Forcing of stratospheric chemistry and dynamics during the Dalton Minimum

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Abstract

The response of atmospheric chemistry and climate to volcanic eruptions and a decrease in solar activity during the Dalton Minimum is investigated with the fully coupled atmosphere-ocean-chemistry general circulation model SOCOL-MPIOM covering the time period 1780 to 1840 AD. We carried out several sensitivity ensemble experiments to separate the effects of (i) reduced solar ultra-violet (UV) irradiance, (ii) reduced solar visible and near infrared irradiance, (iii) enhanced galactic cosmic ray intensity as well as less intensive solar energetic proton events and auroral electron precipitation, and (iv) volcanic aerosols. The introduced changes of UV irradiance and volcanic aerosols significantly influence stratospheric climate in the early 19th century, whereas changes in the visible part of the spectrum and energetic particles have smaller effects. A reduction of UV irradiance by 15% causes global ozone decrease below the stratopause reaching 8% in the midlatitudes at 5 hPa and a significant stratospheric cooling of up to 2°C in the midstratosphere and to 6°C in the lower mesosphere. Changes in energetic particle precipitation lead only to minor changes in the yearly averaged temperature fields in the stratosphere. Volcanic aerosols heat the tropical lower stratosphere allowing more water vapor to enter the tropical stratosphere, which, via HOx reactions, decreases upper stratospheric and mesospheric ozone by roughly 4%. Conversely, heterogeneous chemistry on aerosols reduces stratospheric NOx leading to a 12% ozone increase in the tropics, whereas a decrease in ozone of up to 5% is found over Antarctica in boreal winter. The linear superposition of the different contributions is not equivalent to the response obtained in a simulation when all forcing factors are applied during the DM – this effect is especially well visible for NOx/NOy. Thus, this study highlights the non-linear behavior of the coupled chemistry-climate system. Finally, we conclude that especially UV and volcanic eruptions dominate the changes in the ozone, temperature and dynamics while the NOx field is dominated by the EPP. Visible radiation changes have only very minor effects on both stratospheric dynamics and chemistry.
1 Introduction

The fourth assessment report of the Intergovernmental Panel on Climate Change (Forster et al., 2007) noted that while the scientific understanding of the greenhouse gas (GHG) emissions and volcanic effects on climate is rather high, this is not the case for changes in solar activity. The combined forcings of GHG and tropospheric aerosols is predicted to increase until possible stabilization is reached in the second half of the 21st century. The volcanic effect is unpredictable. Concerning solar activity, it is hypothesized that solar activity will – after a long period of high activity – drop to a new grand minimum in the 21st century (Abreu et al., 2008, 2010). Given this, an assessment of periods in the past with a similar decrease of the solar activity is helpful to understand the mechanism and its implications. As an example, the Dalton Minimum (DM) was a time period lasting from 1790 to 1830 which was characterized by a significant cooling in Europe (Luterbacher et al., 2004) and the extratropical Northern Hemisphere (Ljungqvist, 2010; Auchmann et al., 2012). This unusually cold time coincides with the period of very low solar activity as expressed in low sunspot numbers (Hoyt and Schatten, 1998) and high volcanic activity due to two major volcanic eruptions in 1809 and in 1815. The exact causes of this cooling are not well defined. Some part of it can be explained by downward propagating stratospheric perturbations (e.g. Ineson et al., 2011). We thus decided to study this period and address the solar and volcanic effects on stratospheric climate and chemistry. Up to now, studies of the DM where done to a major part with climate models with coupled interactive oceans. The novelty of our experiment setting was to include interactive chemistry to a GCM coupled to a deep layer ocean. We succeeded thus to include the most important natural forcing in a climate model simulation during the DM: (a) solar irradiance changes, which can be decomposed into a ultraviolet (UV), visible and infrared (IR) part of the spectrum, (b) explosive tropical volcanic eruptions and (c) energetic particle precipitation (EPP).

Solar activity has been monitored for a long time (Wolf, 1861; Hoyt and Schatten, 1998). The influence of the Sun on time scales of up to hundreds of years can first be
divided in two temporal classes: there is a regular, well established 11 yr cycle (Wolf, 1861; Schwabe, 1844) – which can vary in its intensity – and on the longer time scale there are grand minimum and maximum states of the solar activity.

Solar influence can be further classified in terms of where the biggest effects can be observed in the Earth’s atmosphere. This is strongly linked to the part of the spectrum with the largest variability. Kodera and Kuroda (2002) investigated the effects of the 11 yr solar cycle on atmospheric dynamics, focusing on the UV part of the spectrum. Their work suggests a downward propagation of the response in the middle atmosphere caused by heating through UV absorption and ozone increase. In solar active conditions, this additional heating leads to an increasing pole-to-equator temperature gradient, influencing also the stratospheric zonal winds (Kodera and Kuroda, 2002). This process is known as the top-down-mechanism (e.g. Meehl et al., 2009; Gray et al., 2010). A different aspect is to focus rather on the visible spectrum and to follow a bottom-up-approach (Meehl et al., 2009): during solar active conditions more evaporation occurs in the subtropics. This in turn leads to an increase in the precipitation amount, which accelerates the Hadley and Walker cells (Labitzke et al., 2002), finally leading to ENSO-like anomalies and influencing stratospheric circulation. Using reconstructions of the solar irradiance like the ones from Lean et al. (1995), these two processes have been studied extensively in the last years using models of different complexity (see Gray et al., 2010 and references therein).

The fact that volcanoes can influence global climate has already been recognized in Franklin (1784) and Milham (1924). While Franklin (1784) mainly focused on the effect on the troposphere, which, after the Lakagígar eruption, was polluted by a large amount of particles, partly leading to a constant haze, Milham (1924) focused on the Tambora eruption and the uncommon weather pattern following the eruption. In the twentieth century, partly because of four major volcanic eruptions (Agung in 1963, Fuego in 1974, El Chichón in 1982, Pinatubo in 1991), more intense scientific research was done, focusing especially on the radiative effects of stratospheric aerosols (see Hansen et al., 1992; Stenchikov et al., 1998; Robock, 2000, and references therein). The plume of
powerful volcanic eruptions reaches the stratosphere (Halmer et al., 2002). There, \( \text{SO}_2 \) is transformed through a number of chemical reactions to sulfate aerosols. Aerosols at lower stratospheric altitudes (Whitten et al., 1980) are mostly spherical (Tratt and Menzies, 1994) and scatter back to space part of the incoming solar short-wave radiation. On the other hand, sulfate aerosol absorb thermal radiation. The aerosol particles also provide a media for heterogeneous reactions facilitating the removal of reactive nitrogen oxides and activation of halogen radicals. Thus, volcanic aerosols are important for both radiative and chemical processes in the atmosphere.

