the sudden collapse of the boundary layer in the evening. Ideally, one should compare average winds during just the Period of collapse and higher \( \text{NO}_x \); the correlation might come out differently.

Reply: Regarding wind speeds, we looked at the correlation with \( \text{NO}_x \) mixing ratios during the time of collapse of the daytime convective BL, i.e. 1700-1900 LT, and find a slightly stronger negative correlation \( (R=-0.45, p<0.001) \) than when including all data \( (R=-0.37, p<0.001) \).

Revised text 31293, Lines 6-9: As seen previously at Dome C and other locations, \( \text{NO}_x \) mixing ratios were weakly but significantly anti-correlated with wind speed \( (R=-0.37, p<0.001) \), especially when only the time Period of the daily collapse of the convective boundary layer, i.e. 1700-1900 LT,
was considered ($R=-0.45$, $p<0.001$).

31293, Lines 10-19 (continued): There were also significant changes in the behaviour of the wind direction: Early Period III shows a 180 degree rotation of the wind whereas Period II shows a most consistent wind direction centered from the SE. In Period III, when the wind was rotating from SW to SE to N, could there have been contamination from the station? (Frey et al. 2013, Figure 1; also see Gallee this issue, their Figure 3).

**Reply:** Wind directions changed during Period III and rotated through northerly directions, potentially carrying contamination from the generator at Concordia station. However, our data filtering efficiently removes any pollution spikes, which typically exceed 10 ppbv of NO$_x$ (see also Frey et al., 2013). The regular appearing diurnal NO$_x$ maxima are clearly linked to the drastic BL decreases in the early evening, including Period III (Fig. 1). To illustrate this point better we updated Fig. 1 and include also a time series of wind direction (Fig. 1). In the method section we provide more detail on how we removed NO$_x$ data affected by local air pollution using a filtering method described previously (Frey et al., 2013).

**Revised text in section 2.1:** The mean wind direction during the measurement period was from S (176°) with an average speed of 4.0 m s$^{-1}$ (Fig. 1b). During 2.5% of the time winds came from the direction of Concordia station, i.e. the 355-15° sector (see Fig.1 in Frey et al., 2013), carrying potentially polluted air from the station power generator to the measurement site. For example, during Period III winds rotated 4 times through northerly directions (Fig. 1b). Pollution spikes in the raw 1-s data typically exceeded 10 ppbv of NO$_x$ and were effectively removed before computing the 1-min averages by applying a moving 1-min standard deviation filter. Observations were rejected when 1-$\sigma$ of NO and NO$_2$ mixing ratios within a 1-min window exceeded 24 and 90 ppt, respectively.

**Figure 2:** The discussion of this figure might want to include a reference to Argentini et al. 2013 (Annals of Geophysics 56, 5, 2013; 10.4401/ag-6347) which shows the negative heat flux at sunset as well as the decrease in downward longwave radiation for 9 January 2012 (rapid cooling of the surface resulting in a strong shallow surface inversion. That paper also shows fairly graphically, using sodar data, the evolution of the boundary layer on 10 January 2012 it would be nice to have a similar figure for the 9th together with Gallees simulation (note that the Gallee paper in this special issue compares modeled versus sodar observed BLD for 26-28 December 2011) It would be nice if these comparisons could be coordinated and cross referenced between the papers (e.g. the high NO$_x$ Period 12-16 December). Also, Gallees Figure 6 shows a later falloff in BLD in his model than does the sodar does the same result hold for the 9th.

**Reply:** In the discussion of Fig.2 we now include a reference to the observations reported in Argentini et al. (2014). As mentioned above in the reply to the first comment, a comparison with modelled and observed BL heights is possible only during a few days of the OPALE campaign. In particular, no sodar data are available for 9 January 2012 and therefore we compare measured NO$_x$ vertical profiles only to modelled BL heights.

**Added text in section 3.2; 31294, after Line 12:** At Dome C rapid cooling of the surface in the evening results in a strong shallow surface inversion (e.g. Frey et al., 2013), and is illustrated by a decrease in downward long-wave radiation and a negative heat flux, as observed in the evening of 9 January 2012 (Fig.4 in Argentini et al., 2014).

**Section 3.5.2, : This section should probably reference/compare other NOX flux measurements. See Davis et al 2008 and references therein (Onckley et al, Wolff et al., Wang et al. and Neff et al) that discuss the magnitudes, estimates, and boundary depth effects relevant to the NOX flux (esp. Wang et al).**

**Reply:** As suggested we expand the discussion referring to previous estimates of NO$_x$ emission flux ($F_{NO_x}$) from polar snow based on observations and models.

**Added text in section 3.5.2; 31294, before Line 1:** The NO$_x$ flux observed above polar snow is on the order of $10^{12}$ to $10^{13}$ molecule m$^{-2}$ s$^{-1}$ and contributes significantly to the NO$_x$ budget in the polar boundary layer. At the lower end of the range are $F_{NO_x}$ observations at Summit,
Greenland (Honrath et al., 2002) and at Neumayer in coastal Antarctica (Jones et al., 2001) with 2.5 x 10^{12} molecule m^{-2} s^{-1}, whereas on the Antarctic Plateau F_{NO_x} values are up to ten times larger (Oncley et al., 2004; Frey et al., 2013, and this study). The average F_{NO_x} at South Pole during 26-30 November 2000 was 3.9 x 10^{12} molecule m^{-2} s^{-1} (Oncley et al., 2004), whereas at Dome C observed fluxes are 2-6 times larger, with seasonal averages of 8-25 x 10^{12} molecule m^{-2} s^{-1} (Frey et al., 2013, this work). Due to the uncertainties in the processes leading to NO_x production it had been difficult to explain inter-site differences, e.g. by simply scaling F_{NO_x} with UV irradiance and nitrate in the surface snow pack (Davis et al., 2004). Some of the variability in flux values may be due to differences in experimental set up or in the employed flux estimation method (e.g. Davis et al., 2004; Frey et al., 2013). For example, the F_{NO_x} estimates for South Pole are based on measured NO gradients only, inferring NO_x from photochemical equilibrium and using the Bowen ratio method (Oncley et al., 2004), whereas the F_{NO_x} estimates for Dome C are based on observations of both atmospheric nitrogen oxides (NO and NO_2) and the flux-gradient method (Frey et al., 2013).

Model predictions of F_{NO_x} show in general a low bias on the Antarctic Plateau when compared to observations. A first 3-D model study for Antarctica included NO_x snow emissions parameterised as a function of temperature and wind speed to match the observed F_{NO_x} at South Pole (Wang et al., 2007). However, the model under-predicts NO mixing ratios observed above the wider Antarctic Plateau highlighting that the model lacks detail regarding the processes driving the emission flux (Wang et al., 2007). The first model study to calculate F_{NO_x} based on NO_3 photolysis in snow, as described in this work, reports 1.5 x 10^{12} molecule m^{-2} s^{-1} for South Pole in summer (Wolff et al., 2002), about a factor 4 smaller than the observations by Oncley et al. (2004) and up to 16 times smaller than what is needed to explain rapid increases in NO_x mixing ratios over a few hours (Davis et al., 2008, and references therein). Recent model improvements reduced the mismatch with the South Pole flux observations and included the use of updated absorption cross sections and quantum yield of the nitrate ion, as well as e-folding depths measured in surface snow on the Antarctic Plateau, and resulted in a factor 3 increase of flux calculated for South Pole (France et al., 2011). In light of major remaining uncertainties, which include the spatial variability of nitrate in snow and the quantum yield of nitrate photolysis (Frey et al., 2013), we discuss below the variability of F_{NO_x} observed at Dome C.

References


Figure 1: Meteorology and NO$_x$ observations at Dome C in summer 2011–2012 (highlighted Periods I.–IV. as referred to in text and Table 2): (a) air temperature ($T$) at 1.6m and modeled mixing height ($h_z$) (Gallée et al., 2015), (b) wind speed (wspd) and direction (wdir) at 3.3 m (c), 1 min averages of NO$_x$ mixing ratios at 1 m (red line is 1 day running mean) and (d) 10 min averages of observational estimates of NO$_x$ flux ($F_{NO_x}$) between 0.01 and 1 m (red line is 14 day running mean).
Reply to comments of Referee No.2

We thank Referee No.2 for insightful comments which helped to further improve the manuscript. Referee comments (in italics) are addressed below. Revised text, keyed to the ACPD online version, is shown in blue, and is included in the final manuscript we will submit to ACP.

This manuscript focuses on observations of NO and NO\textsubscript{2} from 10 cm into the firn up to 100 m above the snow surface, and measurement of the flux of NO\textsubscript{3} (primarily) out of the snow, made at Dome C during the OPALE campaign. There were abundant supporting measurements available, allowing the authors to put important constraints on the factors controlling variations in the mixing ratios of the nitrogen oxides and the snow to air flux of NO\textsubscript{2} over a range of time scales. At Dome C it is clear that the interplay between the strength of the snow source of NO\textsubscript{3} and vertical mixing exerts primary control over the mixing ratio of NO\textsubscript{2}, both over the course of a day and a season. To first order, the strength of the snow source can also be explained as a combination of the abundance of NO\textsubscript{3} in the snow available for photolysis and the actinic flux in the 300-340 nm range. However, it is evident that variations in the strength of the snow source of NO\textsubscript{3} depend on additional factors that are not fully understood. Authors suggest that the fraction of NO\textsubscript{3} readily photolyzed can change, both within a single season and between years, and suggest more field, lab and model studies are needed to understand what makes some NO\textsubscript{3} - photo-labile while other NO\textsubscript{3} is not. The latter fraction is tentatively labeled photo-stable, I suggest that we do not need a new word and that photo-stable or just stable should be adequate terminology.

Reply: We replaced photo-stable with photo-stable in the text.

In general, the results are clearly presented, and the arguments supporting conclusions are well laid out. I will suggest a few places where I feel that clarity could be improved in the list of detailed minor comments below. However, I feel that more detail is needed in the description of methods. Most importantly, the authors need to explain how the concentration gradients were measured. Seems that the 2-channel CLD allowed NO and NO\textsubscript{2} (and something like NO\textsubscript{3} by difference) to be determined simultaneously, but only one inlet could be sampled at a time. So, what was the cycle between 0.1, 1.0 and 4.0 m sampling heights? How much of each 10 minute interval was spent at each height? Was each height measured several times in the 10 minutes, or was it 0-3.333 minutes on one inlet, 3.333-6.666 on next and then 6.666-10 on the third? One presumes that there had to be some down time for zeroing and calibration, perhaps quite frequently, so did these essential intervals of housekeeping result in gaps during all or most 10-minute gradient measurements, or were they grouped into a longer period of no data once or several times each day?

Reply: The measurement method as well as the duty cycle of the CLD during the gradient measurements were described previously in Frey et al. (2013). However, we repeat some of the requested details in the methods section.

Revised text 31287, after Line 17: The three sample inlets were connected inside the lab shelter to a valve box, which automatically switched the CLD between sampling heights on a 90 s duty cycle. As described below, the 10-minute average concentration difference \( \Delta \text{NO}_x \) between the 0.01 and 1.0 m inlets is used to estimate flux. Therefore, 10-minute mean \( \Delta \text{NO}_x \) values are calculated on average from two sets of two subsequent 90 s intervals, separated by a 90 s interval during which the 4.0 m inlet was measured. Baseline count rates were determined by adding excess ozone to sample air in a pre-chamber so that all electronically excited NO\textsubscript{3} has returned to ground state when reaching the reaction chamber. The baseline was measured for 60 s every 13.5 min alternating between all three inlets. The NO sensitivity of the CLDs was determined every 14 h by standard addition to the sample air mixture of a 1 ppm NO/NO\textsubscript{2} mixture (UK National Physics Laboratory traceable BOC certified), which is further diluted to 4 ppbv of NO. During standard runs also the conversion efficiency (CE) of the photolytic converter was determined by addition of a known mole fraction of NO\textsubscript{2}. This was achieved by gas phase titration of the NO/NO\textsubscript{2} mixture to NO\textsubscript{2} by O\textsubscript{3} generated from a pen-ray lamp, and monitoring the un-titrated NO mole fraction. The instrument artefact originating from NO\textsubscript{2} producing surface reactions in inlets and reaction cells was determined by overflowing the instrument inlet with scrubbed ambient air supplied by a pure air generator (Eco-Physics PAG003). The artefact
was measured every 14 h, offset by 7 h to the calibration runs.

