CIOOCI photolysis at high solar zenith angles: analysis of the RECONCILE self-match flight

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Received: 13 May 2011 – Accepted: 28 May 2011 – Published: 1 July 2011

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Published by Copernicus Publications on behalf of the European Geosciences Union.
Abstract

The photolysis frequency of dichlorine peroxide (ClOOCl) $J_{\text{ClOOCl}}$ is a critical parameter in catalytic cycles destroying ozone in the polar stratosphere. In the atmospherically relevant wavelength region, published laboratory measurements of ClOOCl absorption cross sections and spectra are not in good agreement, resulting in significant discrepancies in $J_{\text{ClOOCl}}$. Previous investigations of the consistency with atmospheric observations of ClO and ClOOCl have focused on the photochemical equilibrium between ClOOCl formation and photolysis, and thus could only constrain the ratio of $J_{\text{ClOOCl}}$ over the rate constant of the ClO recombination reaction $k_{\text{rec}}$. Here, we constrain the atmospherically effective $J_{\text{ClOOCl}}$ independent of $k_{\text{rec}}$ using ClO data sampled in the same air masses before and directly after sunrise. Over sunrise, when the ClO/ClOOCl system comes out of thermal equilibrium and the influence of the ClO recombination reaction is negligible, the rise in ClO concentration is significantly faster than expected from $J_{\text{ClOOCl}}$ based on the absorption spectrum proposed by Pope et al. (2007), but does not warrant cross sections larger than recently published values by Papanastasiou et al. (2009). In particular, the existence of a significant ClOOCl absorption band longwards of 420 nm, is effectively ruled out by our observations. Additionally, the night-time ClO observations show that the ClO/ClOOCl thermal equilibrium constant can not be significantly higher than the one proposed by Plenge et al. (2005).

1 Introduction

Once chlorine is activated on PSCs (Solomon et al., 1986) or background aerosol (Drdla and Müller, 2010) in the polar stratosphere and sunlight is available, photochemical ozone loss occurs essentially via two catalytic cycles, the ClO-dimer cycle (Molina and Molina, 1987):

$$\text{ClO} + \text{ClO} + \text{M} \xrightleftharpoons[k_{\text{diss}}]{k_{\text{rec}}} \text{ClOOCl} + \text{M} \quad (R1)$$
ClOOCI + \( h\nu \xrightarrow{J_{\text{ClOOCI}}} \) Cl + ClOO

ClOO + M → Cl + O₂ + M

2 × (Cl + O₃ → ClO + O₂)

\[ \text{ClOOCl} + h\nu \rightarrow \text{Cl} + \text{ClOO} \]

\[ \text{ClOO} + M \rightarrow \text{Cl} + \text{O}_2 + M \]

\[ 2 \times (\text{Cl} + \text{O}_3 \rightarrow \text{ClO} + \text{O}_2) \]

\[ 2\text{O}_3 + h\nu \rightarrow 3\text{O}_2 \]

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\[ \text{ClOOCl} \rightarrow \text{Cl} + \text{ClOO} \]

\[ \text{ClOO} + \text{M} \rightarrow \text{Cl} + \text{O}_2 + \text{M} \]

\[ 2 \times (\text{Cl} + \text{O}_3 \rightarrow \text{ClO} + \text{O}_2) \]

\[ 2\text{O}_3 + h\nu \rightarrow 3\text{O}_2 \]

\[ \text{ClO} + \text{BrO} \rightarrow \text{Br} + \text{Cl} + \text{O}_2 \]

\[ \text{Br} + \text{O}_3 \rightarrow \text{BrO} + \text{O}_2 \]

\[ \text{Cl} + \text{O}_3 \rightarrow \text{ClO} + \text{O}_2 \]

\[ 2\text{O}_3 + h\nu \rightarrow 3\text{O}_2 \]

\[ \text{BrO} + \text{ClO} \rightarrow \text{Br} + \text{Cl} + \text{O}_2 \]

\[ \text{Br} + \text{O}_3 \rightarrow \text{BrO} + \text{O}_2 \]

\[ \text{Cl} + \text{O}_3 \rightarrow \text{ClO} + \text{O}_2 \]

\[ 2\text{O}_3 + h\nu \rightarrow 3\text{O}_2 \]