Reconstructions of the volcanic forcing (Gao et al., 2008) have been used to model past and modern influences of volcanic events on global climate. Generally, following observations and modeling studies, while the lower stratosphere is heated by absorption of infrared radiation by the aerosols, the troposphere and the surface usually experiences a significant cooling after major volcanic eruptions (Dutton and Christy, 1992; Minnis et al., 1993; Stenchikov et al., 1998; Arfeuille, 2012). The interaction with chemistry is more complex due to the effects of enhanced halogen loading in modern times (Tie and Brasseur, 1995). In a clean preindustrial atmosphere, a significant globally averaged increase of total column ozone can be expected within one to three years after a volcanic eruption, whereas at the equator, ozone column depth is decreasing slightly (Arfeuille, 2012). In a halogens contaminated atmosphere of today, global ozone concentration drops significantly after a volcanic eruption. The resulting heating leads to major changes in the atmospheric dynamics and large-scale oscillation patterns like El Niño, Arctic Oscillation (AO) or North Atlantic Oscillation (NAO) (Robock, 2000; Stenchikov et al., 2002; Yoshimori et al., 2005; Wagner and Zorita, 2005; Christiansen, 2007; Fischer et al., 2007; Spangehl et al., 2010).

The influence of EPP on climate is – compared to the other two aforementioned factors – a rather new subject to science and has been investigated increasingly often during the last twenty years. Its effect is still not well known and a quite controversial issue in the climate change discussion. This disagreement is also a reason why EPP have not been included in important climate model simulation campaigns in support
of WMO and IPCC assessments (WMO, 2010; Forster et al., 2007). The EPP can be divided into three main categories: galactic cosmic rays (GCRs), solar protons and high- and low energetic electrons (HEE, LEE). All of them can ionize neutral molecules in the Earth’s atmosphere.

The GCRs are highly energetic charged particles. They originate from supernova explosions in our galaxy and their flux and energy spectrum at the entry of the heliosphere is very stable, only being modulated by the solar activity, which is shielding the Earth from them via magnetospheric deflection. The observed GCR flux variability thus follows the cycles of the solar magnetic activity. The energy spectrum of GCRs ranges from 1 to $10^{11}$ MeV (Grieder, 2001; Chang et al., 2008). After entering the atmosphere GCRs dissipate their energy mainly by ionization processes. Following a Bragg-peak (Bragg and Kleeman, 1905), the maximal ionization rate by GCR is reached between 15 km and 20 km altitude (Usoskin et al., 2010). The ionization is largest in the polar regions (poleward of $\pm 60^\circ$) where the geomagnetic field has the weakest shielding effect (lowest cut-off rigidity).

Solar protons events (SPEs) – emerge from coronal mass ejections of the sun, which occur very irregularly and are not always directed towards the Earth. Hence, SPEs are very sporadic and of high unpredictability. The solar wind plasma usually reach the Earth’s atmosphere within 1–2 days after the ejection (Kahler, 1992). The charged particles are directed towards the poles, where they follow the lines of the geomagnetic field into the atmosphere. Only in extreme cases – when their energies reach 500 MeV or more – they can propagate down to the stratosphere (Jackman et al., 2008). As a result of the magnetic shielding, the effect of SPEs is strongly latitude dependent with a minimum equatorward of $\pm 20^\circ$ and a maximum poleward of $\pm 60^\circ$.

LEEs and HEEs originate from the interaction of the Earth’s magnetospheric plasmasheet with the solar wind (Brasseur and Solomon, 2005). Solar plasma is kept trapped in the magnetosphere of the Earth and can be accelerated during periods of higher solar wind speeds. The accelerated electrons then rapidly travel along the magnetic field lines to the poles and partly penetrate the uppermost layers of the at-
mosphere (Bazilevskaya et al., 2008). The best evidence for their existence are the aurorae, formed by the excitation of nitrogen and oxygen atoms.

Ionization of oxygen and nitrogen lead to NO$_x$ and HO$_x$ production. While HO$_x$ has a short life-time in the range of minutes to hours and thus affects atmospheric chemistry only locally, NO$_x$ can survive for some days in the stratosphere and will thus be distributed globally by e.g. the Brewer–Dobson circulation (BDC). In the stratosphere, NO$_x$ and HO$_x$ interact with ozone in a significant way, as was found by analyzing important ionization events (Callis et al., 1998; Funke et al., 2011). Changes in ozone concentration inside the polar vortex modify the pole-to-equator temperature gradient and thus can have a significant influence on circulation and weather patterns (Gray et al., 2010). Different modeling studies demonstrated the influence of EPP not only on chemistry (Jackman et al., 2008; Egorova et al., 2011; Calisto et al., 2011; Rozanov et al., 2012) but also on climate (Calisto et al., 2011; Rozanov et al., 2012).

Climate during the DM minimum has already been simulated with general circulation models (GCM) in a number of studies (Bauer et al., 2003; Wagner and Zorita, 2005; Ammann et al., 2007; Spangehl et al., 2007; Arfeuille, 2012). While Bauer et al. (2003) only used a simplified model, Wagner and Zorita (2005) and Spangehl et al. (2007) exploited a coupled atmosphere–ocean GCM (AO-GCM). Arfeuille (2012) used the chemistry-climate model (CCM) SOCOL to simulate the effects of the Tambora volcanic eruption in 1815 on climate and found a strong geopotential height gradient anomaly (around 250 gpm) between 55° N and 75° N at 50 hPa in the first winter after the eruption (November–April) as well as a net radiative forcing anomaly reaching $-8 \text{Wm}^{-2}$ (60° S–60° N) during the first five months following the eruption. Thus, volcanic and solar influences are generally accepted as main drivers for global climate cooling. Wagner and Zorita (2005) also investigated the contribution of the slightly increasing GHG concentrations during the DM and did not find any significant impact.

