Estimate of surface direct radiative forcing of desert dust from atmospheric modulation of the aerosol optical depth

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Abstract

Measurements carried out on the island of Lampedusa, in the central Mediterranean, on 7 September 2005, show the occurrence of a quasi periodic oscillation of aerosol optical depth, column water vapour, and surface irradiance in different spectral bands. The oscillation has a period of about 13 min and is attributed to the propagation of a gravity wave able to modify the vertical structure of the planetary boundary layer. The wave occurred during an event of Saharan dust at Lampedusa. The oscillation amplitude is about 0.1 for the aerosol optical depth, and about 0.4 cm for the column water vapour. The modulation of the downward surface irradiances is in opposition of phase with respect to aerosol optical depth and water vapour column variations. The perturbation to the downward irradiance produced by the aerosols is determined by comparing the measured irradiances with estimated irradiances at a fixed value of the aerosol optical depth, and by correcting for the effect of the water vapour in the shortwave spectral range. The direct radiative forcing efficiency, i.e. the radiative perturbation to the net surface irradiance produced by a unit optical depth aerosol layer, is determined at different solar zenith angles as the slope of the irradiance perturbation versus the aerosol optical depth. The estimated direct surface forcing efficiency at 60° solar zenith angle is \(- (181 \pm 17) \text{Wm}^{-2}\) in the shortwave, and \(- (83 \pm 7) \text{Wm}^{-2}\) in the photosynthetic spectral range. The estimated daily average forcing efficiencies are of about \(-79\) and \(-46 \text{Wm}^{-2}\) for the shortwave and photosynthetic spectral range, respectively.

1 Introduction

Relatively large uncertainties exist in the determination of the role of atmospheric aerosols on the radiative budget, and on the direct forcing produced by aerosols at the Earth’s surface. Satheesh and Ramanathan (2000) suggested the combined use of measurements of aerosol optical depth and shortwave radiation net irradiances at the surface to estimate the shortwave direct aerosol forcing. This approach does not
require the use of radiative transfer models for the evaluation of the aerosol-free short-wave irradiance, and of the assumptions required in the modelling exercise.

The so called direct method has been used by several authors to derive the aerosol forcing in different regions and with different aerosol components and conditions (e.g. Jayaraman et al., 1998; Conant, 2000; Podgorny et al., 2000; Satheesh and Ramanathan, 2000; Meywerk and Ramanathan, 2002; Markowicz et al., 2002; Bush and Valero, 2002; di Sarra et al., 2008; Di Biagio et al., 2009, 2010).

A relatively large dataset, covering at least several months of data, is required to derive the forcing with the direct method. The use of data from a relatively wide temporal interval implies that different classes of particles and, within the same aerosol class, particles with different properties are included simultaneously in the analysis.

di Sarra et al. (2008) have applied this method to the central Mediterranean, and derived estimates of the aerosol forcing for different aerosol classes, namely desert dust, polluted, and biomass burning particles. In a more recent analysis these aerosol types are further divided in different sub-classes, depending on the observed value of the single scattering albedo, SSA (Di Biagio et al., 2009, 2010). The obtained results suggest a dependency on both the aerosol type and the SSA.

In this study we derive the aerosol forcing from high temporal resolution measurements of aerosol optical depth and surface radiative fluxes during the occurrence of a hydrodynamic oscillation of the boundary layer, probably related with the propagation of a gravity wave. The oscillation in the aerosol optical depth associated with the wave induces a quasi-periodic change in the downward shortwave irradiance, and provides the opportunity to derive an estimate of the aerosol direct forcing in a short time period. The method has been applied to one event observed at Lampedusa (35.5° N, 12.6° E), in the central Mediterranean, characterized by transport of Saharan dust.
2 Observations

Observations collected at Lampedusa during the morning of 7 September 2005 display a quasi-periodic modulation of the aerosol optical depth, the column water vapour, and the downward shortwave irradiance. The occurrence of such a phenomenon at Lampedusa in daytime and with cloud-free conditions is seldom observed: out of several years of aerosol optical depth (see e.g. Meloni et al., 2007, 2008) only three cases were clearly identified. One was observed on 22 August 2007; however, in this case the variation in optical depth is small (0.02–0.03) and the retrieval of the forcing is unpractical with the proposed method. A third event took place in October 2011 and will be investigated separately.