I also feel that more detail needs to be provided regarding the measurements of snow nitrate in the field lab. It is stated that samples were collected every few days, but I am curious if they were analyzed right after collection (that day or the next), or allowed to pile up and then run in larger batches several times through the season, or maybe even all in one bunch near the end (this last option might be the best answer, but seems unlikely). In general, this would not seem something to worry about except for the fact that Berhanu et al. also have a manuscript on OPALE in review at ACPD right now, and indicate some uncertainty about nitrate measurements at Dome C during the 2011-12 season. Specifically, they measured what was supposed to be the same snow in 2 different artificial snow pits 12 different times through the season and found a range from 1200 1700 ppb (around a stated true value of 1450 ppb). This variability was not seen in samples run in a single batch, rather was expressed as large shifts between samples run on different days. Were the samples in present study and those reported by Berhanu et al. all run by same technician on the same instrument (commingled in batches)?

Reply: During OPALE the skin layer of surface snow, i.e. the top few mm, was sampled every 3 days. Samples were stored together with the additional snow samples discussed in Berhanu et al. (2014) and then analysed for nitrate in batches by the same operator. There has been a systematic shift in the nitrate standard response in between individual batch runs due to a calibration issue, which may affect the time series of nitrate in surface snow (Berhanu et al., 2014). However we believe the trend during Period II and III (Fig.7c) is real for two reasons: a) all samples were analysed in random order, across several batches, but the temporal trend observed in surface snow concentrations is very similar in both the skin layer (top few mm) and in the top layer of adjacent snow pits (top 2cm) (Fig.4 in Berhanu et al., 2014). And b) nitrate maxima in Antarctic surface snow during summer are a robust feature observed at Dome C over the existing 2009-present period of year-round sampling (Fig. 7b shows 2011-12 and 2009-10), as well as in coastal Antarctica (e.g. Mulvaney et al., 1998). Thus the snow nitrate changes over a week with a typical amplitude of 800-1000 ppbw are repeatable and well above the spatial variability of 20-25% found at Dome C (France et al., 2011; Frey et al., 2013). See revised text further below.

How much would modeled NO_x fluxes change if snow nitrate was adjusted up or down by nearly 20%?

Reply: The NO_x emission flux scales with nitrate concentrations in snow as illustrated by equation 3, and therefore an uncertainty of 20-25% in nitrate concentration, which is in fact typical for the spatial variability seen at Dome C, will translate to a similar variability in F_NO_x (see Fig.7c and discussion in Frey et al., 2013).

Are the higher values of skin nitrate in periods II and III in 2011-12 real, or related to standard drift (sensu Berhanu et al.)?

Reply: As discussed above the higher values during Period II and III represent in our opinion a real temporal trend.

Are the surface snow data in Fig. 7 of this manuscript the same as those in Fig 4 of Berhanu (seem similar, but maybe not identical)? At a minimum, authors need to make a good faith estimate of the precision and accuracy of their own snow nitrate concentrations given the apparent problems in the field during OPALE.

Reply: The nitrate concentrations in the skin layer of surface snow during 2011-12 used in this work (Fig.7b) are the same as presented in Fig.4 of Berhanu et al. (2014). The updated precision is 5% based on replicate standard measurements (see also reply to reviewer No.1 in the online discussion of Berhanu et al. (2014)). The overall accuracy including systematic errors in calibration and collection of just the top few mm of snow is of the order of 20%, comparable to the spatial variability of nitrate in surface snow at Dome C.
Samples were stored together with the additional snow samples discussed in Berhanu et al. (2014) and then analysed for nitrate in batches by the same operator. The precision is 5% based on replicate standard measurements. Due to a systematic shift in the nitrate standard response in between individual batch runs due to a calibration issue (Berhanu et al., 2014) the accuracy is larger than usual. The overall accuracy including systematic errors in calibration and collection of just the top few mm of snow is of the order of 20%, and is therefore comparable to the spatial variability of nitrate in surface snow at Dome C (France et al., 2011). Note that the temporal trend of nitrate concentrations observed in surface snow discussed below is significantly larger, i.e. >50%.

Specific, mostly minor, comments keyed to line numbers in 9 Nov 14 Latex file.

11 interference by pernitric Reply: Done.

21-22 last sentence of abstract seems to clash with the one just before, and kind of comes out of the blue. Paper does develop this idea, but maybe it should just be in conclusions (or it needs to be brought into abstract less abruptly).

Revised text 31283, Lines 23-24: A remaining source of uncertainty and subject of future research is the quantum yield of nitrate photolysis in natural snow, which may change over time as the snow ages.

61 do not need mixing ratios and levels both in this sentence Reply: Corrected.

72-74 agreed that the quantum yield uncertainty is important, but probably not the dominant problem models are facing. Seems premature to highlight this again here. Reply: We removed this sentence.

104-105 only air from the bottom and sides could enter through small holes in the tube (might help to specify that the holes were x cm or mm above the bottom of the probe)

Revised text 31287, Line 6: ... from the bottom and sides could enter, using small horizontal holes at 0-10 cm above the bottom of the tube.

116 CLD employed also converts nitrous Reply: Done.

126-158 to me, it would flow more smoothly to switch the order of these 2 paragraphs, dealing with possible HNO₄ artifact on the NO₂ measurement right after HONO/NO. Then talk about how both possible problems and any vertical gradients might impact the NOₓ gradients.

Reply: Agreed, and we changed the order.

187 given how important mixing height is for much of the discussion, I would like to see some indication that MAR has been validated. Ideally at DOME C (from tower, tether sonde, maybe aircraft profiles) but at least from somewhere on the Antarctic plateau.

Revised text 31290, Line 10: The MAR model has been validated previously over the Antarctic Plateau, focusing on Dome C, during winter (Gallée and Gorodetskaya, 2010) and now also during summer (Gallée et al., 2014).

Equation 3, might mention that this model probably estimates an upper limit for NO₂ flux (if the quantum yield and actinic flux are correct) since it assumes all NO₂ formed escapes the firn before any of it can photolyze, or convert to HNO₃, HONO, HNO₄.

Revised text 31292, after Line 24: For the discussion below it should be borne in mind that the calculated F_NOX is a potential emission flux assuming that NO₂ is vented immediately after release from the snow grain to the air above the snow pack without undergoing any secondary reactions.

265 intra-seasonal trend odd terminology, since it seems you are talking about the week long period
with enhanced mixing ratios, not really a trend through the 2 months

Reply: We replaced intra-seasonal trend by "intra-seasonal variability" throughout the text.

268-269 to late December average (not Nov)

Reply: We clarify this.

Revised text 31293, Lines 15-16: After that NOx mixing ratios gradually dropped over 10 days (Period III-IV) to median concentrations of ~120 pptv, slightly lower than observed in late November (Table 2).

270 2.5 times that Reply: Done.

273 median (range) of 1.6 (0.4-2.9) this is a little misleading. The range shown in Fig 1 D is -1 to 10 x 10^{13}. The smaller range in the text comes from Table 1 which compares season long medians for noon and midnight.

Reply: We believe to assess the range of flux values it is a more conservative measure to state median values at noon and midnight, which are less sensitive to extreme values and the occasional outlier present in relatively noisy flux estimates. We clarify this.

Revised text 31293, Lines 21:

... with a median of 1.6 x10^{13} molecule m^{-2} s^{-1}. Median values of $F_{NO_x}$ at midnight and at noon were 0.4 and 2.9 x10^{13} molecule m^{-2} s^{-1}, respectively (Table 1).

276 almost 5 times (or, about 4.7 times) Revised text 31293, Lines 24: almost 5 times

305-309 Any speculation about why the nitrate profile in the pit under the disk so much different than away from all the activity?

Reply: We have no definite answer to this question. The firn air probe was installed onto untouched snow, and only removed after the end of the atmospheric sampling period. Thus contamination is unlikely, but a local anomaly remains a possibility as pits 5 m next to the lab shelter showed a similar increase of concentration with depth.

Revised text 31295, Line 5: The firn air probe was installed onto untouched snow, and only removed after the end of the atmospheric sampling period. Thus contamination due to local activity appears unlikely, but a local anomaly remains a possibility as snow pits 5 m next to the lab shelter showed a similar increase of concentration with depth (data not shown).

316 the anticorrelation between NO$_2$ and O$_3$ is interesting, but the suggestion that it reflects enhanced nitrate in the snow is not supported. Profile in P3 does not get so deep and neither of the other pits shows a peak near 45 cm

Revised text 31295, Line 11-13: In particular, the drop of O$_3$ mixing ratios by >10pbv at 45 cm depth was not an outlier since collocated NO$_2$ mixing ratios were also significantly elevated compared to adjacent snow layers. However, no snow nitrate measurements were available to further investigate the origin of the NO$_2$ peak.

351-354 seems the details of the MAX DOAS data reduction should have been in Methods

Reply: We moved this part to the method section.

370-371 this statement begs for at least a back of the envelope attempt at quantification. You earlier estimated that HNO$_4$ might contribute 33-66 ppt artifact to NO$_2$, so what would happen if you reduced NO$_2$ by this much in the steady state calculation? Hard to believe this would account for much of the factor of 20 discrepancy.

Reply: This is a very useful comment. In reassessing the potential HNO$_4$ interference we discovered a computational error, which reduces the estimated magnitude of the HNO$_4$ interference by a factor two. NO$_2$ mixing ratios are corrected, assuming that additional NO$_2$ is measured in the CLD from HNO$_4$ thermal decomposition, equivalent to 25% of ambient HNO$_4$ on the order of 130 pptv. We find that average steady-state estimates of oxidant concentrations are still a factor 10 larger than those
observed (RO₂ and BrO). In the revised text we take this into account and put the role of HNO₄ as an interferent into perspective.

**Revised text 31288, Line 29:** HNO₄ present at these values could potentially produce 16-32 pptv of NO₃ in the photolytic converter, equivalent to 8-16% of the average NO₂ mixing ratio measured at 1 m.

**Revised text 31297, Line 9-16:** The same steady-state calculation as described in Frey et al. (2013) was repeated for austral summer 2011-12 and yields an average of 2.5 x10⁴ molecule cm⁻³ or 129 pptv of total radical concentrations \([OX] = [HO₂] + [RO₂] + 2[ XO] \). Observations based on median \([RO₂]\) of 9.9 10⁶ molecule cm⁻³ or 5 pptv (Kukui et al., 2014) and 3 pptv of BrO yield \([OX]\) of about 11 pptv. Hence, the estimated total radical concentration exceeds observations by a factor 12. To estimate the impact of a potential interference by HNO₄ we corrected the NO₂ mixing ratios, assuming that additional NO₂ is measured in the CLD from HNO₄ thermal decomposition, equivalent to 25% (100%) of ambient HNO₄ on the order of 130 pptv. We then find that the average steady-state estimate of oxidant concentrations is still a factor 10 (3) larger than those observed. Thus, a least a part of the inconsistency may be explained by the interference with HNO₄ (not measured).

375-385 this section is a little loose. Starts by saying that period II looks much like 2009-10 with peak 18-20:00 but the figure shows that in 2009-10 the peak lasted later into the evening. Indeed, in all of the intervals except II the evening peak lasts quite a bit past 20:00. Why would that be, since the mixing height is not getting much lower, and the snow source should be weakening.

**Reply:** Thanks for pointing this out. Indeed, NO₂ mixing ratios typically show maxima lasting into the night time hours in 2009-10 and in 2011-12 (except Period II.), whereas NO mixing ratios peak during 1800-2000 LT in 2009-10 (Fig.5 in Frey et al., 2013) and in 2011-12 (data not shown). Assuming no significant changes in BL height after the initial collapse of the convective BL night-time peaks of NO₂ are plausible if the weakening of snow emissions was offset by a corresponding decrease of the chemical sink of NO₂, i.e. the NO₂+OH reaction. This is consistent to a first order taking into account that observed OH concentrations and F-NO₂ vary in a similar way, by up to a factor 5 between local noon and midnight. We revise the text accordingly.

**Revised text 31297, Lines 18-28:** On diurnal time scales NO₂ mixing ratios at Dome C are controlled by the interplay between snowpack source strength and atmospheric physical properties, i.e. turbulent diffusion of heat \(K_h\) and mixing height \(h_m\) of the boundary layer. The median diurnal cycles of NO₂ mixing ratios in 2011-12 show with the exception of Period II (1-8 December) previously described behaviour (Frey et al., 2013), with a strong increase around 1800 LT to maximum values, which last into the night time hours (Fig.6a). Night-time peaks of NO₂ are plausible if the weakening of snow emissions was offset by a corresponding decrease of the chemical sink of NO₂, i.e. the NO₂+OH reaction, assuming no significant change in \(h_m\). This is consistent to a first order taking into account that observed OH concentrations and F-NO₂ vary in a similar way, by up to a factor 5 between local noon and midnight.