In this paper, we investigate the effect of different natural factors on global stratospheric climate during the DM with a fully interactive atmosphere–ocean-chemistry climate model (AO-CCM). To the best of our knowledge, no coupled AO-CCM with
EPP parametrization has yet been used for an in-depth analysis of the climate and chemistry state during the DM so far.

The description of the model framework is done in Sect. 2. In Sect. 3 we describe the chemical and dynamical changes in the stratosphere. In the last chapter, we discuss and summarize the findings of this work.

2 Description of the model and experimental setup

2.1 AO-CCM SOCOL3-MPIOM

The AO-CCM SOCOL3-MPIOM emerges from the coupling of the CCM SOCOL3 (Stenke et al., 2013) and the ocean model MPIOM (Marsland et al., 2003) with the OASIS3 coupler (Valcke, 2013). SOCOL3 consists of the chemistry module MEZON (Model for Evaluation of oZONe trends, Rozanov et al., 1999; Egorova et al., 2003; Schraner et al., 2008) which is coupled to the GCM MA-ECHAM5 (Roeckner et al., 2003). Atmospheric temperature fields are passed to MEZON, which computes the tendencies of 41 gas species, taking into account 200 gas phase, 16 heterogeneous, and 35 photolytical reactions. Once computed, the chemical tendencies are handed back to ECHAM5, which then takes care of the transport of species. The simulations were run in T31 spectral resolution, which is equivalent to a grid spacing of around 3.75°. The vertical spacing is irregular, as the model uses hybrid sigma pressure coordinates on 39 levels from 1000 hPa up to 0.01 hPa (80 km). The chemistry scheme is only called every two hours – simultaneously with the radiative scheme – in order to be computationally efficient.

Due to this relatively coarse vertical resolution the Quasi-Biennial-Oscillation (QBO) is not reproduced autonomously by the GCM. To reproduce the QBO, the equatorial zonal wind field is nudged to reconstructed data in the same manner as described in Giorgetta (1996).
The original ECHAM5 radiation code does not properly treat solar spectral irradiation forcing (Forster et al., 2011), therefore extra-heating correction factors (Zhu, 1994) for the Lyman-α line, the Schumann–Runge, Hartley and Huggins bands as well as for the Herzberg continuum were implemented. The radiation code was also modified in such a way that ECHAM5 reads in spectrally resolved solar irradiance in the six ECHAM5 short-wave bands with varying distribution instead of the standard fixed distribution of the varying total solar irradiance into the six bands.

Parametrization of the different EPPs was done identically to Rozanov et al. (2012) and ref. therein, with the only difference that the code has been modified for use in SOCOLv3. Highly energetic electrons (HEE) was not included in the model. To include the magnetic dependency of the ionization by EPP, a temporal, locally changing dipole magnetic field was implemented in the model using geomagnetic proxy data as input.

2.2 Boundary conditions

The model is forced by several boundary conditions described in the following section. The GHG concentrations for the 1780 to 1840 period of carbon dioxide, methane and nitrous oxide are based on the Paleoclimate Modelling Intercomparison Project Phase III (PMIP3) protocol (Etheridge et al., 1996, 1998; MacFarling-Meure, 2004; Ferretti et al., 2005; MacFarling-Meure et al., 2006). Halogen containing species were kept constant to preindustrial levels.

All solar related drivers were based on the solar modulation potential reconstructions produced from 10Be records from ice cores: for the spectral solar irradiance forcing we use the reconstruction of Shapiro et al. (2011). In Fig. 1, the radiative forcing data is plotted for the six bands of ECHAM5 radiation code. The main difference between this reconstruction and the former ones like Lean et al. (1995) or Bard et al. (2000) is the amplitude of the variability. For example, the difference between the maximum and the minimum TSI value during the DM is roughly 6 Wm$^{-2}$ whereas in e.g. Lean et al. (1995), the drop was only by 2 Wm$^{-2}$. For the photolysis rates, look-up tables are used
which have been generated from the spectral solar irradiance (SSI) of Shapiro et al. (2011).

Several different datasets were used for the energetic particles. For the parameterization of NO\textsubscript{x} influx, Baumgaertner et al. (2009) used the A\textsubscript{p} index which can be reconstructed until the year 1932. The A\textsubscript{p} index can be correlated with the A\textsubscript{a} index, which has a longer history but is only based on two stations. It is available from 1868 to present. Based on sunspot numbers, the A\textsubscript{a} and A\textsubscript{p} indexes can be reconstructed via correlation until the year 1600. SPEs were prescribed from an existing SPE data set (provided by Charles Jackman and covering the period 1963–2008, see Jackman et al., 2009). SPEs are very short-lived (in the order of days), thus such events cannot be reconstructed from classical proxies like \textsuperscript{10}Be. Shea et al. (2006) presented a solution to reconstruct big events like the Carrington event from nitrates deposited in ice cores. This method is however very controversial (Wolff et al., 2008; Schrijver et al., 2012). In our work, SPEs are randomized for the years before 1963 by using a return-period based analysis of the last 45 yr, and weighted this result with the A\textsubscript{p} index. Cosmic rays are based on the solar modulation potential (\(\Phi\)), which has been reconstructed by Steinilber et al. (2008). The dataset compares well with the neutron monitor measurements which are available since the year 1950. \(\Phi\) is an index which describes the solar modulation of the cosmic ray flux, which can be converted into pressure-latitude ionization rates using lookup tables from Usoskin et al. (2010). Paleo-magnetic datasets (C. Finlay, personal communication, 2010) are applied to the model in order to take into account the geomagnetic dependency of the ionization.

Stratospheric aerosol properties are prescribed according to the approach applied by Arfeuille et al. (2013) who used the Atmospheric and Environmental Research Inc. (AER) model (Weisenstein et al., 1997) constrained with ice core-based estimates of sulphate aerosol mass (Gao et al., 2008, 2012). From those simulations, zonally averaged aerosol optical properties (spectrally resolved) and surface area density forcing data for for SOCOL3-MPIOM were extracted. The most important volcanoes during the DM period where two weak eruptions (1 and 3 Mt SO\textsubscript{2}) in 1794 and 1796 with unknown
locations (probably extra-tropical), a strong unknown tropical eruption of 27 Mt SO$_2$ in 1809, the 55 Mt SO$_2$ Tambora eruption in 1815 and two eruptions (8.5 Mt and 20 Mt SO$_2$) – the Babuyan Claro in 1831 and the Cosiguina in 1835.