The aerosol optical depth, Ångström exponent, and column water vapour are derived from multi-filter rotating shadowband radiometer (MFRSR; Harrison et al., 1994) measurements. The MFRSR measures downward global and diffuse irradiances in seven channels. Six channels have a full width half maximum bandwidth of about 10 nm, and are centred at 416, 500, 617, 671, 869, and 940 nm, respectively. The seventh is a broadband channel, which detects radiation between 300 and 1100 nm. In this study we use aerosol optical depths at 500 and 869 nm. The aerosol optical depth is derived from the direct irradiance, calculated as the difference between global and diffuse irradiances, by applying the Bouguer law. The Ångström exponent is calculated from the aerosol optical depths at 500 and 869 nm. Details on the retrieval method and associated errors are given by Pace et al. (2006). The column water vapour is derived from the MFRSR direct irradiance at 940 nm. The signal at 940 nm is calibrated with the modified Langley plot method against integrated water vapour data from radiosondes, and the column water vapour is derived from the measured transmissivity in this channel, as described by Liberti et al. (2010).

Downward shortwave (SW) irradiance is measured with an Eppley Precision Spectral Pyranometer (PSP), while downward photosynthetic active radiation (PAR) is obtained with a Licor LI-190 sensor. The evolution of shortwave and photosynthetic active
irradiances during the morning of 7 September 2005 is shown in Fig. 1. PSP data are corrected for the thermal offset and cosine response as described by Di Biagio et al. (2009). The PSP data are referred to a second PSP calibrated at the PMOD/World Radiation Centre at Davos (Swiss) in September 2005. The PAR radiometer was calibrated at the factory in 2002; its calibration was updated in April 2007 against a freshly calibrated PAR. In this period the calibration constant changed by about 6%, and the adopted value was obtained by temporal interpolation at the time of the measurements.

All irradiance measurements display a quasi periodic variation superimposed on the overall increasing trend occurring in the morning. A periodic change in the diffuse irradiance is also recorded by a TSI-440 total sky imager (TSI), also operational at Lampedusa (not shown). The TSI images do not show presence of clouds, except for a possible very thin cloud close to the Sun during the period 07:19–08:10 UT. The influence of the cloud on the downward irradiance is evident in Fig. 1, and data from this time interval were excluded from the analysis.

The ratio between PAR and SW downward irradiance varies between 0.41 and 0.43; these value are in agreement with measurements of this ratio at other Mediterranean sites (Jacovides et al., 2007).

Figure 2 displays the evolution of the aerosol optical depth at 500 nm, \( \tau \), and of the column water vapour, \( cwv \). The aerosol optical depth varies between 0.29 and 0.41, while \( cwv \) varies between 2.4 and 2.8 cm. The oscillations of \( \tau \) and \( cwv \) are in phase. The period of the oscillation is about 13 min; 17 maxima in \( \tau \) are observed between 05:53 and 09:29 UT. In the same time interval the solar zenith angle varies between 77.5° and 37.2°. The maximum amplitude of the aerosol optical depth oscillation is about 0.1 around 07:30 UT, and progressively reduces afterwards. The surface pressure is measured at Lampedusa every 10 min; the relative maxima and minima (although measurements every 10 min under sample the oscillation) are in phase with the maxima of the aerosol optical depth.

The oscillation observed in the aerosol optical depth, in water vapour, and in atmospheric surface pressure is attributed to the presence of a gravity wave propagating.
through the southern Mediterranean. The radiosonde profiles measured at Tunis (36.8° N, 10.2° E) on 7 September 2005 at 00:00 and at 12:00 UT are compatible with the development of gravity waves in the troposphere, with a phase velocity of about 8 m s⁻¹ from the NW direction. The critical level, where the phase velocity of the wave matches the horizontal velocity of the wind and the Richardson number reaches a minimum value of 0.05 is at about 3200 m a.s.l.