427-436 Another place text could/should be more precise. Assuming the snow nitrate concentrations are valid, the really high levels are only present at the end of II and beginning of III, not through both periods. Cant say much about NO₂ flux in II, but it clearly stays high through nearly all of III, despite an apparent steep drop in nitrate.

**Reply:** We agree and therefore refined the description of the observations in the text accordingly.

**Revised text 31299, Lines 21-24:** Instead changes in \(F_{NO₃}\) can be linked to the temporal variability present in the snow skin layer. During the end of Period II. and beginning of Period III. skin layer NO₃⁻ concentrations were up to two times larger than before and after (Fig.7b). \(F_{NO₃}\) is high during the end of Period II. and beginning of Period III., however drops off a week after the decrease of nitrate concentrations in surface snow.

451 corresponds to days of should this be to No of days?
Revised text 31300, Line 18: quantum yield ... decreased from 0.44 to 0.003 within what corresponds to a few days of UV exposure in Antarctica ...

452 in quantum yield is Reply: Corrected.

453 dont think stabile is a word and stable would probably work Reply: Corrected.

455 Neff and Davis also advocating for different flavors of nitrate in snow, shown on their poster at AICI CASSI, with references to earlier work. Reply: Correct, thus we added a reference to the earlier work by Davis et al. (2008).

480 is an O3 sink Corrected.

491-493 as noted earlier, should estimate how big a part HNO₄ might explain Reply: See reply above.

References


Atmospheric nitrogen oxides (NO and NO₂) at Dome C, East Antarctica, during the OPALE campaign

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Abstract

Mixing ratios of the atmospheric nitrogen oxides NO and NO$_2$ were measured as part of the OPALE (Oxidant Production in Antarctic Lands & Export) campaign at Dome C, East Antarctica (75.1° S, 123.3° E, 3233 m), during December 2011 to January 2012. Profiles of NO$_x$ mixing ratios of the lower 100 m of the atmosphere confirm that, in contrast to South Pole, air chemistry at Dome C is dominated by strongly influenced by large diurnal cycles in solar irradiance and atmospheric stability a sudden collapse of the atmospheric boundary layer in the early evening. Depth profiles of mixing ratios in firn air suggest that the upper snowpack at Dome C holds a significant reservoir of photolytically produced NO$_2$ and is a sink of gas phase ozone (O$_3$). First-time observations of BrO at Dome C suggest 2–3 pptv near the ground, with higher levels in the free troposphere. Assuming steady-state, observed mixing ratios of BrO and RO$_2$ radicals are too low to explain the large NO$_2$ : NO ratios found in ambient air. A previously not considered interference with possible interference by pernitric acid (HO$_2$NO$_2$) may explain part of this inconsistency. During 2011–2012 NO$_x$ mixing ratios and flux were larger than in 2009–2010 consistent with also larger surface O$_3$ mixing ratios resulting from increased net O$_3$ production. Large NO$_x$ mixing ratios arose at Dome C arise from a combination of changes in continuous sun light, shallow mixing height and significant NO$_x$ snow emission flux emissions by surface snow ($F_{NO_x}$). During 23 December 2011–12 January 2012 median $F_{NO_x}$ was twice that during the same period in 2009–2010 due to significantly larger atmospheric turbulence and a slightly stronger snowpack source. A tripling of $F_{NO_x}$ in December 2011 was largely due to changes in snow pack source strength caused primarily by changes in NO$_3^-$ concentrations in the snow skin layer, and only to a secondary order by decrease of total column O$_3$ and associated increase in NO$_3^-$ photolysis rates. Systematic changes in—A source of uncertainty is the quantum yield of photolysis over time may contribute to the observed $F_{NO_x}$ variability nitrate photolysis in natural snow, which may change over time as the snow ages.
1 Introduction

The nitrogen oxides NO and NO₂ (NOₓ = NO + NO₂) play a key role in the polar troposphere in determining its oxidation capacity, defined here as the sum of O₃, HOₓ radicals, and hydrogen peroxide (H₂O₂). The influence is achieved via photolysis of NO₂, the only source for in situ production of tropospheric O₃, through shifting HOₓ radical partitioning towards the hydroxyl radical (OH) via the reaction NO + HO₂ → NO₂ + OH, and finally through reactions with peroxyradicals NO + HO₂ (or RO₂) which compete with the formation of peroxides (H₂O₂ and ROOH).

Atmospheric mixing ratios of NOₓ in the atmospheric boundary layer of coastal Antarctica are small, with average NOₓ values in summer not exceeding 30 pptv (Bauguitte et al., 2012). The build up of large mixing ratios is prevented by gas-phase formation of halogen nitrates (e.g. BrNO₃, INO₃) followed by their heterogeneous loss (Bauguitte et al., 2012). Conversely, mixing ratios of NOₓ on the East Antarctic Plateau are unusually large, similar to those from the mid-latitudes (Davis et al., 2008; Slusher et al., 2010; Frey et al., 2013). Such large mixing ratios of NOₓ were found to arise from a combination of several factors: continuous sunlight, location at the bottom of a large air drainage basin, low temperatures leading to low primary production rates of HOₓ radicals, significant emissions of NOₓ from surface snow, and a shallow boundary layer (Davis et al., 2008; Frey et al., 2013, and refs. therein).

Snow emissions of NOₓ, observed at several polar locations (e.g. Jones et al., 2001; Honrath et al., 2000b), are driven by UV-photolysis of nitrate (NO₃⁻) in snow (Honrath et al., 2000b; Simpson et al., 2002) and are now considered to be an essential component of air-snow cycling of oxidised nitrogen species above the polar ice sheets (Davis et al., 2008; Frey et al., 2009b) and likely also above mid-latitude snow packs (Honrath et al., 2000a; Fisher et al., 2005). Atmospheric dynamics, i.e. vertical mixing strength and mixing height, can explain some of the observed temporal variability and site-specific chemical composition of the lower troposphere at South Pole and Summit, Greenland (Neff et al., 2008) (Neff et al., 2008; Van Dam et al., 2013). Recently, the very strong diurnal cycle of mixing ratios of NOₓ observed at Dome C, East Antarctic Plateau, during
summer was shown to result from the interplay between boundary layer mixing and emissions from the photochemical snow source; during calm periods a minimum of NO\textsubscript{x} mixing ratios occurred around local noon and a maximum in the early evening coinciding with the development and collapse of a convective boundary layer (Frey et al., 2013). A key parameter of the physical atmospheric processes at play is the turbulent diffusivity of the atmosphere, which controls the mixing height, $h_z$, of the atmospheric boundary layer and contributes to the magnitude of the flux of trace chemical species emitted by the snow (e.g. Frey et al., 2013).

The impact of NO\textsubscript{x} emissions from snow on the oxidation capacity of the lower troposphere in summer can be significant. For example, NO\textsubscript{x} snow emissions can result in net \textit{ozone}--O\textsubscript{3} production as observed in the interior of Antarctica (Crawford et al., 2001; Legrand et al., 2009; Slusher et al., 2010) as well as unusually large mixing ratios of hydroxyl radical levels as detected at South Pole (Davis et al., 2008, and refs. therein). Furthermore, in Antarctica the gas phase production of hydrogen peroxide (H\textsubscript{2}O\textsubscript{2}), the only major atmospheric oxidant preserved in ice cores, is sensitive to NO released by the surface snowpack (e.g. Frey et al., 2005, 2009a). A steady-state analysis of ratios of NO\textsubscript{2} : NO at Dome C suggested that mixing ratios of peroxy radicals (not measured at the time) are possibly larger at Dome C than any previous observations in air above polar snow (Frey et al., 2013).

The quantitative understanding of emissions of NO\textsubscript{x} from snow remains incomplete, but it is a research priority to be able to parameterise global models to assess for example global impacts of chemical air-snow exchange on tropospheric \textit{ozone}--O\textsubscript{3} (e.g. Zatko et al., 2013). Emissions of NO\textsubscript{x} from snow at Dome C are among the largest observed above either polar ice sheet, but are typically underestimated by models, especially at large solar zenith angles (Frey et al., 2013). One significant model uncertainty is the quantum yield of nitrate photolysis in snow, which is related to the location of the nitrate ion in snow grains, and needs to be better constrained by observations (Frey et al., 2013; Meusinger et al., 2014).

The study presented here was part of the comprehensive atmospheric chemistry campaign OPALE (Oxidant Production and its Export from Antarctic Lands) in East Antarctica (Preunkert et al., 2012) and provided the opportunity to measure NO\textsubscript{x} mixing ratios and flux
during a second summer season, after a previous campaign in 2009–2010 (Frey et al., 2013). The study objectives were firstly to extend the existing data set with mixing ratio profiles of the lower atmosphere and the firn air (interstitial air) column of the upper snow pack. Secondly, to investigate if observed NO$_2$ : NO ratios are consistent with measurements of hydroxyl and halogen radicals. And thirdly, to analyse the main drivers of the atmospheric NO$_x$ emission flux from snow.

2 Methods

The measurement campaign of 50 days took place at Dome C (75.1° S, 123.3° E, 3233 m) from 23 November 2011 to 12 January 2012. Similar to the 2009–2010 campaign atmospheric sampling was performed from an electrically heated lab shelter (Weatherhaven tent) located in the designated clean-air sector 0.7 km upwind (South) of Concordia station (see map in Frey et al., 2013) (Frey et al., 2013, Fig. 1a). All times are given as local time (LT), equivalent to UTC + 8 h, and during the study period the sun always remained above the horizon.

2.1 NO$_x$ concentration and flux measurements and uncertainties

Three 20 m-long intake lines (Fluoroline 4200 high purity PFA, I.D. 4.0 mm) were mounted on a mast about 15 m from the lab shelter into the prevailing wind to continuously sample air at 0.01, 1.00 and 4.00 m above the natural snow pack. The intake lines were away from the influence of the drifted snow around the lab shelter. On 9 January 2012 vertical profiles of the lower atmosphere were sampled by attaching a 100 m-long intake line to a helium-filled weather balloon, which was then manually raised and lowered. During selected time periods firn air was sampled, to depths 5–100 cm, by means of a custom built probe. The probe consisted of a tube (10 cm diameter) which was lowered vertically into a pre-cored hole to the chosen snow depth, passing through a disc (1 m diameter) resting on the snow surface. The disk had a lip of 10 cm protruding into the snow. The lip and
disk minimised preferential pumping of ambient air along the tube walls. The air intake was mounted at the bottom end of the vertical tube so that only air from the bottom and sides could enter, using small horizontal holes, could enter at 0–10 cm above the open bottom end of the vertical tube. All probe components were made from UV-transparent plastic (Plexiglas Sunactive GS 2458). Furthermore, 2 × 3 m sheets of UV-opaque (Acrylite OP-3) and UV-transparent (Acrylite OP-4) plexiglass, mounted on aluminium frames at 1 m above the snow surface, were used to deduce the effect of UV radiation on the mixing ratio of NOx in the interstitial air and avoid at the same time any temperature effect altering the snow surface.

To measure NOx, the same 2-channel chemiluminescence detector (CLD) and experimental set up as during the 2009–2010 campaign were used (instrument schematic in Frey et al., 2013) (Frey et al., 2013, Fig.1b). Channel one of the CLD measured atmospheric mixing ratios of NO whereas the other channel determined the sum of the mixing ratios of NO and NO originating from the quantitative photolytic conversion of NO2. The difference between the two channels was used to calculate atmospheric mixing ratios of NO2. The three sample inlets were connected inside the lab shelter to a valve box, which automatically switched the CLD between sampling heights on a 90 s duty cycle. As described below, the 10-minute average concentration difference ΔNOx between the 0.01 and 1.0 m inlets is used to estimate flux. Therefore, 10-minute mean ΔNOx values are calculated on average from two sets of two subsequent 90 s intervals, separated by a 90 s interval during which the 4.0 m inlet was measured. Baseline count rates were determined by adding excess O3 to sample air in a pre-chamber so that all electronically excited NO2 has returned to ground state when reaching the reaction chamber. The baseline was measured for 60 s every 13.5 min alternating between all three inlets. The NO sensitivity of the CLDs was determined every 14 h by standard addition to the sample air matrix of a 1 ppm NO/NO2 mixture (UK National Physics Laboratory traceable BOC certified), which is further diluted to 4 ppbv of NO. During standard runs also the conversion efficiency (CE) of the photolytic converter was determined by addition of a known mole fraction of NO2. This was achieved by gas phase titration of the NO/NO2 mixture to NO2 by O3 generated
from a pen-ray lamp, and monitoring the un-titrated NO mole fraction. The instrument artefact originating from NOx producing surface reactions in inlets and reaction cells was determined by overflowing the instrument inlet with scrubbed ambient air supplied by a pure air generator (Eco-Physics PAG003). The artefact was measured every 14 h, offset by 7 h to the calibration runs.