The forward and backward Reaction (R1) of the ClO-dimer cycle, i.e. the ClO recombination and the ClO dimer (ClOOCI) thermal dissociation, govern the partitioning between ClO and ClOOCl in darkness with a thermal equilibrium constant

\[ K_{\text{eq}} = \frac{k_{\text{rec}}}{k_{\text{diss}}} = \frac{[\text{ClOOCl}]}{[\text{ClO}]^2}. \]

At daylight, Reactions (R2)–(R4) shift the ClOₙ (≈ ClO + 2ClOOCl) partitioning towards the monomer and drive chemical ozone loss (Reaction R4). The ClOOCI photodissociation Reaction (R2) is the rate-limiting step of the ClO-dimer cycle under twilight conditions that prevail throughout most of the winter in the polar stratosphere. By governing the ClO concentration, Reaction (R2) also limits the rate of Reaction (R5)
in the ClO-BrO cycle, making the ClOOCI photolysis frequency \( J_{\text{ClOOCI}} \) by far the most critical kinetic parameter for the overall ozone loss rate (von Hobe et al., 2007).

\( J_{\text{ClOOCI}} \) is a product of the ClOOCI photolysis cross section (absorption cross section \( \sigma_{\text{ClOOCI}}(\lambda) \) times the photolysis quantum yield \( \phi(\lambda) \), i.e. the fraction of absorbed photons leading to photodissociation; plotted in Fig. 1a) and the actinic flux \( I(\lambda) \) (Fig. 1b) integrated over wavelengths:

\[
J_{\text{ClOOCI}} = \int \sigma_{\text{ClOOCI}}(\lambda) \cdot \phi(\lambda) \cdot I(\lambda) \cdot d\lambda
\]  

Photolysis cross sections have been determined by monitoring the disappearance of ClOOCI (Chen et al., 2009; Lien et al., 2009; Jin et al., 2010) or the appearance of Cl atoms (Wilmouth et al., 2009) in photolysis experiments at discrete wavelengths. For the wavelength range relevant for ClOOCI photolysis in the stratosphere (\( \lambda > 310 \text{ nm} \)), only absorption cross sections (Cox and Hayman, 1988; DeMore and Tschuikow-Roux, 1990; Burkholder et al., 1990; Bloss et al., 2001; Papanastasiou et al., 2009) or relative absorption spectra (Huder and DeMore, 1995; von Hobe et al., 2009; Pope et al., 2007) have been measured by UV/Vis absorption spectroscopy. Because excited states of ClOOCI are rapidly dissociative (Birk et al., 1989; Kaledin and Morokuma, 2000; Toniolo et al., 2001; Peterson and Francisco, 2004), absorption and photolysis cross sections are often assumed to be equivalent, i.e. \( \phi(\lambda) \sim 1 \).

The laboratory studies, in which absorption spectra have been measured, are lacking the knowledge of the amount of ClOOCI present in the absorption cell and the relative spectra have to be scaled to one of the following absolute cross section measurements. The studies of Cox and Hayman (1988), DeMore and Tschuikow-Roux (1990) and Burkholder et al. (1990) indicated the absorption peak value of ClOOCI at 245 nm between \( 5.8 \times 10^{-18} \) and \( 7.2 \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \). According to the results of Papanastasiou et al. (2009), the absorption peak is at 244.5 nm and has a value of \( 7.6 \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \), Wilmouth et al. (2009) placed the photolysis peak of \( 6.6 \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \) at 248.0 nm, whereas studies of Lien et al. (2009) showed even higher photolysis cross section (\( 8.85 \times 10^{-18} \text{ cm}^2 \text{ molecule}^{-1} \)) near to the
assumed peak, at 248.4 nm. Here, we chose to scale all relative absorption spectra to the cross section of Lien et al. (2009) for two reasons:

- Lien et al. (2009) have used a novel approach in which a molecular beam of ClOOCl is attenuated through photolysis by laser light. The method is insensitive to UV absorption interference by impurities and does not require information on the absolute ClOOCl concentration,

- it seems unreasonable to scale a relative absorption spectrum to an absolute cross section associated with a different spectral shape.

The upper panel of Fig. 1 shows that the existing laboratory studies on $\sigma_{\text{ClOOCl}}$ display large uncertainties especially in the long wavelength tail, i.e. in the region of most relevance in the atmosphere. These uncertainties propagate directly into the atmospheric $J_{\text{ClOOCl}}$ values derived from the cross section measurements. Figure 1c reveals the significance of the photolysis longward of 400 nm at high solar zenith angles (SZA) and thus emphasizes the importance of the increasingly small cross sections in this wavelength region. The inability to tightly constrain $J_{\text{ClOOCl}}$, that governs the ClO-dimer cycle, is one of the key issues preventing full quantitative understanding of the ozone depletion in the winter/early spring polar atmosphere.