Tropospheric aerosol properties were constructed by scaling the existing CAM3.5 simulations with a bulk aerosol model driven by CCSM3 (CMIP4) sea-surface temperatures and the 1850–2000 CMIP5 emissions (S. Bauer, personal communication, 2011). For data before the year 1850, the applied scaling is a function of the world population except for 10% of the presumed 1990 biomass burning aerosols which were considered natural.

The model was forced by the standard (Hagemann et al., 1999; Hagemann, 2002) land surface datasets provided with the ECHAM5 package.

Finally, equatorial zonal mean zonal winds for nudging the QBO were generated from backward extension of the Brönnimann et al. (2007) reconstructions using an idealized QBO cycle plus a seasonal anomaly cycle.

2.3 Experiments

To investigate the influence of solar, volcanic and the EPP forcings, we perform a series of three-members ensemble sensitivity experiments described in Table 1. We initialize our runs in the year 1780 from a transient simulation starting in 1388 AD. While the first ensemble member is run with unperturbed initial 1780 conditions, the two following members are initialized with an ocean field of the year 1781 and 1779. Every experiment covers 60 yr to reach December 1840. The analyzed period is chosen from 1805 to 1825 in order to reduce the noise and strengthen the signal from volcanic, solar and particle forcings. The small number of ensemble members is chosen to reduce the computational time needed for the simulations.

To address the relative roles of the UV, visible and infrared radiation as well as the extra heating and the photolysis rates, the two experiments called DM-TD (top-down) and DM-BU (bottom-up) are designed. In DM-TD, all forcing data is kept constant except (i) the first radiation band (UV) of ECHAM5 radiation code, (ii) the coefficients of extra-
heating parameterization as well as (iii) the photolysis rates. Hence, all forcing comes from the stratosphere only because the response of the heating rate in the second band of ECHAM5 radiation code (240–440 nm) to the solar variability is very small (Forster et al., 2011). The opposite experiment, DM-BU, is designed in a way that all forcing is kept constant except in the 2–6 bands of the ECHAM5 radiation code. Hence, DM-BU does not include any stratospheric heating or ozone production changes, meaning that all extra radiation is absorbed mostly in the troposphere and by the surface. In turn, the DM-VOLC runs are driven with all forcings kept to constant 1780 conditions except the volcanic forcing. The DM-CTRL1780 runs are performed with perpetual 1780 conditions, whereas the DM-ALL runs are driven with all forcings in transient conditions. To address the effect of energetic particles, we carry out the DM-EPP experiment which is forced only by the parameterizations for GCR, SPE and LEE, while all other forcings are set constant or switched to background aerosol concentrations (volcanic forcing).

2.4 Method of comparison

In the next section we analyze 60 yr long time series of annual zonal mean quantities constructed from the results of three 20 yr long ensemble runs for each experiment. The statistical significances was calculated using the two-tailed Student’s t test using the 5 % significance level, comparing all 60 yr long time series for each experiment. All figures illustrate the relative or abolute deviation of the results of the experiment runs relative to the control run (DM-CTRL1780). On all plots, the yellow line indicates the height of the dynamical WMO tropopause. The nonlinearities are computed in a following way: Differences of DM-TD, DM-BU, DM-VOLC and DM-EPP relative to DM-CTRL1780 are computed and added. This field is then compared to the difference field between DM-ALL and DM-CTRL1780. A positive value in the nonlinearity plot would mean, that the stacked relative differences are greater than the combined from DM-ALL.

If not noted differently, the upper left figure of every panel illustrates the overall effect of all factors (DM-ALL), followed by the effects of UV solar irradiance (DM-TD), volcanic
aerosols (DM-VOLC) and energetic particle precipitation (DM-EPP). As a reduction of the visible and infrared radiation in the DM-BU experiment has small effects on the stratospheric chemistry, these results are not shown in the chemical section.

3 Results

3.1 Atmospheric chemistry

In this section, we focus on four species: ozone, water vapor, HO\textsubscript{x} and NO\textsubscript{x}, as they show the most pronounced response to the considered factors.

3.1.1 Ozone

Figure 2a shows the relative effect of DM-ALL including all factors – namely reduction of solar radiation, volcanic eruptions and EPP – with respect to the control simulation (DM-CTRL1780). Substantial ozone depletion is found almost everywhere reaching its maximum (−8%) in the upper tropical mesosphere and middle stratosphere over the high latitudes. However, the opposite response is simulated in the polar upper mesosphere and tropical upper troposphere/lower stratosphere (UT/LS) regions where the ozone mixing ratio increases by up to 15%.

Figure 2b shows that mainly the effects of the solar UV reduction in DM-TD are responsible for the ozone loss at ozone layer height and for the gain of ozone in the polar upper mesosphere. These ozone changes in the atmosphere can be explained mostly by three factors:

1. The decrease in solar UV irradiance which reduces the ozone production via oxygen photolysis in the stratosphere and NO\textsubscript{2} photolysis in the troposphere,
2. The increase of NO\textsubscript{x} (see Fig. 5, Sect. 3.1.3), which facilitates the intensification of the NO\textsubscript{x} cycle of ozone oxidation (Reactions 1–3),

\[
\begin{align*}
\text{NO} + \text{O}_3 &\rightarrow \text{NO}_2 + \text{O}_2 \quad \text{(R1)} \\
\text{NO}_2 + \text{O} &\rightarrow \text{NO} + \text{O}_2 \quad \text{(R2)} \\
\text{Net:} & \\
\text{O}_3 + \text{O} &\rightarrow 2\text{O}_2 \quad \text{(R3)}
\end{align*}
\]

3. The slight compensation of the abovementioned ozone depletion processes due to the stratospheric cooling caused by reduced solar UV and ozone mixing ratio, slowing down the ozone destruction cycles (see Sect. 3.2.1).

At the poles in lower mesospheric height (60–80 km) a surplus of ozone by up to 20\% is explained by the fact that at these heights, the UV radiation acts like a sink rather than a source of ozone. Thus, with less UV radiation, near the mesopause, ozone destruction is suppressed.