The CLD employed converts also mean wind direction during the measurement period was from S (176°) with an average speed of 4.0 m s⁻¹ (Fig. 1b). During 2.5% of the time winds came from the direction of Concordia station, i.e. the 355-15° sector (Frey et al., 2013, Fig.1a), potentially carrying polluted air from the station power generator to the measurement site. For example, during Period III. winds rotated 4 times through northerly directions (Fig. 1b). Pollution spikes in the raw 1-s data typically exceeded 10 ppbv of NOx and were effectively removed before computing the 1-min averages by applying a moving 1-min standard deviation (σ) filter. Observations were rejected when 1-σ of NO and NO2 mixing ratios within a 1-min window exceeded 24 and 90 pptv, respectively.

The CLD employed also converts nitrous acid (HONO) to NO in the photolytic converter and thus HONO sampled by the CLD is an interference, as discussed previously (Frey et al., 2013). Average mixing ratios of HONO at 1 m above the snowpack measured with the LOPAP (Long Path Absorption Photometer) technique were ∼35 pptv (Legrand et al., 2014). The corresponding downward correction for NO2 at 1 m above the snowpack is ∼5%. However the LOPAP technique may overestimate the mixing ratio of HONO owing to an interference with pernitric acid (HO2NO2) (Legrand et al., 2014). True corrections of NO2 inferred from modelled HONO mixing ratios (Legrand et al., 2014) are more likely to be on the order of <1.5%. Due to the uncertainty in absolute mixing ratios of HONO, no correction of NOx values for the HONO interference was applied.

The presence of strong gradients in mixing ratios of inferred by Legrand et al. (2014) can potentially lead to an overestimate of concentration differences between 0.01 and 1.0 used below to derive the vertical flux. During the OPALE campaign the atmospheric life time of NOx, τ(NOx), ranged between 3(12:00LT) and 7(00:00LT), whereas that of HONO, τ(HONO), ranged between 4.5(12:00LT) and 24(00:00LT) (Legrand et al., 2014). The life time of is comparable to
the typical transport times of \sim 10 between the surface and 1 at Dome C in summer (Frey et al., 2013). Hence, ratios as well as corresponding corrections required for are not constant with height above the snow surface. No gradients of mixing ratios were measured but modelled values were 18.8 and 10.2 at noon, and 15.3 and 12 at midnight, at 0.1 and 1.0, respectively (Legrand et al., 2014). Corresponding corrections of mean mixing ratios for are 1.3–1.5% with a maximum difference of 0.2% between 0.1 and 1.0. Thus, at Dome C a strong gradient in the mixing ratios of was a negligible effect on the mixing ratios of measured at 0.1 and 1.0 and thus a negligible effect on the estimated flux.

The thermal decomposition of HO$_2$NO$_2$ in the sample lines or photolytic converter of the CLD could also cause a positive bias of NO$_x$. Spike tests showed that the sample air residence time in the total volume of inlets and CLD is \sim 4 s (Frey et al., 2013). At a sample flow rate of 5.0 STP – L min$^{-1}$ the residence time in the combined volume of photolytic converter and CLD reaction cell is estimated to be < 2 s. Atmospheric lifetimes of $\tau_{\text{HO}_2\text{NO}_2}$, with respect to thermal decomposition to HO$_2$+NO$_2$ were calculated at mean ambient pressure (645 mb) using rate coefficients after Jacobson (1999). $\tau_{\text{HO}_2\text{NO}_2}$ decreases from 8.6 h at mean ambient temperature assumed in the sample intake lines (–30 °C) to 7 s at the maximum observed temperature in the photolytic converter (30 °C). Therefore, NO$_2$ production from HO$_2$NO$_2$ thermal decomposition is negligible in the sample intake lines, but approximately 25% of all HO$_2$NO$_2$ present may be converted to NO$_2$ in the photolytic converter. A recent airborne campaign above the East Antarctic Plateau showed mean summertime atmospheric mixing ratios of HO$_2$NO$_2$ between 0 and 50 m of 65 pptv with maxima about twice as large (Slusher et al., 2010). HO$_2$NO$_2$ present at these values could potentially produce 33–6616–32 pptv of NO$_2$ in the photolytic converter equivalent to 16–328–16% of the average NO$_2$ mixing ratio measured at 1 m. We on 5 January 2012 attempted to test for the presence of HO$_2$NO$_2$ by passing ambient air through a 50 m intake heated to 50 °C before it entered the CLD. However, during the tests no significant change in NO$_2$ was detected.

The presence of strong gradients in mixing ratios of HONO inferred by Legrand et al. (2014) can potentially lead to an overestimate of the NO$_x$ concentration
differences between 0.01 and 1.0 m used below to derive the vertical NO$_x$ flux. During the OPALE campaign the atmospheric life time of NO$_x$, $\tau_{NO_x}$, ranged between 3 h (12:00 LT) and 7 h (00:00 LT), whereas that of HONO, $\tau_{HONO}$, ranged between 4.5 min (12:00 LT) and 24 min (00:00 LT) (Legrand et al., 2014). The life time of HONO is comparable to the typical transport times of $\sim$10 min between the surface and 1 m at Dome C in summer (Frey et al., 2013). Hence, HONO: NO$_x$ ratios as well as corresponding corrections required for NO$_2$ are not constant with height above the snow surface. No gradients of HONO mixing ratios were measured but modelled values were 18.8 and 10.2 pptv at noon, and 15.3 and 12 pptv at midnight, at 0.1 and 1.0 m, respectively (Legrand et al., 2014). Corresponding corrections of mean NO$_2$ mixing ratios for HONO are 1.3–1.5 % with a maximum difference of 0.2 % between 0.1 and 1.0 m. Thus, at Dome C a strong gradient in the mixing ratios of HONO was a negligible effect on the mixing ratios of NO$_x$ measured at 0.1 and 1.0 m and thus a negligible effect on the estimated NO$_x$ flux.

2.2 NO$_x$ flux estimates

The turbulent flux of NO$_x$, $F_{NO_x}$, was estimated using the integrated flux gradient method (e.g. Lenschow, 1995) and mixing ratios of NO$_x$ measured at 0.01 and 1.0 m. $F_{NO_x}$ in the surface layer is parameterised according to the Monin–Obukhov similarity theory (MOST) whose predictions of flux-profile relationships at Halley, an Antarctic coastal site of the same latitude as Dome C, agree well with observations (Anderson and Neff, 2008, and references therein):

$$F_{NO_x} = -\frac{\kappa u_* z}{\Phi_h \left( \frac{z}{L} \right)} \frac{\partial c}{\partial z}$$

(1)

with the von Karman constant $\kappa$ (set to 0.40), friction velocity $u_*$, measurement height $z$, concentration gradient $\partial c/\partial z$, and $\Phi_h \left( \frac{z}{L} \right)$ an empirically determined stability function for heat with $L$ as the Monin–Obukhov length. Assuming constant flux across the layer between
the two measurement heights \( z_1 \) and \( z_2 \) allows the integration to be solved and yields:

\[
F_{\text{NO}_x} = - \frac{\int_{c_1}^{c_2} \kappa u_* \partial c}{\int_{z_1}^{z_2} \Phi_h \left( \frac{z}{L} \right) \frac{\partial z}{z}} = - \frac{\kappa u_* [c(z_2) - c(z_1)]}{\int_{z_1}^{z_2} \Phi_h \left( \frac{z}{L} \right) \frac{\partial z}{z}}
\]  

(2)

Stability functions \( \Phi_h \) used are given in Frey et al. (2013), while their integrated forms can be found in Jacobson (1999). Friction velocity \( u_* \) and \( L \) were computed from the three-dimensional wind components \((u, v, w)\) and temperature measured at 25 Hz by a sonic anemometer (Metek USA-1) mounted next to the uppermost \( \text{NO}_x \) intake line, at 4 m above the snow surface. Processing of raw data in 10 min blocks included temperature cross-wind correction and a double coordinate rotation to force mean \( w \) to zero (Kaimal and Finnigan, 1994; Van Dijk et al., 2006). Equation (2) implies that a positive flux is in upward direction, equivalent to snow pack emissions and a negative flux is in downward direction, equivalent to deposition.

The application of MOST requires the following conditions to be met: (a) flux is constant between measurement heights \( z_1 \) and \( z_2 \), (b) the lower inlet height \( z_1 \) is well above the aerodynamic roughness length of the surface, (c) the upper inlet height \( z_2 \) is within the surface layer, i.e. below 10% of the boundary layer height \( h_z \) (Stull, 1988), and (d) \( z_1 \) and \( z_2 \) are far enough apart to allow for detection of a significant concentration difference \([c(z_2) - c(z_1)]\). During summer at Dome C conditions——

Condition (a) is met in the surface layer if the chemical lifetime \( \tau_{\text{chem}} \) of \( \text{NO}_x \) is much longer than the turbulent transport time scale \( \tau_{\text{trans}} \). Based on observed \( \text{OH} \) and \( \text{HO}_2 \) the \( \tau_{\text{chem}} \) for \( \text{NO}_x \) is estimated to be 3 h at 1200 LT and 7 h at 0000 LT during OPALE (Legrand et al., 2014). Estimating \( \tau_{\text{trans}} \) following the approach described previously (Frey et al., 2013, Eq. 6 and 7) yields 0.6, 1.7 and \((b)\) are 2.5 min during the day (0900-1700 LT), the typical time of BL collapse (1700-1900 LT) and during the night (1900-0900 LT), respectively. Thus, \( \tau_{\text{chem}} \) exceeds \( \tau_{\text{trans}} \) by at least a factor 100, confirming that vertical mixing always dominates over the gas phase photochemical sink and flux can be assumed constant between the two inlets. Condition (b) is met as discussed in Frey et al. (2013). For (c) the upper inlet height of 1 m is compared to estimates of mix-
ing height $h_z$ from the MAR model (Gallée et al., 2015). The MAR model has been validated previously over the Antarctic Plateau, focusing on Dome C, during winter (Gallée and Gorodetskaya, 2010) and now also during summer (Gallée et al., 2015). Calculated flux values of NO$_x$ were removed when $h_z < 10$ m resulting in the removal of 22% (773 values) of all available 10 min flux averages. Flux estimates are removed specifically during the evening and night, when the BL is shallow. Hence, fluxes during night time are less well constrained, but nevertheless support a significant diurnal cycle (Frey et al., 2013, Fig. 6b,g and Fig. 9). For (d) 10 min averages of $[c(z_2) - c(z_1)]$ not significantly different from zero, i.e. smaller than their respective 1-σ standard error, were not included in the calculation of the flux of NO$_x$. The 1-σ standard error in $[c(z_2) - c(z_1)]$ was determined by error propagation of the 1-σ standard error of NO$_x$ mixing ratios. A total of 8% (303 values) of all available 10 min flux averages were not significantly different from zero and thus removed.

In summary, the restrictions imposed by MOST and NO$_x$ measurement uncertainty justify placing inlets at 0.01 and 1.0 m and lead to the removal of 30% (1076 values) of all available flux estimates. The total uncertainty of the 10 min NO$_x$ flux values due to random error in $[c(z_2) - c(z_1)]$ (31%), $u_*$ (3% after Bauguitte et al., 2012) and measurement height (error in $\ln(z_2/z_1)$ of $\sim$7%) amounts to 32%.

### 2.3 MAX-DOAS observations

Scattered sunlight was observed by a ground-based UV-visible spectrometer, in order to retrieve bromine oxide (BrO) column amounts. The instrument was contained in a small temperature-controlled box, which was mounted onto a tripod at 1 m above the snow surface. An external gearbox and motor scanned the box in elevation (so-called Multiple Axis). Spectra were analysed by Differential Optical Absorption Spectroscopy (DOAS), the combination being known as the MAX-DOAS technique. See Roscoe et al. (2014) for more details of apparatus and analysis. Briefly, the observed spectrum contains Fraunhofer lines from the Sun’s atmosphere, which interfere with absorption lines in the Earth’s atmosphere and are removed by dividing by a reference spectrum. The amounts of absorbers in the Earth’s
atmosphere are found by fitting laboratory cross-sections to the ratio of observed to reference spectra, after applying a high-pass filter in wavelength (the DOAS technique).