Here, we investigate four sets of ClOOCl absorption cross sections/scaled spectra (Sect. 2) and their influence on [ClO] rise at high SZA. For comparison, we use pairs of in-situ ClO measurements in the same air mass just before and after sunrise obtained during a so called self-match flight. The self-match flight pattern and the measurements are described in Sects. 3.1 and 3.2. Section 3.3 describes in detail the approach employed to test the consistency of the four different $J_{\text{ClOOCl}}(\lambda)$ values with the observed rise in ClO concentration over the sunrise. An assessment of the success in resampling the same air in the matches based on observations of a chemically and dynamically conserved tracer and the results of the ClO/ClOOCl kinetic studies will be presented in Sect. 4.
2 Investigated $\sigma_{\text{ClOOCI}}$

In further analysis two relative spectra, absolute cross sections and an artificial spectrum are considered:

i. Pope et al. (2007) presented ClOOCI spectrum with the lowest ever published absorption cross sections at $\lambda > 310\, \text{nm}$ resulting in relatively small $J_{\text{ClOOCI}}$ values. They carried out an innovative experiment purifying ClOOCI prior to the absorption measurement, thus removing all impurities other than Cl$_2$. However, in their spectral analysis they probably overcorrected for the Cl$_2$ impurity as has been shown by von Hobe et al. (2009) (see below). The spectrum is scaled to measurements of Lien et al. (2009) to obtain cross sections.

ii. von Hobe et al. (2009) used the same technique proposed by Pope et al. (2007) to prepare and purify ClOOCI, but measured the spectrum of sample isolated in a solid noble gas matrix. In the atmospherically relevant region, their spectrum lies considerably higher than the Pope et al. (2007) spectrum but compares rather well to a Pope et al. (2007) gas phase spectrum uncorrected for Cl$_2$. The von Hobe et al. (2009) spectrum is also scaled to cross section of Lien et al. (2009).

iii. Papanastasiou et al. (2009) measured absolute cross sections over a wavelength range from the ClOOCI absorption peak into the visible. The new measurement refines an earlier study from the same laboratory (Burkholder et al., 1990). The two spectra have slightly different shapes but yield similar atmospheric $J_{\text{ClOOCI}}$ values. Burkholder et al. (1990) results generally produced the best agreement in a number of studies comparing modelled and observed ClO$_x$ partitioning (Stimpfle et al., 2004; von Hobe, 2007; Schofield et al., 2008; Kremser et al., 2010) as well as ozone loss (Santee et al., 2003; Chipperfield et al., 2005; Frieler et al., 2006; Tripathi et al., 2006).

iv. An artificial spectrum based on Papanastasiou et al. (2009), but with a simulated $\sigma_{\text{ClOOCI}}$ set constant to $7.12 \times 10^{-21}\, \text{cm}^2\, \text{molecule}^{-1}$ above $420\, \text{nm}$ is used to test
the possible existence of an additional ClOOCl absorption band in the visible.

\( J_{\text{ClOOCl}} \) values derived from the cross sections i–iv and the actinic flux from Fig. 1b are shown in Fig. 1c.

\( I(\lambda) \) used in Fig. 1b and in further analysis was calculated using a HALOE \( O_3 \) climatology (Groß and Russell, 2005) and CLaMS photolysis code (Becker et al., 2000) for an altitude of 18.4 km (corresponding to the flight altitude of the RECONCILE self-match flight, cf. Sect. 3.1). To compare the ClO rises resulting from the considered photolysis rates with ClO observed during the self-match flight, we use the overall \( J_{\text{ClOOCl}} \) effective in the atmosphere obtained by integration of the areas under the curves plotted in Fig. 1c.

3 Experiment description

3.1 The self-match flight

The idea behind a self-match flight is to sample the same air masses twice during outbound and inbound flight legs of the same flight as has been described by Schofield et al. (2008). The flight was carried out on 30 January 2010 from Kiruna (67° 49′ N, 20° 20′ E), Sweden, as part of the large aircraft field campaign within the European project RECONCILE (von Hobe et al., 2011), employing the Russian research aircraft M-55 Geophysica with a ceiling altitude of 20 km. Flight track and match-pair trajectories are shown in Fig. 2.