We assume that the propagation of a gravity wave produced a modulation of the troposphere, which, in its turn, may have influenced the column amount of atmospheric constituents, such as water vapour and aerosols. In fact in an extremely simplified treatment of a gravity wave propagating in a layer of monodisperse aerosol it can be shown that the perturbation of the optical depth associated to the gravity wave can be written as

\[ \tau' \cong \Re \left\{ e^{i(k_x \cdot x - \omega t)} \beta_{\text{ext}} \int_0^{\infty} \frac{d}{dz'} n_0(z') \left\{ \frac{1}{(\omega - k_x U_0(z))} W(z) dz \right\} \right\} \]  

(1)

where \( n_0(z) \) and \( \beta_{\text{ext}} \) are the background number density and the extinction coefficient of the aerosol, \( U_0(z) \) is the background atmospheric horizontal wind velocity projected in the same direction of the wave phase velocity, \( \omega \) and \( k_x \) the frequency (in principle complex) and horizontal wavenumber of the wave, respectively, and \( W(z) \) is the complex amplitude of the vertical velocity perturbation. It is easily verified that the optical depth will oscillate in time, with a possible growth/damping rate given by the imaginary part of \( \omega \), only if the integral in Eq. (1) is not null. Such condition depends critically on the vertical structure of the wave and on the aerosol background profile. It must be noticed that the derivative of the background aerosol number density acts as a weighting function which can be assumed monotonic and approaching fast to zero above the boundary layer; thus the integration interval can be reduced to a few kilometres, possibly below the critical level, with a negligible error. With such an assumption and with the additional reasonable assumption that the imaginary part of \( \omega \) is small (but not so small
to introduce singularities), the optical depth will oscillate if the vertical integrals up to
the critical level, $z_c$, of the real and imaginary parts of the vertical velocity perturbations
$W_r$ and $W_i$ do not both vanish.

This condition is usually verified in stability analyses present in the literature, where
the profiles of the vertical velocity disturbance roughly display a semi oscillation be-
tween the ground where they are both zero due to boundary conditions (first node)
and the critical level (occasionally close to a second node); see for example Mastran-	onio et al. (1976) and Fua and Einaudi (1984). A more detailed study is needed and
will be the subject of a future paper. Here we may state that it is highly probable that
the observed oscillations in the aerosol optical depth and column water vapour are
indeed produced by a gravity wave propagating through the atmosphere at that time.
Observations of column water vapour oscillations, with development of clouds, associ-
ated with the propagation of gravity waves have been previously reported by Reinking
et al. (2000). The influence of gravity waves on vertically integrated quantities for a non-
rotating model of the tropical atmosphere is studied by Raymond and Fuchs (2007).

The backward airmass trajectories reaching Lampedusa on 7 September 2005 over-
pass Northern Libya (trajectories arriving below 2000 m), Tunisia and Northern Algeria
(trajectories arriving above 2000 m), carrying Saharan dust particles to Lampedusa.
The trajectories are calculated at 08:00 UT, using the Hybrid particle dispersion model
(HYSPLIT, Draxler and Rolph, 2012). The aerosol optical properties measured dur-
during the morning of 7 September, 2005, are typical of cases dominated by Saharan
dust (e.g. Pace et al., 2006; Meloni et al., 2006, 2007): the Ångström exponent of the
aerosol during the morning is comprised between 0.5 and 0.6. The change in Ångström
exponent is small and not correlated with the oscillation of the optical depth. Thus, we
assume that the changes in aerosol microphysical properties are small throughout the
morning and changes in optical depth are primarily due to changes in aerosol number
density.
As will be discussed in detail below, the variations in downward irradiance shown in Fig. 1 are in opposition of phase to changes in aerosol optical depth and water vapour column amount, suggesting a negative aerosol shortwave forcing at the ground.

The combined modulation of downward irradiance, water vapour column, and aerosol optical depth is used to estimate the aerosol effect on the irradiance. From the irradiance-optical depth relation, the aerosol surface direct forcing efficiency is derived, under the assumption that during the measurement interval the aerosol microphysical properties do not appreciably change. The water vapour column amount is used to take into account and correct for the influence of the water vapour changes on the shortwave irradiance. The derivation of the forcing efficiency is described in detail in the next section.

3 Aerosol forcing

Using the direct method, the direct aerosol forcing efficiency can be derived uniquely from observations by plotting the net irradiance versus the aerosol optical depth, and calculating the slope of the fitting linear relationship. The slope is equal to the forcing produced by an aerosol layer whose optical depth is equal to 1.