In our case the spectral fit was from 341 to 356 nm, and the interfering gases $O_3$, $O_4$ (oxygen dimer) and $NO_2$ were included with BrO. The analysis was done with two reference spectra, one from near the start of the campaign in December, the other following the addition of a snow excluder in January, necessary because it also contained a blue glass filter with very different spectral shape. The analysis was restricted to cloud-free days or part-days. In MAX-DOAS geometry, the stratospheric light path is almost identical in low-elevation and zenith views, so stratospheric absorption is removed by subtracting simultaneous zenith amounts from low-elevation slant amounts, important for BrO as there is much in the stratosphere.

To find the vertical amounts of BrO radicals the MAX-DOAS measurements were evaluated as follows: we divided by the ratio of the slant path length to the vertical (the Air Mass Factor, AMF), calculated by radiative transfer code (Mayer and Kylling, 2005), assuming all the BrO was in the lowest 200 m.

2.4 Ancillary measurements and data

Other co-located atmospheric measurements included mixing ratios of OH radicals and the sum of peroxy radicals ($RO_2$) at 3 m using chemical ionisation mass spectrometry (Kukui et al., 2014) and mixing ratios of $O_3$ at 1 m with a UV absorption monitor (Thermo Electron Corporation model 49I, Franklin, Massachusetts). Photolysis rate coefficients, $J$, were determined based on actinic flux, $I$, measured at $\sim 3.50$ m above the snow surface using a Met-Con $2\pi$ spectral radiometer equipped with a CCD detector and a spectral range from 285 to 700 nm (further details in Kukui et al., 2014). Total column $O_3$ above Dome C was taken from ground based SAOZ (Système d’Analyse par Observation Zenitale) observations (http://saoz.obs.ujf-grenoble.fr/SAOZ_consol_v2.html). Standard meteorology was available from an automatic weather station (AWS) at 0.5 km distance and included air temperature (Vaisala PT100 DTS12 at 1.6 m), relative humidity (at 1.6 m), wind speed and direction (Vaisala WAA 15A at 3.3 m). Photolysis rate coefficients, $J$, were determined based on actinic flux, $I$, measured at $\sim 3.50$ above the snow surface with a Met-Con $2\pi$ spectral radiometer equipped...
with a CCD detector and a spectral range from 285 to 700 (see also Kukui et al., 2014). The mixing height \( h_z \) of the atmospheric boundary layer was calculated from simulations with the MAR model as the height where the turbulent kinetic energy decreases below 5% of the value of the lowest layer of the model (Gallée et al., 2015).

During this study NO\(_3^-\) concentrations in snow were measured every 2–3 days in the surface skin layer, i.e. in the top 0.5 cm of the snowpack, as well as in shallow snow pits within the clean-air sector. Snow NO\(_3^-\) concentrations were determined using clean sampling procedures and a continuous flow analysis technique (e.g. Frey et al., 2009b). Total column above Dome C was taken from ground-based SAOZ (Système d’Analyse par Observation Zenitale) observations. The mixing height \( h_z \) of the atmospheric boundary layer was calculated from simulations with the MAR model as the height where the turbulent kinetic energy decreases below 5%. Samples were stored together with the additional snow samples discussed in Berhanu et al. (2014) and then analysed for NO\(_3^-\) in batches by the same operator. The precision is 5% of the value of the lowest layer of the model (Berhanu et al., 2014) based on replicate standard measurements. Due to a systematic shift in the NO\(_3^-\) standard response in between individual batch runs due to a calibration issue (Berhanu et al., 2014) the accuracy is larger than usual. The overall accuracy including systematic errors in calibration and collection of just the top few mm of snow is of the order of 20%, and is therefore comparable to the spatial variability of NO\(_3^-\) in surface snow at Dome C (France et al., 2011). For the discussion below it should be borne in mind that temporal changes of NO\(_3^-\) concentrations observed in surface snow are >50% (Fig. 7b) and therefore significantly larger than the measurement accuracy.

### 2.5 Modelling NO\(_3^-\) photolysis

The flux of NO\(_2\), \( F_{NO_2} \), from the snowpack owing to photolysis of the NO\(_3^-\) anion in the snowpack can be estimated as the depth-integrated photolysis rate of NO\(_3^-\)

\[
F_{NO_2} = \int_{z=0m}^{z=1m} [NO_3^-]_z J_z(NO_3^-) \, dz \tag{3}
\]
where $J_z(\text{NO}_3^-)$ is the photolysis rate coefficient of reaction $\text{NO}_3^- + h\nu \rightarrow \text{NO}_2 + \text{O}^-$ at depth, $z$, in the snowpack. $[\text{NO}_3^-]_z$ is the amount of nitrate--\text{NO}_3^- per unit volume of snow at depth, $z$, in the snowpack. $J(\text{NO}_3^-) - J_z(\text{NO}_3^-)$ is calculated as described in France et al. (2010) using a radiative transfer model, TUV-snow (Lee-Taylor and Madronich, 2002), to calculate irradiances within the snowpack as a function of depth. The optical properties and detailed description of the Dome C snowpack are reported in France et al. (2011). Values of depth-integrated flux were calculated as a function of solar zenith angle and scaled by values of $J(\text{NO}_3^-)$ measured by the Met-Con $2\pi$ spectral radiometer described above to account for changing sky conditions. Scaling by a measured value of $J(\text{NO}_3^-)$ is more accurate than previous efforts of scaling with a broad band UV instrument (e.g. France et al., 2011). The quantum yield and the absorption spectrum for nitrate--\text{NO}_3^- photolysis in snow were taken from Chu and Anastasio (2003). For the discussion below it should be borne in mind that the calculated $F_{\text{NO}_2}$ is a potential emission flux assuming that $\text{NO}_2$ is vented immediately after release from the snow grain to the air above the snow pack without undergoing any secondary reactions.

3 Results and discussion

3.1 NO$_x$ observations in ambient and firn air

In summer 2011–2012 atmospheric mixing ratios of NO$_x$ with strong diurnal variability were observed (Fig. 1c), similar to the 2009–2010 season, and showed maximum median levels in firn air of $\sim 3837$ pptv, which rapidly decreased to 319 pptv at 0.01 m and 213 pptv at 1.0 m (Table 1). As seen previously at Dome C and other locations, NO$_x$ mixing ratios were weakly but significantly anti-correlated with wind speed (e.g. at 1.0 m $R^2 = 0.14$, $p < 0.001$) $R = -0.37$, $p < 0.001$), especially when only the time period of the daily collapse of the convective boundary layer, i.e. 1700–1900 LT, was considered ($R = -0.45$, $p < 0.001$), and their diurnal cycle was dampened during storms (Fig. 1b–c).
The two main differences between summer 2011–2012 and summer 2009–2010 are a strong intra-seasonal trend and larger atmospheric mixing ratios. A significant increase of NO\textsubscript{x} mixing ratios at 1.0 m from low values in late November 2011 (Period I. (23–30 November 2011)) occurred in two steps: a small rise in the first week of December (Period II. (1–8 December 2011)), followed by a strong increase of daily averages from 300 to 1200 pptv during at the beginning of Period III. (9–11 December 2011 (Period III.)) (Fig. 1c). After that NO\textsubscript{x} mixing ratios gradually dropped over 10 days (Periods III.–IV.) to late November average concentrations of \(\overset{\sim}{-120}\) median concentrations of \(\sim120\) pptv, slightly lower than observed in late November (Fig. 1c, Table 2). During Period III. (9–22 December 2011 (Period III.)) the median concentration of NO\textsubscript{x} at 1.0 m was 451 pptv, about 2.5 times that during the same time period in 2009, but similar thereafter (Fig. 1c, Table 2).

The NO\textsubscript{x} fluxes, \(F_{\text{NO}_x}\), between 0.01 and 1.0 m were mostly emissions from the snow surface, with a median (range) of 1.6 (0.4–2.9) \(\times 10^{13}\) molecule m\(^{-2}\) s\(^{-1}\). Median values of \(F_{\text{NO}_x}\) at midnight and at noon were 0.4 and 2.9 \(\times 10^{13}\) molecule m\(^{-2}\) s\(^{-1}\), respectively (Table 1). During Period III. \(F_{\text{NO}_x}\) showed an increase by a factor 3, approximately around the same time when atmospheric mixing ratios of NO\textsubscript{x} increased (Period III.) (Fig. 1d, Table 1). The median flux of NO\textsubscript{x} during 9–22 December 2011 reached 3.1 \(\times 10^{13}\) molecule m\(^{-2}\) s\(^{-1}\), about almost 5 times the season median from 2009–2010. During 23 December to 12 January (Period IV.) the median flux of NO\textsubscript{x} in 2011–2012 was about twice that observed in 2009–2010 (Table 2). Potential causes of significant variability in mixing ratios and flux on seasonal time scales are discussed in Sect. 3.5.

### 3.2 The lower atmosphere–firn air profile

On 9 January 2012 a total of 12 vertical atmospheric profiles of NO\textsubscript{x} mixing ratios were measured between 11:30 and 23:30 LT. The lower 100 m of the atmosphere appear well mixed throughout the afternoon, with modelled mixing heights \(h_z\) of 200–550 m and observed turbulent diffusion coefficients of heat \(K_h\) of \(\sim0.1\) m\(^2\) s\(^{-1}\) (Fig. 2). However, in the late afternoon \(K_h\) values decreased gradually over a few hours to reach in the evening levels half those during the day thereby giving evidence of strongly reduced vertical mixing. Furthermore,
around 18:30 LT modelled $h_z$ values decreased within minutes from 550 to < 15 m height (Fig. 2a) illustrating the collapse of the convective boundary layer typically observed at Dome C in the early evening during summer (King et al., 2006). At Dome C rapid cooling of the surface in the evening results in a strong shallow surface inversion (e.g. Frey et al., 2013), and is illustrated by a decrease in downward long-wave radiation and a negative heat flux, as observed in the evening of 9 January 2012 (Argentini et al., 2014, Fig.4). It follows that NO$_x$ snow emissions are trapped near the surface and caused, which then leads to a significant increase in NO$_x$ mixing ratios below 15 m height measured almost immediately after collapse of the boundary layer (Fig. 2). During 22:20–22:40 LT a small increase in $K_v$, due to the nightly increase in wind shear (see Frey et al., 2013), was sufficient to cause upward mixing of NO$_x$ accumulated near the surface to $\sim$ 35 m height (Fig. 2). The vertical balloon soundings further underline the unique geographical setting of Dome C or other sites of similar latitude on the East Antarctic Plateau where air chemistry is dominated by strong diurnal cycles, both in down-welling solar radiation and atmospheric stability, contrasting South Pole where diurnal changes are absent and changes are more due to synoptic variability (Neff et al., 2008).

A vertical profile of mixing ratios of NO$_x$ and O$_3$ in firn air was measured on 12 January 2012 between 10:00 and 18:00 LT, for which depths were sampled in random order for 30–60 min each. Mixing ratio maxima of NO and NO$_2$ were $\sim$ 1 and 4 ppbv, respectively, about one order of magnitude above ambient air levels (Table 1), and occurred at 10–15 cm depth, slightly below the typical e-folding depth of 10 cm of wind pack snow at Dome C (France et al., 2011) (Fig. 3a). NO dropped off quickly with depth, reaching 55 pptv at 85 cm, whereas NO$_2$ decreased asymptotically approaching $\sim$ 2 ppbv (Fig. 3a). NO$_3^-$ concentrations in snow under the firn air probe did not follow the exponential decrease with depth typically observed at Dome C (e.g. Erbland et al., 2013), but. The firn air probe was installed onto untouched snow, and only removed after the end of the atmospheric sampling period. Thus contamination due to local activity appears unlikely, but a local anomaly remains a possibility as snow pits 5 m next to the lab shelter showed a similar increase of concentration with depth (data not shown). But NO$_3^-$ values within one e-folding depth were still in the
range measured further away (Profiles P1–P3 in Fig. 3a), justifying a discussion of vertical trends of mixing ratios.

$O_3$ mixing ratios in firn air were always below ambient air levels, suggesting the snow pack to be an $O_3$ sink as observed previously for the snowpack on the Greenland ice sheet (Peterson and Honrath, 2001), and showed a significant anti-correlation with $NO_2$ ($R^2 = 0.7, R = -0.84, p < 0.001$). This is further evidence for significant release of $NO_x$ by the snow matrix into the interstitial air, which then titrates $O_3$ through the reaction $NO + O_3 \rightarrow NO_2 + O_2$ (Fig. 3). In particular, the elevated concentration drop of $O_3$ mixing ratios by $>10$ ppbv at $45$ cm depth was not an outlier, as indicated by the collocated drop of $NO_2$ mixing ratios were also significantly elevated compared to adjacent snow layers, but was possibly due to snow with large concentrations (Fig. 3a). However, no $NO_3^-$ measurements were available to further investigate the origin of the $NO_2$ peak. The observed vertical trends in $NO_x$ suggest that below a few e-folding depths the open pore space of the upper snowpack holds a significant reservoir of $NO_2$ produced photolytically above, as hypothesized previously (Frey et al., 2013). In contrast, $NO$ disappears at depths devoid of UV irradiance as it reacts with $O_3$.