To analyse the SZA dependence of \( J_{\text{ClOOCl}} \), the flight was carried out over sunrise, with the first measurements still in darkness (i.e. with no active photolysis reactions and the chemical system expected to be in thermal equilibrium), and the second “matching” measurement in daylight, moving from higher to lower solar zenith angles (i.e. increasing solar radiation and photolysis rates) over the course of the inbound flight leg as is shown in Fig. 2. This was achieved by a nearly sun-synchronous flight pattern in westward direction for the outbound flight leg and the inbound flight leg coming
back in eastward direction. Location and timing of the flight were planned as to direct the aircraft into a region of elevated \( \text{ClO}_x \). This was identified prior to the flight using the chemistry-transport model CLaMS (see Sect. 3.3) in the forecast mode. Air mass trajectories for the motion of individual air parcels between the two encounters were calculated based on data from the European Centre for Medium-Range Weather Forecasts (ECMWF).

### 3.2 Measurements

Chlorine monoxide (ClO) was measured by the HALOX (HALogen OXide monitor) instrument situated in a pod underneath the left wing of the Geophysica aircraft (von Hobe et al., 2005). Through an inlet optimised for radical measurements, ambient air is pumped through two measurement ducts where periodically (10 s cycles) NO is added, converting ClO to Cl atoms that are detected downstream by resonance fluorescence in the vacuum UV. HALOX has a detection limit of 5 ppt and an accuracy of about 18%. Most critical for the analysis presented below is precision, which varies between 4% and 8% depending on the stability of the light source (helium discharge lamps containing chlorine) and on the observed detector noise. Random noise is reduced and signal-to-noise-ratio increased by integrating the signal over six NO addition cycles, yielding a time resolution of 1 min.

Temperature and pressure were measured using commercial Rosemount sensors and geolocation data were provided by the M-55 avionic system.

To examine the accuracy of matches (cf. Sect. 4.1), we used measurements of nitrous oxide \( (\text{N}_2\text{O}) \), which is expected to be unaffected by chemical processes on the time scale of the self-match flight and thus should show conserved mixing ratios. \( \text{N}_2\text{O} \) was measured by the 2-channel gas chromatograph HAGAR (High Altitude Gas Analyser, Volk et al., 2000) with a time resolution of 90 s. Mean precision and accuracy are better than 1% and 2% of the tropospheric background value, respectively.

Number concentrations of the particles with diameters of 6 nm, 11 nm and 15 nm were measured by COPAS (COndensation PArticle Counting System; Weigel et al., 2000).
2809), which consists of an aerosol inlet and two dual-channel continuous flow Con-
28densation Particle Counters (CPCs). The particles are detected in three channels op-
erated with distinct temperature differences between the saturator and the condenser
block. An other channel is operated with an aerosol heating line for a determination of
the number of non-volatile particles. In the COPAS data, the exhaust gases of Geo-
physica were observed several times on the inbound flight leg confirming the good
performance of the self-match flight.

3.3 CLaMS simulations

To define the match trajectories and to compare the ClO observations to values ex-
pected from various combinations of kinetic parameters, the Chemical Lagrangian
Model of the Stratosphere (CLaMS) (McKenna et al., 2002a,b; Grooß et al., 2005) was
used in a box-trajectory model mode, with trajectory calculations based on ECMWF
wind fields and temperatures.

94 match trajectory start points were defined by the one-minute time intervals at
which ClO was measured on the outbound flight leg. Forward trajectories were cal-
culated until 12:00 UTC of 30 January with 1 min time resolution. The points with the
lowest distance to the inbound flight track were then defined as pairing match points
(cf. Sect. 4.1).

Chemistry calculations were also carried out on these trajectories with various com-
binations of $J_{\text{ClOOCl}}$, $k_{\text{rec}}$, and $K_{\text{eq}}$ values (cf. Sect. 4.4). For initialisation of the chemistry
calculations, trace gas mixing ratios measured on the Geophysica (full list of sampled
species published by von Hobe et al., 2011) during the outbound flight leg where used.
Species not measured were taken from two-day back trajectory calculations initialised
from a full hemispheric CLaMS simulation over the entire winter. The analysis of the
backward trajectories revealed that the investigated air masses had been in darkness
longer than 15 h prior to the flight, justifying the assumption of thermal equilibrium for
the outbound flight leg (von Hobe et al., 2005). During these 15 h, temperature did not
fall below 196.6 K on any of the trajectories, ruling out significant contribution of ClO
rise from Cl₂ released during very recent heterogeneous activation.