In our case, the determination of the aerosol forcing is based on the estimate of the variation in the downward irradiance produced by a variation in aerosol optical depth. The downward irradiance depends largely on the solar zenith angle, and consequently, on time. In order to identify the aerosol effect, a reference downward irradiance which includes the effect of changes in solar zenith angle, but not aerosol optical depth and water vapour column, needs to be defined.
3.1 Water vapour column

All PSP measurements are scaled to the column water vapour amount of 2.6 cm, using the expression:

\[ F(\tau, \text{cwv} = 2.6 \text{ cm}, \theta) = F(\tau, \text{cwv}, \theta)[1 + f(\theta)(2.6 - \text{cwv})] \]

where \( f \) gives the relative change in downward irradiance per change in \( \text{cwv} \), and is estimated using the SBDART model (Ricchiazzi et al., 1998) for different solar zenith angles, \( \theta \), and water vapour amounts. The value of \( f \) ranges between 0.02 (at low solar zenith angle) and 0.03 (at 75° solar zenith angle) cm\(^{-1}\). The value of \((2.6 - \text{cwv})\) is comprised between −0.2 and 0.2 cm in the investigation period, and the correction to the downward irradiances due to the water vapour changes are always smaller than 0.6 % (less than 3.5 Wm\(^{-2}\) at 35° solar zenith angle, and less than 1.4 Wm\(^{-2}\) at 75°).

3.2 Aerosol optical depth

The reference irradiance \( R(\theta) \) is derived by fitting the water vapour corrected PSP irradiances taken at a fixed value of the optical depth with the following expression:

\[ R(\theta, \tau_f) = a_1 + a_2 \cos(\theta) + a_3[\cos(\theta)]^2. \]

Two values of \( \tau_f \) were chosen in two different time intervals: \( \tau_f = 0.315 \) in the time interval 05:50–08:59 UT (having excluded the interval 07:09–08:23), and \( \tau_f = 0.295 \) in the interval 08:39–09:57. The \( a_1, a_2, \) and \( a_3 \) coefficients are derived for each of the two intervals, allowing to derive the downward irradiances at fixed water vapour and aerosol optical depth as a function of the solar zenith angle or time. The two curves curve fit well the experimental data in both intervals.

The difference between the measured downward irradiance and \( R(\theta) \) (fit for \( \tau_f = 0.315 \)) is also displayed in Fig. 2. It must be noted that this difference includes the effects of both aerosol and water vapour, which are included in the measurements.
while $R(\theta)$ is calculated at fixed cwv and $\tau$. The differences in downward irradiance are evidently linked and in opposition of phase with respect to the aerosol and water vapour variations.

### 3.3 Forcing estimate

The deviations of $F(\tau, \text{cwv} = 2.6 \text{cm})$, $\Delta F$, from the fitting curves (providing the downward irradiances at fixed water vapour and aerosol optical depth) are due to the influence of the aerosol variations from the reference value.

The values of $\Delta F$ are averaged over 1 min intervals and are plotted against the 1-min average of the aerosol optical depth in different solar zenith angle intervals, with the aim of taking into account the dependence of the forcing efficiency on the solar position. Solar zenith angle interval of five degrees around 75°, 70°, 65°, 45°, 40°, and 35° solar zenith angles, were adopted. Due primarily to the limited variability in aerosol optical depth, a poor linear correlation between $\Delta F$ and $\tau$ is found for the intervals centred at 75°, 45°, 40°, and 35°. Reliable results could be conversely obtained by aggregating the two intervals centred at 35° and 45°, while the data around 75° and 45° were not used for the fit.

Figure 3 shows the behaviour of the SW radiation residuals reduced to a cwv of 2.6 cm versus the observed aerosol optical depth, separately for the three solar zenith angle intervals. The linear fits to the data, whose slope is $b$, are also displayed.

The slope of the fit $b$ is used to derive the forcing efficiency $FE$ with the expression:

$$FE = b(1 - a)$$

where the term $(1 - a)$ is needed to derive the surface net irradiance from the downward irradiances. The value of the albedo $a$ is calculated following Jin et al. (2004) taking into account the measured wind speed, and is a function of the solar zenith angle.