The surplus of ozone at the tropical tropopause can be explained by volcanic effects with the DM-VOLC experiment (see Fig. 2c). The main reason for the ozone increase in the tropical UT/LS after volcanic eruptions is the deactivation of HO\textsubscript{x} and NO\textsubscript{x} via heterogeneous reactions on/in the sulfuric acid particles, formed in the stratosphere from the products of the volcanic eruptions. In the present day atmosphere, ozone depletion was observed after major volcanic eruptions and attributed to catalytic reactions involving reactive halogens. However, a potentially significant background of natural chlorine and bromine existed. The effect of this background, to cause ozone depletion, would have been dominated by nitrogen deactivation on sulphate aerosol. Therefore the heterogeneous conversion of HO\textsubscript{x} and NO\textsubscript{x} to relatively passive reservoir species leads to ozone increase to up to 16\%. The decrease of ozone in the mesosphere is related to the strong increase of the HO\textsubscript{x} species at that height (see Sect. 3.1.2.) leading to an acceleration of the ozone depletion cycle.

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Energetic particles can influence the ozone concentration, as shown in Rozanov et al. (2012). Although our NO\textsubscript{x} field looks very similar to the one from the cited work (see Sect. 3.1.3), the ozone response to EPP in the polar mesosphere is much weaker in our simulations. The main reason for this finding is that background temperatures in the involved regions are different in SOCOLv3 from SOCOLv2. Though, the reaction of Nitrogen with oxygen is highly temperature dependent (Funke et al., 2011). Thus, only a minor part of the signal seen in Fig. 2a can be attributed to EPP: the annual mean ozone anomaly shows an ozone decrease of up to 2\% in the southern extratropics, which is due to the ionization of nitrogen by GCRs. The change is however only significant on a 10\% level. Over the poles, at lower mesosphere height, a significant increase of ozone of up to 2\% is simulated due to the lower ionization rates of both SPEs and LEEs, leading to less NO\textsubscript{x} (see later). Seasonal variations are visible. The biggest effect is modeled in austral spring (SON, not shown), where significant losses of ozone at UT/LS level of up to 4\% are found in the southern polar latitudes.

The nonlinearities in the ozone field are hard to distinguish and show only small areas in the ozone layer with significant changes (not shown).

### 3.1.2 HO\textsubscript{x} and water vapor

In Figs. 3a and 4a the differences in water vapor and HO\textsubscript{x} between DM-ALL and DM-CTRL1780 are illustrated. While water vapor concentration increases dramatically above the tropopause, HO\textsubscript{x} is experiencing an increase in tropical UT/LS and a decrease in the mesosphere and middle tropical stratosphere.

The results of the DM-TD experiment illustrated in Fig. 3b help to attribute the H\textsubscript{2}O increase and HO\textsubscript{x} loss in the mesosphere to the introduced decrease of solar UV irradiance. A strong (by up to 25\%) HO\textsubscript{x} decrease in the mesosphere coinciding with pronounced increase of H\textsubscript{2}O is driven by less intensive water vapour photolysis in Lyman-alpha line and Schumann–Runge bands. When looking at the highest levels of the model atmosphere, one recognizes that HO\textsubscript{x} decreases less in the lower mesosphere than at stratopause levels. This can be explained by looking at the increase
of the water vapour content in Fig. 4b. As water vapour is more prominent in periods of decreasing UV radiation above 60 km (Fig. 3b), due to the decrease in photodissociation, production of OH via reaction with O(1D) is more likely. (Fig. 3b). Moreover a cooling of the upper troposphere/lower stratosphere (UT/LS) of 0.1 K decreases the stratospheric water content by 2%. Hence, as the mean decrease of temperature at UT/LS height (shown in the dynamics section) in our DM-TD run is of around 0.2 K, a decrease of roughly 4% of the stratospheric water content is to be expected – and modeled. This drop is the reason for the observed decrease in HOx below 65 km down to the tropopause of 4% on average.

To explain the strong increase of water vapour above the tropopause, a look at the results of DM-VOLC experiment (see Fig. 4c) is needed. The volcanic eruptions lead to a strong (of up to 14%) increase in the stratospheric water content even when the results are averaged over 20 yr long period. It is interesting to note that in the two years after Tambora, simulated water vapor contents raised by up to 60% at the tropical tropopause (not shown). Such a strong increase in water vapor content leads to an acceleration of reactions

\[ \text{H}_2\text{O} + \text{O}(^1\text{D}) \rightarrow \text{OH} + \text{OH} \]  
\[ \text{H}_2\text{O} + h\nu \rightarrow \text{H} + \text{OH} \]

in the lower mesosphere and stratosphere. Hence, an increase in HOx throughout the whole stratosphere and mesosphere (see Fig. 4c) is observed, with peak increases over the equatorial tropopause. The increase in HOx leads to a speed-up in oxidation of long-lived species like methane or CO (not shown).

In a clean and unpolluted atmosphere, a surplus of nitrogen oxides automatically leads to the drop of HOx concentrations. In Fig. 4d the additional NOx produced (see next subsection) from the GCRs decrease the amount of HOx slightly. These changes are marginally significant on a 5% level but highly significant on a 10% level in the tropical regions of the largest ionization rates (around 100 hPa). A decrease of HOx of up to 3% – during the absolute minimum of the DM even of up to 8% – is simulated by
our model. This decrease is supported by the slight additional decrease of 0.5 to 1% in the stratospheric water vapour content (Fig. 3d).

The nonlinearity analysis for both species show that significant changes only happen in the northern polar troposphere and are hence not illustrated here.