4 Results

4.1 Matches

Match points on the inbound flight leg corresponding to those on the outbound leg were found by minimising the horizontal radius around the coordinates of the Geophysica track and finding the coordinates fitting in time on the CLaMS forward trajectories. As shown in Fig. 3, the match radii \( R_{\text{match}} \) were always smaller than 8.5 km. To define a successful match pair, the vertical difference, in terms of potential temperature, between the locations of observations on the out- and inbound flight parts \( \Delta \Theta \) also needs to be examined. Here, we only consider matches with \( \Delta \Theta < 2.0 \) K. The vertical ClO gradient measured during descent was 15 ppt K\(^{-1}\).

To further test the accuracy of the matches, observations of the long-lived tracer \( \text{N}_2\text{O} \) by the HAGAR instrument were used. Match pairs, in which the difference between \( \text{N}_2\text{O} \) measured on the outbound and inbound flight legs was larger than the precision of HAGAR for the relevant data points (\( \approx 1.6 \) ppb), were rejected. With few exceptions, the tracer measurements confirmed an excellent match performance (Fig. 3, lower panel), as do several encounters of the Geophysica exhaust on the inbound flight leg identified by COPAS observations.

41 match pairs fulfilled all the conditions and are marked by large black circles in Fig. 3.

4.2 Night-time ClO\(_x\) chemistry

Figure 4 shows ClO\(_x\) mixing ratios calculated from the night-time ClO measurements and \( K_{\text{aq}} \) published by Plenge et al. (2005), Ferracci and Rowley (2010) and taken from JPL 2009 (Sander et al., 2009). An upper limit for ClO\(_x\) is given by the total inorganic
chlorine Cl\textsubscript{y}, which was obtained from the tracer-tracer correlation with N\textsubscript{2}O published by Grooß et al. (2002). The N\textsubscript{2}O increase of 2.4 % over the past 10 yr was taken into account while no significant change in stratospheric Cl\textsubscript{y} was assumed.

\[ Cl\textsubscript{y} = 3.394 + 0.001648 \cdot N\textsubscript{2}O - 7.659^{-5} \cdot N\textsubscript{2}O^2 + 2.646^{-7} \cdot N\textsubscript{2}O^3 - 4.469^{-10} \cdot N\textsubscript{2}O^4 \]  

N\textsubscript{2}O and Cl\textsubscript{y} are given in ppb and the relation is valid for N\textsubscript{2}O values between 17–320 ppb.

Only the \( K_{eq} \) from Plenge et al. (2005) is consistent with observed night-time ClO. \( K_{eq} \) taken from studies of Ferracci and Rowley (2010) and JPL 2009 would result in ClO\textsubscript{x} exceeding Cl\textsubscript{y}. Therefore in our study \( K_{eq} \) from Plenge et al. (2005) is used to estimate ClOOCl for the further analysis.

4.3 Constraints on \( J_{\text{ClOOCl}} \) from the observed increase in ClO

The overall rate of change in ClO concentration resulting from the reactions system (R1)–(R4) is given by:

\[ 0.5 \frac{d[ClO]}{dt} = J_{\text{ClOOCl}}[ClOOCl] + k_{\text{diss}}[M][ClOOCl] - k_{\text{rec}}[M][ClO]^2 \]  