### 3.3 Response to UV irradiance

Changes in surface downwelling UV irradiance lead to a quick response of mixing ratios and speciation of $NO_x$ in ambient and firn air as observed during a partial solar eclipse and during a shading experiment (Fig. 4). The solar eclipse occurred early in the season, on 25 November 2011, and caused a decrease in ambient NO mixing ratios at $1.0$ m by about $10$ pptv or $10\%$, whereas $NO_2$ mixing ratios did not change significantly (Fig. 4a and b). The NO gas phase source, UV photolysis of $NO_2$, is reduced during the solar eclipse. But the sink of NO, the fast titration with $O_3$, is unaffected by the reduction in UV irradiance. During the shading experiment on 11 January 2012 plastic sheets were placed at $1$ m above the snow surface, alternating in $30$ min intervals between UV-opaque and UV-transparent materials. The impact of blocking incident UV irradiance (wavelengths $< 380$ nm) on firn air mixing ratios at $10$ cm snow depth was up to $300$ pptv or $30\%$ decrease in mixing ratios.
of NO, whereas mixing ratios of NO\(_2\) increased at the same time by \(\sim 150\) pptv or 5\%, although often not statistically significant (Fig. 4c and d). Similar to the solar eclipse, the behavior of NO\(_x\) mixing ratios in firn air is in accordance with a disruption of the fast gas phase interconversion of NO\(_x\) species. Decrease of NO and increase of NO\(_2\) mixing ratios are consistent with the suppression of NO\(_2\) photolysis, which is both a NO source and a NO\(_2\) sink.

Most importantly varying incident UV irradiance in the wavelength region of NO\(_3^-\) absorption (action spectrum maximum at 320 nm) over half-hourly time scales does not cause a depletion of NO\(_2\) in firn air even though NO\(_2\) is the main product of NO\(_3^-\) photolysis in the snowpack. A dampened UV response of NO\(_2\) mixing ratios suggests that the NO\(_x\) reservoir present in the open pore space of the upper snow pack discussed above must be large as it is not depleted during 30 min filter changes at the sample pump rates used. One implication is that the impact of changes in incident UV irradiance on the snow source and thus NO\(_x\) flux and mixing ratios is only observable on diurnal and seasonal time scales.

3.4 NO\(_2\) : NO ratios, peroxy and halogen radicals

In 2011–2012 the NO\(_2\) : NO ratios at 1.0 m were up to 3 times larger than in 2009–2010 (Table 2). A previous steady-state analysis indicated that high peroxy and possibly halogen radical levels must be present to explain large deviations from the simple Leighton steady-state (Frey et al., 2013). During summer 2011–2012 median concentrations of RO\(_2\) radicals at 3 m, thought to consist mainly of HO\(_2\) and CH\(_3\)O\(_2\), were \(9.9 \times 10^7\) molecule cm\(^{-3}\) (Kukui et al., 2014). To find the vertical amounts of radicals the MAX-DOAS measurements were evaluated as follows: we divided by the ratio of the slant path length to the vertical (the Air Mass Factor, AMF), calculated by radiative transfer code (Mayer and Kylling, 2005), assuming all the was in the lowest 200.

Figure 5 shows the BrO results, where the apparent vertical amounts at 15\(^\circ\) are much larger than those at lower elevations – the AMFs are incorrect, and interestingly, as at Halley in 2007 (Roscoe et al., 2014), much of the BrO must be in the free troposphere. The average of BrO at the three elevations is about \(0.8 \times 10^{13}\) molecule cm\(^{-2}\), with a slight decrease
during the campaign. The average at Halley in 2007 was about $2.5 \times 10^{13}$ molecule cm$^{-2}$, so mixing ratios of BrO at Dome C are about a third those at Halley. The Dome C data were not inverted to determine the mixing ratio near the surface, but the changes in slant column with elevation angle are similar to those at Halley in 2007 (Roscoe et al., 2014). So if the Halley inversion results are simply divided by 3 the Dome C values imply 2 to 3 pptv of BrO near the surface.

The same steady-state the total radical concentration $[\text{OX}] = [\text{HO}_2] + [\text{RO}_2] + 2[\text{XO}]$, with $\text{XO} = \text{BrO, ClO}$, can be calculated based on observed NO$_2$ : NO ratios and $\mathcal{J}(\text{NO}_2)$ (Ridley et al., 2000). Repeating the calculation as described by Frey et al. (2013) was repeated for austral summer 2011–2012 and yields an average of 2.6 in Frey et al. (2013) for 19 December 2011 to 9 January 2012 yields a median $[\text{OX}]$ of $2.2 \times 10^9$ molecule cm$^{-3}$ or 134186 pptv of total radical concentrations (Ridley et al., 2000). The steady-state estimates of Median $[\text{RO}_2]$ radicals are therefore $\sim$20 times those observed at mid-day by Kukui et al. (2014). While observed and radicals are too low to explain the large $\cdot [\text{HO}_2]$ of $9.9 \times 10^7$ molecule cm$^{-3}$ or 5 ratios it is possible that at least part of this pptv observed during the same period (Kukui et al., 2014) and 3 pptv of BrO yield $[\text{OX}]$ of 11 pptv. Hence, the estimated total radical concentration exceeds observations by a factor 10.3. To estimate the impact of a potential interference by HO$_2$NO$_2$ we corrected the NO$_2$ mixing ratios, assuming that additional NO$_2$ is measured in the CLD from HO$_2$NO$_2$ thermal decomposition, equivalent to 25% (100%) of ambient HO$_2$NO$_2$ on the order of 130 pptv. We then find that the median steady-state estimate of total oxidant concentrations is still a factor 9.6 (3.3) larger than that observed. Thus, only a part of the inconsistency may be due to overestimated from a potential interference with explained by the interference with HO$_2$NO$_2$ (not measured) as discussed above.

3.5 Drivers of seasonal NO$_x$ variability

On diurnal time scales NO$_x$ mixing ratios at Dome C are controlled by the interplay between snowpack source strength and atmospheric physical properties, i.e. turbulent diffusion coefficient of heat $K_h$ of heat $K_h$ and mixing height $h_z$ of the boundary
Layer (Frey et al., 2013). Indeed, the median diurnal cycles of NO$_x$ mixing ratios during 1–8 December 2011 (in 2011-12 show with the exception of Period II.) showed previously described behavior (Frey et al., 2013) with a strong maximum during 18:00–20:00 (1-8 December) previously described behavior (Frey et al., 2013), that is a strong increase around 1800 LT and a minimum at 01:00 LT, then during to maximum values, which last into the night time hours (Fig. 6a). Night-time peaks of NO$_x$ are plausible if the weakening of snow emissions is offset by a corresponding decrease of the chemical sink of NO$_x$, i.e., the NO$_2$ + OH reaction, assuming no significant change in $h_z$. This is consistent to a first order taking into account that observed OH concentrations (Kukui et al., 2014) and $F_{NO_x}$ vary in a similar way, by up to a factor 5 between local noon and midnight.

During Period III. (9–22 December 2011 (Period III.) at similar noon time values further increase of the primary maximum and generally large mixing ratios are similar to Period II. but the increase in the evening has a larger amplitude and generally larger mixing ratios prevail during night time (Fig. 6a). During 23 December 2011–2012 January 2012 (Period IV.) the diurnal cycle of increased NO$_x$ mixing ratios returned to low values and small diurnal amplitudes (Fig. 6a). As expected the during Period III. are consistent with the observed NO$_x$ emission flux $F_{NO_x}$ between 0.01 and 1.0, which always peaked at local noon, but showed during 9–22 December 2011 (also showed during Period III.) a strong increase at all times of the day with a near doubling of the noon time median (Fig. 6b). During Period IV. (23 December 2011–2012 January 2012) the diurnal cycles of both NO$_x$ mixing ratios and $F_{NO_x}$ returned to low values and small diurnal amplitudes (Fig. 6a–b).

Below we evaluate potential causes of the unusual variability in NO$_x$ mixing ratios and flux observed on seasonal time scales.

### 3.5.1 Atmospheric mixing vs. snow source strength

Similar to explaining diurnal NO$_x$ cycles at Dome C the seasonal trends–variability of daily mean NO$_x$ mixing ratios during the first half of December 2011 can be attributed to a combination of changes in $F_{NO_x}$ and $h_z$ (Fig. 1). The strong increase of NO$_x$ around 11 December 2011 falls into a period–Period when $F_{NO_x}$ almost tripled, but while wind speeds
slightly decreased and shallow boundary layer depths prevailed with daily $h_z$ maxima below 100–200 heights prevailed (Fig. 1). After, Table 2). For example, on 12 December and 13 December the modelled diurnal ranges of $h_z$ were 3.4–224 m and 3.6–251 m, respectively, while sodar observations yielded 10–150 m and 5–125 m, respectively (Gallée et al., 2015). After 13 December 2011 $F_{NO_x}$ remained at high values, thus, the decrease of NO$_x$ mixing ratios appears to be primarily caused by daily maximum stronger upward mixing into a larger volume, i.e. wind speeds increased and daily $h_z$ increasing to $\gg 600$ maxima grew, exceeding 600 m on 18 December –2011 (Fig. 1). After 23 December NO$_x$ mixing ratios drop to low levels, due to smaller $F_{NO_x}$ and a deep boundary layer (Fig. 1).

Trends in $F_{NO_x}$ are controlled by variability in depends on atmospheric turbulence ($K_h$) and concentration differences (\(\Delta NO_x\)), which in turn are determined by the strength of the photolytic snow pack source at a given $K_h$ (Eq. 21–2). However, the relative importance of $K_h$ and snow pack source strength can vary. For example, during Period IV. (23 December 2011–12 January 2012) the median $F_{NO_x}$ was $1.3 \times 10^{13}$ molecule m$^{-2}$ s$^{-1}$, about twice that observed during the same period in 2009–2010 (Fig. 6g; Table 2). The inter-seasonal difference can be explained by both, significantly larger atmospheric turbulence and more negative $\Delta NO_x$ during all times of the day 2011–2012 (Fig. 6h and i). Median $K_h$ was 0.08 m$^2$ s$^{-1}$, double that in 2009–2010, and median $\Delta NO_x$ was $-51$ pptv compared to $-32$ pptv in 2009–2010 (Table 2).

In contrast, during 2011–2012 the observed intra-seasonal variability of $F_{NO_x}$ is dominated by changes in the snow pack source strength. During Period III. (9–22 December 2011) median $K_h$ values ($\sim 0.05$ m$^2$ s$^{-1}$) and diurnal cycles were smaller than thereafter (Fig. 6c; Table 2), while $\Delta NO_x$ values were among the largest observed so far at Dome C, about three times those during the rest of the season, and therefore primarily caused the tripling of $F_{NO_x}$ (Fig. 6d and i). In section 3.5.2 we’ll discuss underlying causes of changes in the strength of the snow source.

Previously, non-linear HO$_x$-NO$_x$ chemistry and the associated increase in NO$_x$ lifetime were suggested to be an additional factor needed to explain large increases in NO$_x$ mixing ratios observed at South Pole (Davis et al., 2008, and references therein). In order to assess
the relevance of this factor at Dome C we apply a simple box model to estimate net $\text{NO}_x$ production rates as done previously (Frey et al., 2013). It is assumed that mixing is uniform and instantaneous, that the snow emission flux $F_{\text{NO}_x}$ is the main $\text{NO}_x$ source and the reaction with the OH radical is the dominant $\text{NO}_x$ sink and

$$\frac{d[\text{NO}_x]}{dt} \sim \frac{F_{\text{NO}_x}}{h_z} - k[\text{NO}_2][\text{OH}]$$

where $k$ is the respective reaction rate coefficient. In 2009–10 no OH observations were available at Dome C and average values from South Pole were used instead. In 2009–10 estimated net production rates of $\text{NO}_x$ at night were on the order of $100\,\text{pptv h}^{-1}$ and therefore explained the average increase in $\text{NO}_x$ from 110 to 300 pptv observed from 1700 to 1900 LT (Frey et al., 2013). In 2011–12 the same analysis is repeated using OH measurements available for most of Period IV. (Kukui et al., 2014) as well as $h_z$ calculated with the MAR model (Gallée et al., 2015). Resulting night time values of net $\text{NO}_x$ production rates are with about $40\,\text{pptv h}^{-1}$ smaller than in 2009-10 but again to a first order consistent with a smaller observed increase in $\text{NO}_x$ mixing ratios in the evening hours; i.e. during Period IV. median $\text{NO}_x$ increased between 1630 and 1930 LT from 114 to $242\,\text{pptv}$ ((Fig. 6a,f). The above model is oversimplified as the very likely presence of $\text{HO}_2\text{NO}_2$ will modulate the diurnal variability of $\text{NO}_x$ sinks and sources with an impact on $\text{NO}_x$ lifetime as suggested by Davis et al. (2008). However without any information on the diurnal cycle of $\text{HO}_2\text{NO}_2$ at Dome C further modelling is not warranted.