3.1.3 \( \text{NO}_x \)

In the DM-ALL experiment, the \( \text{NO}_x \) mixing ratio dramatically decreases in the polar mesosphere by up to 70% and – with a much smaller magnitude – also in the tropical middle stratosphere. In the tropical upper troposphere and in the stratosphere, increase in \( \text{NO}_x \) is found (see Fig. 5a) reaching its maximum in the upper tropical stratosphere/lower mesosphere. The latter can be explained by the smaller solar UV forcing in the DM-TD experiment (see Fig. 5b). The reduction in the solar UV irradiance leads to a pronounced decrease of the photolysis rates for all species including nitrogen oxide (NO). The NO absorption bands overlap with the oxygen Schumann–Runge bands (170–200 nm) and the introduced decline of the solar irradiance in this interval is one of the most pronounced (Shapiro et al., 2011). The NO photolysis \( (\text{NO} + h\nu \rightarrow \text{N} + \text{O}) \) plays a crucial role in the \( \text{NO}_y \) budget providing pure loss of \( \text{NO}_y \) via the subsequent cannibalistic reaction \( (\text{N} + \text{NO} \rightarrow \text{N}_2 + \text{O}) \) which explain the overall \( \text{NO}_x \) increase in the DM-TD experiment.

The dipole structure in the tropical UT/LS is explained by the influence of the volcanic eruptions. The volcanic sulfate aerosols provide a media for a number of fast heterogeneous reactions. For the clean stratosphere during the DM, the most important reactions was the \( \text{N}_2\text{O}_5 \) hydrolysis which facilitates the conversion of active nitrogen oxides to rather passive nitric acid. This effect is shown in Fig. 5c which illustrates the results of the DM-VOLC experiment. A significant \( \text{NO}_x \) decrease over the DM period is observed in this experiment above the tropopause over the tropics and high latitudes where the aerosol abundance is at maximum. The causes for a small \( \text{NO}_x \) increase in the tropical upper troposphere are not clear, probably it is related to the ozone increase in this area which leads to an enhanced \( \text{NO}_x \) production via \( \text{N}_2\text{O} + \text{O}(^1\text{D}) \rightarrow \text{NO} + \text{NO}. \)
As expected, most of the changes in NO$_x$ seen in Fig. 5a is dominated by energetic particles (see Fig. 5d). The NO$_x$ influx parameterized as the function of $A_p$ index weakened in intensity during the DM leading to NO$_x$ decrease by up to 80\%. Particles with higher energies – to a large part GCRs, whose flux was higher during the DM – penetrate deeper into the atmosphere. At tropopause levels, up to 6\% more NO$_x$ at the poles and of up to 2\% more NO$_x$ at the equator is produced by GCRs. While higher NO$_x$ concentrations at the poles in high altitudes above 50 km do not have a harmful effect on life, such a NO$_x$ production at lower altitudes leads to a slightly accelerated destruction of ozone via reactions 1–3. This insight could be of high importance for the possible oncoming new grand solar minimum in the current century. The NO$_x$-anomalies compare well to those found in Rozanov et al. (2012). The reason why the positive change in NO$_x$ is not reflected in Fig. 5a is that conversion to NO$_y$ occurs due to the additional amount of stratospheric aerosols from the volcanic eruptions (not shown). While in DM-EPP, only a slight significant increase of NO$_y$ is recognizable over the poles (not shown), a strong signal is seen in the DM-ALL experiment over the whole tropopause region.

The NO$_x$ field is a good example to show the non-linear behavior of atmospheric chemistry (see Fig. 6a). Superimposing all relative differences of all experiments, the mesospheric polar regions from the stacked DM-BU, DM-TD, DM-VOLC and DM-EPP result in significant differences of up to 15\% more NO$_x$ compared to DM-ALL. The NO$_x$ field in the lower stratosphere over the northern extratropics shows also a significant positive anomaly of 2–4\% more NO$_x$. The analysis of the NO$_y$ field in Fig. 6b shows an even more pronounced anomaly when superimposing all differences of all contributions together and comparing this result to the DM-ALL field: values of up to 25\% more NO$_y$ in the mesospheric polar atmosphere and in the northern polar midstratospheric region are reached. These differences stem partly from the NO-photolysis, which was kept to constant 1780 values during the DM-EPP run. As well, the additional cooling during the DM-VOLC run resulted in NO$_x$ deactivation over the poles, which could not happen during the DM-TD run.
3.2 Stratospheric dynamics

In this section we present the changes of the stratospheric temperature and winds during the DM and identify the contribution of all considered factors.

3.2.1 Temperature

As it is shown in Fig. 7a, during the DM the model simulates cooling in the entire atmosphere except the lower tropical stratosphere. The cooling gradually increases with altitude from 1 K in the middle stratosphere around 25 km up to 8 K near the mesopause. A slight enhancement of the cooling is also visible in the lower polar stratosphere. Weaker cooling of up to 0.6 K occurs below 100 hPa maximizing in the tropical upper troposphere.

Figure 7b shows the temperature changes due to implied decrease of visible and infrared solar irradiance (experiment DM-BU) and demonstrates that this factor is producing a weak cooling in the troposphere and in the upper stratosphere. The tropospheric cooling is explained by less energy income to the surface, while the cooling in the stratosphere is most likely caused by the decrease in available solar radiation for the ozone absorption in the Chappuis band. The temperature changes due to solar UV irradiance (experiment DM-TD) are illustrated in Fig. 7c which shows that this factor plays the dominant role in the cooling of the atmosphere above 25 km. The results shown in Fig. 7d demonstrate that the influence of volcanic eruptions has a more complicated spatial pattern. The volcanic aerosols produce a strong warming of up to 2 K at around 20 km in the tropical and subtropical region. The volcanic aerosol is able to absorb infrared solar and terrestrial radiation (Stenchikov et al., 1998, e.g.). The obtained strong warming in the lower tropical stratosphere means that the increased absorption of the terrestrial radiation by volcanic aerosols dominates over the absorption of the solar radiation, which should lead to a cooling due to the introduced decrease of the solar activity during the DM. This explains the strong increase of water vapour in the stratosphere which was illustrated in Sect. 3.1.2 The cooling in the
upper part of the model domain is explained by the blocking of the outgoing terrestrial radiation by the aerosol layer leading to a decrease of the energy incoming in these layers. The dipole-like structure of the temperature changes over the polar regions would hint on the intensification of the polar night jets: this suspicion is confirmed when analyzing seasonal means, which show a strong statistical significant acceleration on the 5\% level of the north polar night jet and a significant increase on the 10\% level of the southern polar night jet (not shown). In the troposphere, the blocking of the solar visible and infrared radiation by volcanic aerosols leads to a cooling in the stratosphere.