The accuracy of the determination of the forcing with this method depends primarily on the differences among irradiances and optical depths measured with the same instrument at short time intervals. Although all the instruments are regularly calibrated...
and controlled, the overall calibration of the sensors has a minor effect on the results. In addition, since the reference curve at fixed aerosol optical depth is determined from the same observations, the influence of possible instrumental effects, such as the residual influence of the angular response of the radiometers, are minimized. The instrumental stability within the measurement time interval is very good, and differences in irradiances and optical depths, although small, are significant. Consequently, the uncertainty on FE is estimated taking into account only the uncertainty on the fit, calculated following Higbie (1991), and the estimated error on the correction for the water vapour changes. The influence of this correction may be relevant, due to the relatively large uncertainty on the water vapour column measurement, of the order of 15%. The uncertainty on the slope is thus derived by taking into account the effect of a 15% variation in cwv, which contributes to the total uncertainty by about 12 W m\(^{-2}\), and the uncertainty associated with the fit.

The same procedure, except for the water vapour correction which is not required in the visible spectral range, was applied to the PAR measurements. Better linear fits are generally obtained from the PAR signals, and estimates of FE could be derived also for the solar zenith angle interval around 45°. Figure 4 shows the obtained fits.

### 4 Results and discussion

Table 1 reports the values of the FE for the SW and the photosynthetically active spectral ranges, for the selected solar zenith angle intervals.

The average ratio between PAR and SW FEs is 0.55, with an estimated uncertainty of the order of 0.1. The lowest ratio is found for large solar zenith angles, while the highest for about 35°. At small solar zenith angle the ratio between the PAR and SW forcing efficiencies is larger than the ratio between PAR and SW irradiance. This effect is attributed to three factors: the aerosol optical depth decreases with wavelength, and also its radiative effect; the presence of absorption bands, mainly by water vapour, at the longer wavelengths in the shortwave range leads to a reduced aerosol influence.
on the downward irradiance; the dust single scattering albedo is relatively low in the visible, and increases at longer wavelengths (e.g. Sokolik and Toon, 1996), producing a stronger aerosol impact on PAR than at longer wavelengths. The aerosol effect in the SW and PAR spectral ranges were also derived using radiative transfer model simulations carried out with the MODTRAN 4 (Berk et al., 2003) for the model desert dust properties and at the same solar zenith angles. The average ratio between the dust forcing in the PAR and SW spectral ranges is 0.62, which is higher than the modelled PAR to SW irradiance ratio (0.44). As expected, the ratio depends on the dust optical properties.

The reader is reminded that the FEs reported in Table 1 are instantaneous values. A polynomial curve fitted to the values of Table 1 is calculated at different times and averaged over a whole day to derive an estimate of the daily average FE. The derived estimates of the SW daily average FE is $-79 \text{ Wm}^{-2}$, and the PAR daily average FE is $-46 \text{ Wm}^{-2}$. Over the whole daily cycle the ratio between PAR and SW FE is 0.58.

The daily SW forcing for the selected day, with an average aerosol optical depth of about 0.3, is about $-24 \text{ Wm}^{-2}$ in the SW, and is about $-14 \text{ Wm}^{-2}$ for PAR.

Estimates of the Saharan dust direct shortwave forcing efficiency over the Mediterranean sea by different authors are listed by Di Biagio et al. (2010). The reported daily average values span from $-50$ and $-80 \text{ Wm}^{-2}$ (Zhou et al., 2005). Previous studies at Lampedusa suggested values around $-68 \text{ Wm}^{-2}$ (di Sarra et al., 2008; Di Biagio et al., 2009, 2010). More recently, di Sarra et al. (2011) estimated an FE of about $-55 \text{ Wm}^{-2}$ for a very intense dust event.

The determinations of the instantaneous value of the SW FE are also within the range of values found in the literature. Previous determinations obtained at Lampedusa give values between $-140$ and $-220 \text{ Wm}^{-2}$ for $35^\circ$ solar zenith angle (di Sarra et al., 2008, 2011; Di Biagio et al., 2009, 2010). In a recent modelling study Gómez-Amo et al. (2011) derive values between $-175$ and $-250 \text{ Wm}^{-2}$ at $60^\circ$ solar zenith angle.

Meloni et al. (2005) estimated the radiative effect produced in the visible spectral range by two Saharan dust events occurred at Lampedusa in July 2002. They found
a daily average surface FE of about $-30 \text{ Wm}^{-2}$ for a case characterized by high aerosol single scattering albedo, and of about $-43 \text{ Wm}^{-2}$ for a case with low aerosol SSA.