Figure 5 illustrates the SZA dependence of \( J_{\text{ClOOCl}} \) for the \( \sigma_{\text{ClOOCl}} \) i–iv for the conditions encountered during the RECONCILE self-match flight. The ClOOCl photolysis term is zero in darkness and rises more or less sharply, depending on the choice of \( \sigma_{\text{ClOOCl}} \), as SZA decreases. Also shown is the ClOOCl thermal decomposition rate constant \( k_{\text{diss}}[M] \) for the mean temperature and pressure observed during the flight. In thermal equilibrium, the rate of ClO production from this process is expected to equal the removal rate via the ClO self-reaction, i.e. \( k_{\text{diss}}[M][ClOOCl] = k_{\text{rec}}[M][ClO]^2 \) and \( d[ClO]/dt = 0 \). Thus, when the sun rises, ClO mixing ratios are expected to increase as soon as \( J_{\text{ClOOCl}} \) becomes large enough so that the photolysis reaction presents a significant additional ClO production term compared to the ClOOCl thermal decomposition. This point occurs at different SZA for the different \( \sigma_{\text{ClOOCl}} \) tested and can be compared to the SZA when a significant rise of ClO mixing ratios is actually observed.
Figure 6 shows that observed ClO starts to rise close to a SZA of 92°. Also shown is the simulated rise in ClO expected for the photolysis rates i–iv. This was calculated making two assumptions: [ClOOCl] is assumed to remain constant at its night-time concentration (calculated from observed night-time ClO concentrations and \( K_{eq} \) according to Plenge et al., 2005), and ClO removal from the self-reaction (Reaction R1) is ignored. Strictly, these assumptions hold true only at sunrise, i.e. when the first ClOOCl molecule is photolysed. As SZA decreases, [ClOOCl] will become smaller and ClO removal via the self-reaction faster, both leading to a reduced overall rate of ClO increase. Thus, the observed ClO increase and the integrated simulated ClO increase from ClOOCl photolysis are expected to start rising simultaneously at the threshold SZA with the observations then falling more and more below the theoretical line as SZA decreases.

Simulated ClO rise calculated for i) (scaled absorption spectra published by Pope et al., 2007) increases too late and proceeds too slow when compared with observed ClO. On the other end of the spectrum, a significant ClOOCl absorption band in the visible (case iv) would lead to an even earlier increase in ClO and is thus not warranted based on our observations. \( J_{\text{ClOOCl}} \) values based on ii) and iii) produce a reasonably good agreement between observed and expected ClO rise within the given uncertainties. They can probably be regarded as maximum and minimum of plausible photolysis cross sections effective in the atmosphere.

4.4 Chemistry simulations along match trajectories

In this analysis, box model runs with full chemistry calculations are initialised with the measurements made on the outbound flight leg. For calculation of ClOOCl mixing ratios, HALOX ClO and \( K_{eq} \) from the publication of Plenge et al. (2005) are used (cf. Sect. 4.2). The data sets resulting from the chemistry simulation runs along the trajectories between the points on the outbound and the inbound flight leg are compared with values observed on the inbound flight leg. The results are shown in Fig. 7.
Simulations were carried out for combinations of the four $J_{\text{ClOOCl}}$ parametrisations discussed above and two parametrisations of $k_{\text{rec}}$ (JPL 2009 and Nickolaisen et al., 1994). $K_{\text{eq}}$ was taken from Plenge et al. (2005) (cf. Sect. 4.2).

Irrespective of the agreement between modelled and observed ClO, Fig. 7 illustrates the point made earlier, that immediately after sunrise, ClO mixing ratios are much more sensitive towards $J_{\text{ClOOCl}}$ than towards $k_{\text{rec}}$. Only as the sun rises higher and ClO builds up, the influence of $k_{\text{rec}}$ becomes important. Clearly, at high SZA at sunrise, the steady state approximation $J_{\text{ClOOCl}}/k_{\text{rec}} \approx [\text{ClO}]^2/[\text{ClOOCl}]$ is not valid.

In agreement with the results presented in Sect. 4.3, for the scaled cross sections of Pope et al. (2007), simulated ClO falls below observed mixing ratios as soon as they start to rise. In case ii), the model consistently follows the observations at very high SZA over sunrise, but predicts significantly smaller ClO mixing ratios than observed as SZA decreases further and photochemical steady state is approached. At these relatively low SZAs, simulations iii) and iv) show relatively good agreement with the observations, but they overestimate ClO earlier, especially cross section iv) with the artificial absorption band in the visible.

Provided that there are no additional reactions other than Reactions (R1)–(R4) governing the ClO-dimer cycle, the differences between model and observations for cases i) and iv) are too large for the underlying cross sections to be realistic in the atmosphere. For cases ii) and iii), the behaviour of the scaled von Hobe et al. (2009) parametrisation underestimating observed ClO at SZA < 90° and Papanastasiou et al. (2009) overestimating ClO at SZA > 90°, suggests ClOOCl cross sections within the range between those two. The HALOX measurements indicate an SZA threshold of significant ClO increase due to ClOOCl photolysis between 91.5° and 92° (cf. Fig. 6). That rules out the exponential extrapolation of $\sigma_{\text{ClOOCl}}$ of Papanastasiou et al. (2009), because of its moderate slope at wavelengths larger than 350 nm, which would cause more than observed ClOOCl photolysis at too high solar zenith angles. The spectrum fully fitting our observed data is supposed to have the steep slope and shape of von Hobe et al. (2009) spectrum, but the scaling should be higher providing larger
photolysis at $\lambda < 350$ nm.