Due to dilution and gravitational settling as well as washout processes, the volcanic aerosols concentration is decreasing over time. This has an implication on the temperature anomalies: we find that lower stratospheric temperatures are going back to an average level around two years after the eruption (not shown), which is in agreement to the general findings (Robock, 2000).

The annual mean temperature changes from EPP are small and not statistically significant. Seasonal means however show significant differences: the austral winter seasonal mean show a dipole pattern over the southern pole (Fig. 9a). A significant drop in temperatures of up to 0.7 K between 100 hPa and 5 hPa is modeled in winter time, deepening in spring time to a cooling of up to 1 K (not shown). At the same time, a heating of approximately the same amplitude in a height between 1 hPa and 0.05 hPa is modeled. We explain the pattern over the southern pole by a strengthening of the polar vortex during austral winter (see next subsection) as well as a significant positive ozone anomaly of up to 3\% at 5 hPa (Fig. 2d). Ozone at these heights act as a radiative coolant. The positive temperature anomaly at mesosphere heights is due to a faster descent of air masses (BDC), leading to an increase in diabatic heating. No significant major changes in temperature can be observed during the boreal winter season.

The analyzed nonlinearities in temperatures are only significant in the troposphere and hence not shown here.
3.2.2 Wind and general circulation

The combined effect of all considered factors shown in Fig. 8a consists of a strong although only partly significant acceleration of the zonal winds in the subtropical stratosphere from 20 to 60 km and in the tropical stratosphere at around 50 km height. On the other hand, a significant deceleration of the tropical jets and a decrease of the mesospheric extra-tropical zonal winds are found.

Because the introduced decrease of solar irradiance does not have any significant influence on the annual mean zonal wind and only minor upper stratospheric influence at the southern polar region in austral winter time (not shown), all changes are attributed to the influence of the volcanic eruptions.

The model result shows a strong and significant deceleration of the zonal winds from the subtropical middle troposphere down to the surface of up to 0.8 m s\(^{-1}\). The downward propagation of the signal is observed in both hemispheres. This effect comes from the strong, extended warming of the entire tropical lower stratosphere by volcanic aerosols (see Fig. 8d), leading to a smaller temperature gradient from the equator to the extratropics. This weakens the subtropical jets. On the other hand, a significant strengthening of the polar jets up to 2.8 m s\(^{-1}\) is modeled, coming from the increased temperature gradient between the tropical tropopause and the polar tropopause (\(\Delta T = 3.2\) K). Moreover, in Fig. 10, a strong acceleration in vertical residual circulation (positive numbers are upwards) is observed after the 1809 and 1815 volcanic eruptions. Thus, the BDC is accelerated right after major volcanic eruptions. We explain this result by the finding that immediately after the volcanic eruption, cooling in the upper troposphere occurs. This favors the dissipation of gravity waves through the tropopause, leading to an additional gravity wave drag in the lower stratosphere and hence an acceleration of the BDC. It is only once the stratosphere heats up after that aerosols form that the vertical residual circulation drops due to strengthening of the temperature gradient at the tropopause.
The DM-EPP experiment indicates positive – however not significant – changes in annual zonal mean winds of up to 0.8 m s\(^{-1}\) at the stratospheric southern polar extratropics. This anomaly becomes highly significant in austral winter seasonal mean (Fig. 9b) and reaches values of up to 1.2 m s\(^{-1}\). The origin of this finding is the increase of adiabatic heating by descending air masses of the BDC in austral winter. The residual vertical circulation (not shown) shows a significant increase in down-ward motion of the air masses by up to 0.8 mm s\(^{-1}\). This in turn forms a positive temperature anomaly, leading to an increase in the pole-to-equator gradient at 60 km of height. As a consequence, the zonal wind increases. During boreal winter, a similar but less strong and non-significant pattern is found (not shown).

Nonlinearities in the zonal wind field are especially well visible at the southern extratropical stratopause, where a positive bias of up to 2 m s\(^{-1}\) is found (not shown). We assign this anomaly to the complex interaction between the DM-VOLC temperature gradient changes and the overall DM-TD negative temperature anomaly, which could have led to significant higher temperature gradient changes especially at the stratopause.

4 Conclusions

We present in this paper a modeling study of the different forcing which could have led to the dynamical and chemical changes in the stratosphere during the DM from 1805 to 1825 AD. The contributions, analyzed with four sensitivity experiments, include decrease in visible and near infrared radiation (DM-BU), UV radiation (DM-TD), volcanic eruptions (DM-VOLC) and energetic particles, the latter including an increase of galactic cosmic rays ionization and a decrease of solar proton events and low energetic electron precipitation (DM-EPP). A comparison of DM-TD, DM-VOLC and DM-EPP to the control run showed major significant changes. However, when comparing DM-BU to the control run, we did not succeed in identifying any noticeable changes neither in stratospheric chemistry nor in the stratospheric dynamics – excepted for the mid-
stratospheric temperature field. This is mainly due to the drop of only 1% in the visible radiation band from the Shapiro et al. (2011) reconstruction.

When isolating the different contributions, our simulations show following effects on temperatures: when reducing the solar radiation in the 185–250 nm band and the photolysis rates, a temperature drop, reaching higher negative anomalies at the mesosphere than at the tropopause is modeled. The anomalies reach 0.2 K at the tropopause and quasi-uniformly drop to values down to 6 K at the lower mesosphere. The main reason is the lower absorption of radiation by ozone. A significant cooling of up to 0.6 K is observed in the middle stratosphere when reducing the visible part of the spectrum (250–690 nm). In the volcanic scenario, the tropical tropopause is heated by up to 2.2 K due to radiative absorption by the aerosols while a significant cooling of up to 1.2 K is modeled around the stratopause. In the DM-EPP scenario, only a seasonal significant change in temperatures could be modeled in JJA, leading to a dipole-like structure: a cold anomaly of 0.5 K in the southern polar middle stratosphere and a warming of around the same amount in the southern polar lower mesosphere.

The zonal winds did not change significantly neither in the DM-BU nor in the DM-TD experiment. However, a strong significant change is modeled in the subtropics and extratropics in both hemispheres in the lower and upper stratosphere by up to 2.8 ms$^{-1}$ when forcing the model only with volcanic aerosols. This effect is highest in the two years following an important volcanic eruption. Same as for the temperature, changes in zonal winds in the DM-EPP scenario are only significant in the DJF seasonal mean, showing a zonal wind increase of up to 1.2 ms$^{-1}$ in the southern extratropical stratosphere, which is due to the dipole-like temperature anomalies.