## 5 Conclusions

Observations collected at Lampedusa, in the central Mediterranean, during an oscillation of the troposphere associated with the propagation of a gravity wave were used to derive an estimate of the Saharan dust radiative forcing. The oscillation occurred during a Saharan dust event on 7 September 2005, and is observed in the aerosol optical depth, column water vapour, and downward shortwave and PAR irradiances.

The oscillation of the downward irradiances are in opposition of phase with respect to changes in aerosol optical depth and water vapour column, and are attributed to the combined effect of these two factors. All SW irradiances are reported to the same cwv content, by using radiative transfer model calculations. Reference curves describing the downward SW and PAR irradiances are constructed by using measurements obtained at a fixed value of the aerosol optical depth. The radiative perturbation produced by dust is thus obtained as the difference between the measured irradiances, reported at cwv = 2.6 cm, and the reference curve at the same solar zenith angle.

The radiative forcing efficiencies are then determined at some values of the solar zenith angle as the slope of the linear fit of the radiative perturbation versus the measured aerosol optical depth.

The estimated value of the dust FE is $-188 \text{ Wm}^{-2}$ for the SW and $-93 \text{ Wm}^{-2}$ for PAR at 70° solar zenith angle, $-163 \text{ Wm}^{-2}$ for the SW and $-112 \text{ Wm}^{-2}$ for PAR at 45°. The estimated daily average FE is $-79 \text{ Wm}^{-2}$ for the SW range, and $-46 \text{ Wm}^{-2}$ for PAR. The obtained values are in good agreement with previous estimates for desert dust over the ocean. The ratio between PAR and SW FEs is higher than the ratio of the PAR to SW irradiance. This effect, which is confirmed by radiative transfer model calculations, is attributed to spectral changes of the dust optical properties, and the occurrence of strong water vapour absorption bands at the longer wavelength range.
The proposed method allows to estimate the radiative effect of the aerosols, in the presence of a natural oscillation of the atmosphere, by using a limited dataset; the method is robust with respect to a detailed instrumental characterization, since is based mainly on direct observations from the same set of instruments. Although these oscillatory events are infrequent at Lampedusa, they may occur more frequently in regions with a different geographical setting; in these cases the proposed method might allow a better insight on the aerosol radiative effects in different conditions.

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References


**Table 1.** Estimated dust forcing efficiencies derived in the shortwave and PAR spectral ranges and in different time/solar zenith angle intervals.

<table>
<thead>
<tr>
<th>Time interval</th>
<th>Solar zenith angle interval</th>
<th>SW FE (Wm$^{-2}$)</th>
<th>PAR FE (Wm$^{-2}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>06:14–06:45 UT</td>
<td>73°–67°</td>
<td>−188 ± 18</td>
<td>−93 ± 7</td>
</tr>
<tr>
<td>06:40–07:10 UT</td>
<td>67°–62°</td>
<td>−181 ± 17</td>
<td>−83 ± 7</td>
</tr>
<tr>
<td>08:23–08:58 UT</td>
<td>47°–42°</td>
<td>−</td>
<td>−107 ± 5</td>
</tr>
<tr>
<td>08:51–10:11 UT</td>
<td>43°–32°</td>
<td>−163 ± 16</td>
<td>−112 ± 3</td>
</tr>
</tbody>
</table>
Fig. 1. Evolution of the downward shortwave and PAR irradiances measured at Lampedusa during the morning of 7 September 2005.
Fig. 2. Evolution of aerosol optical depth at 500 nm, column water vapour, and of the difference between the measured downward shortwave irradiance and the irradiance estimated at fixed aerosol optical depth and column water vapour (see text) during the morning of 5 September 2007. The region between the two vertical lines corresponds to the presence of thin clouds in the sky imager pictures.
Fig. 3. Deviations of the measured downward shortwave irradiance at fixed column water vapour from the estimated irradiance at fixed water vapour and aerosol optical depth versus the measured aerosol optical depth, around (a) 70°, (b) 65°, and (c), 37.5° solar zenith angle.
Fig. 4. Deviations of the measured downward PAR irradiance from the estimated PAR irradiance at fixed aerosol optical depth versus the measured aerosol optical depth, around (a) 70°, (b) 65°, (c) 45°, and (d) 37.5° solar zenith angle.