### 3.5.2 Snow source strength

The $\text{NO}_x$ flux observed above polar snow is on the order of $10^{12}$ to $10^{13}$ molecule m$^{-2}$ s$^{-1}$ and contributes significantly to the $\text{NO}_x$ budget in the polar boundary layer. At the lower end of the range are $F_{\text{NO}_x}$ observations at Summit, Greenland (Honrath et al., 2002) and at Neumayer in coastal Antarctica (Jones et al., 2001) with $2.5 \times 10^{12}$ molecule m$^{-2}$ s$^{-1}$, whereas on the Antarctic Plateau $F_{\text{NO}_x}$ values are up to ten times larger. For example, the
average $F_{\text{NO}_x}$ at South Pole during 26-30 November 2000 was $3.9 \times 10^{12}$ molecule m$^{-2}$ s$^{-1}$ (Oncley et al., 2004), whereas at Dome C observed fluxes are 2-6 times larger, with seasonal averages of $8-25 \times 10^{12}$ molecule m$^{-2}$ s$^{-1}$ (Frey et al., 2013, this work). Due to the uncertainties in the processes leading to NO$_x$ production it had been difficult to explain inter-site differences, e.g. by simply scaling $F_{\text{NO}_x}$ with UV irradiance and NO$_3^-$ in the surface snow pack (Davis et al., 2004). Some of the variability in flux values may be due to differences in experimental set up or in the employed flux estimation method (e.g. Davis et al., 2004; Frey et al., 2013). For example, the $F_{\text{NO}_x}$ estimates for South Pole are based on measured NO gradients only, inferring NO$_x$ from photochemical equilibrium and using the Bowen ratio method (Oncley et al., 2004), whereas the $F_{\text{NO}_x}$ estimates for Dome C are based on observations of both atmospheric nitrogen oxides (NO and NO$_2$) and the flux-gradient method (Frey et al., 2013).

Model predictions of $F_{\text{NO}_x}$ show in general a low bias on the Antarctic Plateau when compared to observations. A first 3-D model study for Antarctica included NO$_x$ snow emissions parameterised as a function of temperature and wind speed to match the observed $F_{\text{NO}_x}$ at South Pole (Wang et al., 2007). However, the model under-predicts NO mixing ratios observed above the wider Antarctic Plateau highlighting that the model lacks detail regarding the processes driving the emission flux (Wang et al., 2007). The first model study to calculate $F_{\text{NO}_x}$ based on NO$_3^-$ photolysis in snow, as described in this work, reports $1-1.5 \times 10^{12}$ molecule m$^{-2}$ s$^{-1}$ for South Pole in summer (Wolff et al., 2002), about a factor 4 smaller than the observations by Oncley et al. (2004) and up to 16 times smaller than what is needed to explain rapid increases in NO$_x$ mixing ratios over a few hours (Davis et al., 2008, and references therein). Recent model improvements reduced the mismatch with the South Pole flux observations and included the use of updated absorption cross sections and quantum yield of the NO$_3^-$ ion, as well as e-folding depths measured in surface snow on the Antarctic Plateau, and resulted in a factor 3 increase of flux calculated for South Pole (France et al., 2011). In light of major remaining uncertainties, which include the spatial variability of NO$_3^-$ in snow and the quantum yield of NO$_3^-$ photolysis (Frey et al., 2013), we discuss below the variability of $F_{\text{NO}_x}$ observed at Dome C.
A number of factors may contribute to changes in snow source strength of NO\textsubscript{x}. One possibility to explain increases in $F_{\text{NO}_x}$ is that the NO\textsubscript{2} reservoir in the open pore space of the upper snowpack discussed above may undergo venting upon changes in atmospheric pressure. However, no statistically significant relationship between $F_{\text{NO}_x}$ and atmospheric pressure is found (data not shown). The main cause of large $F_{\text{NO}_x}$ values appears rather to be related to changes in snow production rates of NO\textsubscript{x} from NO\textsubscript{3} photolysis, which depend on the NO\textsubscript{3} photolysis rate coefficient $J_{\text{NO}_3}$ and the NO\textsubscript{3} concentration in the photic zone of the snow pack (Eq. 3).

Trends in down-welling UV irradiance due to stratospheric O\textsubscript{3} depletion were suggested previously to drive $J_{\text{NO}_3}$ and therefore $F_{\text{NO}_x}$ and the associated increase in net production of surface O\textsubscript{3} observed at South Pole in summer since the 1990's (Jones and Wolff, 2003). At Dome C the observed increase in $F_{\text{NO}_x}$ and strongly negative $\Delta$NO\textsubscript{x} values coincided with a period when total column O\textsubscript{3} declined from $> 300$ to about 250 DU (Fig. 7a and c). During Period III. (9–22 December 2011) the median column O\textsubscript{3} was about 8% lower than during the time periods before and after (Table 2). However, associated changes in $J_{\text{NO}_3}$ on the order of $\sim 10\%$ are too small to account alone for the observed tripling in $F_{\text{NO}_x}$ (Fig. 6e; Table 2).

Instead changes in $F_{\text{NO}_x}$ can be linked to the temporal variability of NO\textsubscript{3} present in the snow skin layer. During time periods the end of Period II. and beginning of Period III. skin layer NO\textsubscript{3} concentrations were up to two times larger than before and after, coinciding with increased $F_{\text{NO}_x}$ (Fig. 7bandc). $F_{\text{NO}_x}$ is high during the end of Period II. and beginning of Period III., however drops off one week after the decrease of nitrate concentrations in surface snow (Fig. 7c). To confirm this the link between NO\textsubscript{x} emissions and NO\textsubscript{3} in snow $F_{\text{NO}_2}$ values were modelled (Eq. 3) based on observed $J_{\text{NO}_3}$, daily sampling of skin layer NO\textsubscript{3} and two depth profiles, at 100 m (P1) and 5 km (P2) distance from the lab shelter, in order to account for spatial and temporal variability of NO\textsubscript{3} concentrations of surface in snow. Modelled $F_{\text{NO}_2}$ capture some of the temporal trends in observational estimates of $F_{\text{NO}_x}$ confirming the link with $J_{\text{NO}_3}$ and NO\textsubscript{3} concentrations (Fig. 7c). However, median
ratios of observed $F_{\text{NO}_x}$ and modelled $F_{\text{NO}_2}$ values are 30–50 during Period III. and 15–30 during Period IV. (Fig. 7c).

Disagreement between model and observations was previously attributed to uncertainties in the quantum yield of NO$_3^-$ photolysis in natural snow (Frey et al., 2013). The model employed here uses a constant quantum yield, i.e. its value at the mean ambient temperature at Dome C ($-30$ °C) of 0.0019 (Chu and Anastasio, 2003). However, quantum yield may vary with time, as the same lab study reports a positive relationship between quantum yield and temperature (Chu and Anastasio, 2003). Comparison of time periods before and after 18 December 2011 shows an increase of mean air temperature from $-34.2$ °C to $-27.7$ °C and a decrease of its mean diurnal amplitude from 13 to 9.7 K (Fig. 1a). However, observations of $F_{\text{NO}_x}$ showed behaviour opposite to that expected from a temperature driven quantum yield, i.e. $F_{\text{NO}_x}$ values decreased as air temperature increased (Fig. 1a and d). Yet, the large diurnal amplitude of air temperature at Dome C could explain diurnal changes of $F_{\text{NO}_x}$ by a factor 1.5–1.75. The temperature effect is however small when compared to the observed up to 20-fold change between night and day in $F_{\text{NO}_x}$, which are driven by actinic flux. A recent lab study found that the quantum yield of photolytic loss of nitrate–NO$_3^-$ from snow samples collected at Dome C decreased from 0.44 to 0.003 within what corresponds to a few days of UV exposure in Antarctica (Meusinger et al., 2014). The authors argue that the observed decrease in quantum yield is due to nitrate–NO$_3^-$ being made of a photo-labile and a photo-stable fraction, confirming a previous hypothesis that the range of quantum yields reflects the location of NO$_3^-$ within the snow grain and therefore availability to photolysis (Frey et al., 2013) (Davis et al., 2008; Frey et al., 2013). Thus, observed the $F_{\text{NO}_x}$ values observed at Dome C fall well within the range of predictions based on quantum yield values measured in Dome C snow samples from the same site, which exceed that used in the current model by a factor 2–200. A systematic decrease in quantum yield due to depletion of photo-labile NO$_3^-$ in surface snow may have contributed to the observed decrease in $F_{\text{NO}_x}$ after 22 December 2011. However, a lack of information on snow grain morphology or NO$_3^-$ location within the snow grain limits further exploration of the impact of a time variable quantum yield on $F_{\text{NO}_x}$. It should be noted that during 2009–2010 large
skin layer NO$_3^-$ values did not result in $F_{\text{NO}_x}$ values comparable to those in 2011–2012 which may be due to a different partitioning between \textit{photolabile} and \textit{photostabile} \textit{photo-labile} and \textit{photo-stable} NO$_3^-$ in surface snow (Fig. 7b and c; Table 2).

The consequences of large NO$_x$ fluxes consist not only in contributing to high NO$_x$ mixing ratios but also in influencing local O$_3$ production, as suggested by significantly higher surface O$_3$ mixing ratios ($>$ 30 ppbv) during 9–22 December in 2011–2012 (Period III.) \textit{than compared to} 25 ppbv in 2009–2010 (Fig. 7d).

4 Conclusions

Measurements of NO$_x$ mixing ratios and flux carried out as part of the OPALE campaign at Dome C in 2011–2012 allowed to extend the existing data set from a previous campaign in 2009–2010.

Vertical profiles of the lower 100 m of the atmosphere confirm that at Dome C \textit{strong diurnal cycles}, both in down-welling solar radiation and atmospheric stability, \textit{large diurnal cycles in solar irradiance and a sudden collapse of the atmospheric boundary layer in the early evening} control the variability of NO$_x$ mixing ratios and flux. In contrast, at South Pole diurnal cycles are absent and changes more due to synoptic variability (Neff et al., 2008). \textit{Large mixing ratios of NO$_x$ at Dome C arise from a combination of several factors: continuous sunlight, large NO$_x$ emissions from surface snow and shallow mixing depths after the evening collapse of the convective boundary layer. Unlike at South Pole it is not necessary to invoke non-linear HO$_x$-NO$_x$ chemistry to explain increases in NO$_x$ mixing ratios. However, uncertainties remain regarding atmospheric levels of HO$_2$NO$_2$ and its impact on NO$_x$ life time being a temporary NO$_x$ reservoir.} Understanding atmospheric composition and air-snow interactions in inner Antarctica requires studies at both sites as they together encompass the spectrum of diurnal variability expected across the East Antarctic Plateau (King et al., 2006; Frey et al., 2013).

Firn air profiles suggest that the upper snow pack at Dome C is \textit{an O$_3$ sink and holds below a few e-folding depths a significant reservoir of NO$_2$ produced photolytically above,}
whereas NO disappears at depths devoid of UV as it reacts with O$_3$. Shading experiments showed that the presence of such a NO$_2$ reservoir dampens the response of NO$_x$ mixing ratios above or within the snowpack due to changes in down-welling UV irradiance on hourly time scales. Thus, systematic changes in NO$_x$ mixing ratios and flux due to the impact of UV on the snow source are only observable on diurnal and seasonal time scales.

First-time observations of BrO at Dome C suggest 2–3 pptv near the ground, with higher levels in the free troposphere similar to Halley, possibly originating from a sea ice source in coastal Antarctica (Theys et al., 2011) or from stratospheric descent (Salawitch et al., 2010). Assuming steady-state observed mixing ratios of BrO and RO$_2$ radicals are about a factor ten too low to explain the large NO$_2$ : NO ratios found measured in ambient air. It is possible that likely present The likely presence of HO$_2$NO$_2$ at Dome C but not measured during OPALE, (not measured) may cause an overestimate of NO$_2$ with the detection method employed and may therefore explain explain a part of this inconsistency.