The analysis of the stratospheric chemistry lead to following conclusions: ozone drops by up to 8% in the ozone layer, HO$_x$ decreases by up to 20% at the stratopause and water vapour content decreases in the low stratosphere by up to 3.6% but increases by up to 40% at the lower mesosphere in the DM-TD scenario. In the DM-VOLC scenario, ozone increases by up to 16% at the tropical tropopause but decreases by up to 6% at the stratopause, HO$_x$ increases all over by up to 25% as
does the water vapour amount by up to 14%. These effects are highest in the two subsequent years after a major volcanic eruption (up to 60% more water vapour in the lower stratosphere in the 2 yr after Tambora). The DM-EPP experiment showed highly significant changes in the NO$_x$ field: while decrease of up to 80% is modeled at the mesospheric poles, increase of up to 4% is simulated at the polar tropopause. However, no or very low effects are modeled on ozone, HO$_x$ and water vapour.

By considering the changes in dynamics and chemistry, we conclude that only due to the complex interaction of volcanic, UV solar spectral and EPP forcing, these contributions induce changes in dynamics and chemistry of the stratosphere during the DM. The reduction of the visible radiation plays only a minor role in most of the fields except for the temperature. Thus, for future modeling studies, including an interactive chemistry with separate treatment of the different spectral bands is of great importance in order to get the climate responses on solar- and volcanic forcing as realistic as possible.

Concluding, the ozone decrease was predominantly influenced by the decrease in UV radiation in the polar mesosphere and at ozone layer height, whereas the volcanic eruptions influenced ozone concentrations at tropical tropopause height. EPP influenced only in a minor part ozone concentrations in the polar mesosphere. HO$_x$ and water vapor increase were affected primarily by volcanic eruptions in the stratosphere and by UV in the lower mesosphere. NO$_x$ fields were most notably influenced by EPP in the polar mesosphere and by UV in the upper stratosphere. Stratospheric winds were influenced to a major part from volcanic eruptions. Temperatures were mainly influenced by volcanic eruptions and UV reduction, leading to a significant warming at the tropical tropopause and to a cooling in the remaining of the atmosphere.

For an outlook, the drop of ozone by up to 7% at ozone layer height when reducing the UV radiation should be kept in mind when looking at the possible future grand solar minimum of the 21st century. A similar – or an even greater – decrease due to ozone depleting substances and UV radiation reduction then gets a possible treat to life on Earth. As well, the effects of reduction of UV, volcanic eruptions and increase of oxida-
tion by GCRs should be thoroughly investigated in future research of the 21st century with an AOC-GCM due to halogen and anthropogenic NO\textsubscript{x} loading. The evolution of the ozone layer might be of utmost interest to the whole scientific community, as e.g. crop yields or skin cancer could increase due to GHG interactions with stratospheric chemistry. Moreover, the observed changes of the zonal winds and in temperatures in the troposphere during all sensitivity experiments should be analyzed carefully in a future work, as they might have influenced directly the climate on the Earth's surface.

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Table 1. Experiments for DM sensitivity runs: CONST values are 1780 monthly mean values. BCKGRD means that only background aerosol emissions were enabled while volcanic eruptions were turned off. TRANS means transient forcing.

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<th>EPP</th>
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Fig. 1. Spectral solar irradiance during the DM adapted from Shapiro et al. (2011). From left to right, top to bottom: hard UV (185–250 nm), soft UV (250–440 nm), visible (440–690 nm), IR-A (690–1190 nm), IR-B (1190–2380 nm) and IR-C (2380–4000 nm).
Fig. 2. Relative differences of yearly mean ozone of the DM-ALL, DM-TD, DM-VOLC and DM-EPP experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student’s $t$ test with $\alpha = 5\%$. The yellow line illustrates the height of the WMO tropopause.
**Fig. 3.** Relative differences of yearly averaged water vapour of the DM-ALL, DM-TD, DM-VOLC and DM-EPP experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student’s t test with $\alpha = 5\%$. The yellow line illustrates the height of the WMO tropopause.
Fig. 4. Relative differences of yearly averaged HO$_x$ of the DM-ALL, DM-TD, DM-VOLC and DM-EPP experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student’s $t$ test with $\alpha = 5\%$. The yellow line illustrates the height of the WMO tropopause.
Fig. 5. Relative differences of NO$_x$ of the DM-ALL, DM-TD, DM-VOLC and DM-EPP experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student’s $t$ test with $\alpha = 5\%$. The yellow line illustrates the height of the WMO tropopause.
Fig. 6. Differences between the DM-ALL vs. DM-CTRL1780 field and the (DM-TD + DM-BU + DM-VOLC + DM-EPP) vs. DM-CTRL1780 field. Positive values show a positive anomaly of the stacked differences over the DM-ALL difference field. Hatched areas are significantly different on a Student’s t test with $\alpha = 10\%$. The yellow line illustrates the height of the WMO tropopause.
Fig. 7. Absolute differences in temperatures of the DM-ALL, DM-BU, DM-TD and DM-VOLC experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student’s $t$ test with $\alpha = 5\%$. The yellow line illustrates the WMO tropopause height.
Fig. 8. Absolute differences in mean zonal wind of the DM-ALL and the DM-VOLC experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student's $t$ test with $\alpha = 5\%$. The yellow line illustrates the WMO tropopause height.
Fig. 9. Absolute DJF seasonal differences in temperature (left) and mean zonal wind (right) of the DM-EPP experiments with relation to the DM-CTRL1780 forcing run. Hatched areas are significantly different on a Student's t test with $\alpha = 5\%$. The yellow line illustrates the WMO tropopause height.
Fig. 10. Residual vertical velocity (calculation as in Andrews et al. (1987), Eqs. 3.5.1 to 3.5.3) from 1805 to 1820 for the DM-CTRL1780, DM-VOLC and DM-TD simulations at 100 hPa height averaged from 20° N/20° S. The two major volcanic eruptions in 1809 (unknown) and 1815 (Tambora) are marked with two vertical orange lines. The data is not smoothed.