During 2011–2012 NO$_x$ mixing ratios and flux were larger than in 2009–2010 consistent with also larger surface O$_3$ mixing ratios resulting from increased net O$_3$ production. Large NO$_x$ mixing ratios and significant variability during December 2011 were attributed to a combination of changes in mixing height and NO$_x$ snow emission flux $F_{NO_x}$. Trends in $F_{NO_x}$ were found to be controlled by atmospheric turbulence and the strength of the photolytic snowpack source, of which the relative importance may vary in time. Larger median $F_{NO_x}$ values in 2011–2012 than those during the same period in 2009–2010 can be explained by both, significantly larger atmospheric turbulence and a slightly stronger snowpack source. However, the tripling of $F_{NO_x}$ in December 2011 was largely due to changes in snow pack source strength driven primarily by changes in NO$_3^-$ concentrations in the snow skin layer, and only to a secondary order by decrease of total column O$_3$ and associated increase in NO$_3^-$ photolysis rates. Median ratios of observed $F_{NO_x}$ and modelled $F_{NO_2}$ values ranged from 15 to 50 using the quantum yield of NO$_3^-$ photolysis reported by Chu and Anastasio (2003). Model predictions based on quantum yield values measured in a recent lab study on Dome C snow samples (Meusinger et al., 2014) yield 2–200 fold larger $F_{NO_2}$ values encompassing observed $F_{NO_x}$. In particular, a decrease in quantum yield due to depletion of photo-labile NO$_3^-$ in sur-
face snow may have contributed to the observed decrease in $F_{\text{NO}_x}$ after 22 December 2011. Yet in 2009–2010 large skin layer $\text{NO}_3^-$ values did not result in elevated $F_{\text{NO}_x}$ values as seen in 2011–2012 possibly due to different partitioning of $\text{NO}_3^-$ between a photo-labile and photo-stabile fraction. In summary the seasonal variability of $\text{NO}_x$ snow emissions important to understand atmospheric composition above the East Antarctic Plateau depends not only on atmospheric mixing but also critically on $\text{NO}_3^-$ concentrations and availability to photolysis in surface snow, as well as incident UV irradiance. Future studies need to address on the Antarctic Plateau need to reduce uncertainties in $\text{NO}_2$ measurements, obtain also observations of $\text{HO}_2\text{NO}_2$ and assess how quantum yield of $\text{NO}_3^-$ photolysis varies in time in snow varies as a function of snow chemical and physical properties: thereby obtaining a more detailed view on the dynamics in the vertical redistribution of across the sunlit snowpack driven by photolysis and redeposition (e.g. Frey et al., 2009b). This is important to be able to close the mass budget of reactive nitrogen species between atmosphere and snow above Antarctica.

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Table 1. NO\textsubscript{x} mixing ratios and flux at Dome C during 23 November 2011–12 January 2012.

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<tr>
<th>Parameter</th>
<th>$z$, m</th>
<th>mean ±1σ</th>
<th>median</th>
<th>$t_{\text{total, days}}^a$</th>
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<tbody>
<tr>
<td>NO, pptv</td>
<td>$-0.1^b$</td>
<td>1097 ± 795</td>
<td>879</td>
<td>2.9</td>
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<tr>
<td></td>
<td>0.01</td>
<td>121 ± 102</td>
<td>94</td>
<td>18.6</td>
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<tr>
<td></td>
<td>1.0</td>
<td>98 ± 80</td>
<td>77</td>
<td>24.4</td>
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<tr>
<td></td>
<td>4.0</td>
<td>93 ± 68</td>
<td>78</td>
<td>13.7</td>
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<tr>
<td>NO\textsubscript{2}, pptv</td>
<td>$-0.1^b$</td>
<td>4145 ± 2667</td>
<td>2990</td>
<td>2.6</td>
</tr>
<tr>
<td></td>
<td>0.01</td>
<td>328 ± 340</td>
<td>222</td>
<td>17.6</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
<td>211 ± 247</td>
<td>137</td>
<td>23.2</td>
</tr>
<tr>
<td></td>
<td>4.0</td>
<td>210 ± 199</td>
<td>159</td>
<td>12.8</td>
</tr>
<tr>
<td>NO\textsubscript{x}, pptv</td>
<td>$-0.1^b$</td>
<td>5144 ± 3271</td>
<td>3837</td>
<td>2.6</td>
</tr>
<tr>
<td></td>
<td>0.01</td>
<td>447 ± 432</td>
<td>319</td>
<td>17.5</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
<td>306 ± 316</td>
<td>213</td>
<td>23.2</td>
</tr>
<tr>
<td></td>
<td>4.0</td>
<td>302 ± 259</td>
<td>241</td>
<td>12.8</td>
</tr>
<tr>
<td>$F$–NO\textsubscript{x} \times 10^{13}$ molecule m\textsuperscript{-2}s\textsuperscript{-1}$ ^c</td>
<td>0.01–1.0</td>
<td>2.5 ± 8.2</td>
<td>1.6</td>
<td>17.4</td>
</tr>
<tr>
<td>$F$–NO\textsubscript{x} \times 10^{13}$ molecule m\textsuperscript{-2}s\textsuperscript{-1}, local noon</td>
<td>0.01–1.0</td>
<td>5.0 ± 8.2</td>
<td>2.9</td>
<td>1.1</td>
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<tr>
<td>$F$–NO\textsubscript{x} \times 10^{13}$ molecule m\textsuperscript{-2}s\textsuperscript{-1}, local midnight</td>
<td>0.01–1.0</td>
<td>0.3 ± 1.6</td>
<td>0.4</td>
<td>0.2</td>
</tr>
</tbody>
</table>

$^a$ Total sample time estimated as the sum of all 1 min intervals.
$^c$ 1 December 2011–12 January 2012.
Table 2. Seasonal evolution of median NO$_x$ mixing ratios and flux along with relevant environmental parameters at Dome C in summer 2011–2012 (time periods I.–IV. highlighted in Fig. 1 and 7) and comparison to summer 2009–2010 (from Frey et al., 2013).

<table>
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<tr>
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<td>NO$_x$ (pptv)$^a$</td>
<td>180</td>
<td>324</td>
<td>451</td>
<td>122</td>
<td>183</td>
<td>145</td>
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<td>F×NO$_x$ × 10$^{13}$ (molecule m$^{-2}$ s$^{-1}$)$^b$</td>
<td>–</td>
<td>0.94</td>
<td>3.10</td>
<td>1.30</td>
<td>–</td>
<td>0.66</td>
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<tr>
<td>ΔNO$_x$ (pptv)$^b$</td>
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<td>-63</td>
<td>-153</td>
<td>-51</td>
<td>–</td>
<td>-32</td>
</tr>
<tr>
<td>NO$_2$ : NO$^a$</td>
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<td>1.5</td>
<td>2.8</td>
<td>2.0</td>
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<td>-31.0</td>
<td>-27.4</td>
<td>-31.5</td>
<td>-30.9</td>
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<td>wind speed (m s$^{-1}$)</td>
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<td>3.0</td>
<td>2.5</td>
<td>3.8</td>
<td>2.4</td>
<td>2.2</td>
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<td>$K_h$ (m$^2$ s$^{-1}$)</td>
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<td>0.046</td>
<td>0.049</td>
<td>0.080</td>
<td>-</td>
<td>0.043</td>
</tr>
<tr>
<td>$h_z$ (m)$^c$</td>
<td>-</td>
<td>19</td>
<td>20</td>
<td>36</td>
<td>6–59</td>
<td>18–25</td>
</tr>
<tr>
<td>$J_{NO_x}$ × 10$^{-8}$ (s$^{-1}$)</td>
<td>-</td>
<td>-</td>
<td>2.93</td>
<td>2.68</td>
<td>-</td>
<td>-</td>
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<tr>
<td>SZA ($^\circ$)</td>
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<td>68.1</td>
<td>67.6</td>
<td>67.9</td>
<td>67.6</td>
<td>67.9</td>
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<td>column O$_3$ (DU)</td>
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<td>294</td>
<td>272</td>
<td>297</td>
<td>311</td>
<td>309</td>
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<tr>
<td>NO$_3$ skin layer (ng g$^{-1}$)$^d$</td>
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<td>764</td>
<td>1090</td>
<td>439</td>
<td>866</td>
<td>1212</td>
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<td>O$_3$ (ppbv)</td>
<td>34.2</td>
<td>35.7</td>
<td>31.9</td>
<td>21.1</td>
<td>24.6</td>
<td>22.6</td>
</tr>
</tbody>
</table>

$^a$ At 1 m above the snow surface.
$^b$ Based on concentrations at 1.0 and 0.01 m above the snow surface.
$^c$ Model estimates.
$^d$ From daily sampling of the top 0.5 cm of snow.
Figure 1. Meteorology and NO$_x$ observations at Dome C in summer 2011–2012 (highlighted periods I.–IV. as referred to in text and Table 2): (a) air temperature ($T$) at 1.6 m and modeled mixing height ($h_z$) (Gallée et al., 2015), (b) wind speed (wspd) and direction (wdir) at 3.3 m (c), 1 min averages of NO$_x$ mixing ratios at 1 m (red line is 1 day running mean) and (d) 10 min averages of observational estimates of NO$_x$ flux ($F_{NO_x}$) between 0.01 and 1 m (red line is 14 day running mean).
Figure 2. Balloon profiles (vertical dashed lines) from 9 January 2012: (a) modelled mixing height $h_z$ (10 min running mean) and observed turbulent diffusion coefficient of heat $K_h$ at 1 m (symbols: 10 min averages and black line: 30 min running mean) at 1. (b) interpolated vertical profiles of NO$_x$ mixing ratios with contour lines representing 60 pptv intervals. The lower 100 m appear well mixed during the day, while after collapse of the convective boundary layer in the early evening snow emissions of NO$_x$ are trapped near the surface causing a strong increase in mixing ratios near the ground.
Figure 3. Firn air mixing ratios of (a) NO\textsubscript{x} and (b) O\textsubscript{3}, observed on 12 January 2012. Symbols represent 30 min averages. Solid and dashed lines are results from a 20 m and 50 m long intake lines, respectively. Shown are also NO\textsubscript{3} concentrations in snow at 100 m (P1) and 5 km (P2) distance from the lab shelter as well as from under the firn probe (P3).
Figure 4. The impact of rapid changes in incident solar radiation on atmospheric NO\textsubscript{x} mixing ratios (1 min values). (a–b) ambient concentrations at 1 m during a partial solar eclipse on 25 November 2011 (shaded area) with black lines representing the 10 min running mean. (c–d) firn air concentrations at 10 cm depth during a shading experiment using UV-filters on 11 January 2012. Square symbols and error bars represent interval averages and SD, respectively. Shaded areas and filled squares indicate time periods when the UV filter was in place.
Figure 5. Median daily values of MAX-DOAS BrO vertical amounts from Dome C during sunny days or part-days only, after subtracting zenith amounts (see text). Reference spectrum from near-noon on 18 December until 6 January, then from near noon on 7 January. The apparently larger vertical amounts at higher elevations show that much of the BrO is in the free troposphere.
Figure 6. Observed median diurnal cycles during selected intervals in 2011–2012 (a–e) and 2009–2010 (f–i). Shown are (a, f) NO$_x$ mixing ratios at 1 m (a, f), (b, g) NO$_x$ flux (F-NO$_x$) between 0.01 and 1 m (b, g), (c, h) the turbulent diffusion coefficient of heat ($K_h$) at 1 m (c, h), (d, i) the difference in NO$_x$ mixing ratios ($\Delta$NO$_x$) between 1.0 and 0.01 m (d, i), and (e) the $2\pi$ downwelling nitrate photolysis rate coefficient ($J_{NO_3^-}$) (e). Note comparable observations of $J_{NO_3^-}$ are not available from 2009–2010.
Figure 7. (a) Total column O$_3$ above Dome C. (b) NO$_3^-$ concentrations in the skin layer of surface snow (top 0.5 cm). (c) Observational estimates of NO$_x$ flux ($F_{NO_x}$) between 0.01 and 1 m (10 min averages) and modelled $F_{NO_2}$ (multiplied by 10) based on NO$_3^-$ in the skin layer and depth profiles observed at 100 m (P1) and 5 km (P2) distance from the lab shelter (see Fig. 3a); the 1 day running mean of $F_{NOx}$ during 2009–2010 is shown for comparison (from Frey et al., 2013) (d) Atmospheric O$_3$ mixing ratios. Highlighted periods I.–IV. as referred to in text and Table 2.