5 Conclusions

CIO measurements during the RECONCILE self-match flight on 30 January 2010 were used to examine the plausibility of published ClOOCl absorption cross sections measured in the laboratory and to test the hypothesis of an additional ClOOCl absorption in the visible. The analysis of the increase in ClO concentration over sunrise as well as CLaMS chemistry simulations suggest a ClOOCl photolysis frequency in the range resulting from absorption cross sections published by von Hobe et al. (2009) scaled to Lien et al. (2009) and by Papanastasiou et al. (2009). Cross sections published by Pope et al. (2007) are clearly inconsistent with our observations unless there are additional unknown processes converting ClOOCl to ClO in the atmosphere. The observed behaviour of ClO mixing ratios starting to increase only at SZA below 92° also rules out an additional absorption band of ClOOCl at wavelengths $> 420$ nm.

Additionally, the analysis of night-time ClO indicates that our observations are inconsistent with equilibrium constants higher than $K_{eq}$ published by Plenge et al. (2005).

Acknowledgements. The authors thank the M-55 Geophysica team and the MDB (Myasishchev Design Bureau, Moscow, Russia) for planning and carrying out the flights and ENVISCOPE GmbH for logistical support during the RECONCILE campaign. We are also grateful to the mission scientist H. Schlager from DLR Oberpfaffenhofen for organisation and flight planning and to ECMWF for providing the meteorological analysis.

The RECONCILE project is funded under the European Commission Seventh Framework Programme (FP7) under the Grant number RECONCILE-226365-FP7-ENV-2008-1. The Ph.D. student position of Olga Sumińska-Ebersoldt is financed within the ENVIVAL-LIFE project through a grant from Bundesministerium für Wirtschaft (BMWi) via DLR, Germany.
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Fig. 1. Panel (a) presents a comparison of UV/Vis absorption spectra, absorption cross-sections and photolysis cross sections of ClOOCl from various studies. The dashed extensions show the exponential extrapolation of chosen spectra in the actinic region. Spectra i–iv used in this study are shown with thicker lines, with the artificial spectrum for examination of the hypothesis of a ClOOCl absorption band in the visible represented by a dashed orange line. Panel (b) shows spectral actinic flux $I(\lambda)$ calculated for four solar zenith angles. The photolysis rate constants estimated for absorption cross sections i–iv and $I(\lambda)$ are shown in panel (c).
Fig. 2. Flight path of the RECONCILE self-match flight on 30 January 2010. The solar zenith angle on the track is represented by the colour scale. The black lines between outbound and inbound parts of flight represent the trajectories of the match pairs fulfilling the strict match conditions (cf. Sect. 4.1).
Fig. 3. Performance of the match-flight. The upper panel and the middle panels show the horizontal and the vertical match accuracy in form of match radii $R_{\text{match}}$ (in km) and potential temperature difference $\Delta \Theta$ (in K), respectively. Differences in the mixing ratios of N$_2$O are shown in the lower panel. The red, blue and green lines correspond to the COPAS measurements. The increase of particle number concentrations indicates the Geophysica exhaust trail encounters.
Fig. 4. ClO\textsubscript{x} mixing ratios derived from HALOX ClO night measurements in combination with various $K_{eq}$ parametrisations. The red line corresponds to Cl\textsubscript{y} calculated from the correlation with N\textsubscript{2}O (Eq. 3).
The green line shows $k_{\text{diss}}[\text{M}]$ for comparison. $k_{\text{diss}}$ was determined from the recombination constant taken from JPL 2009 and thermal equilibrium constant of Plenge et al. (2005). The number density $[\text{M}]$ was calculated for the mean temperatures and pressure on the track of the Geophysica during the self-match flight.

**Fig. 5.** Photolysis rate based on cross sections i–iv in dependence of solar zenith angle. The green line shows $k_{\text{diss}}[\text{M}]$ for comparison. $k_{\text{diss}}$ was determined from the recombination constant taken from JPL 2009 and thermal equilibrium constant of Plenge et al. (2005). The number density $[\text{M}]$ was calculated for the mean temperatures and pressure on the track of the Geophysica during the self-match flight.
**Fig. 6.** Comparison of the observed rise of ClO concentrations (marked by the purple symbols) with SZA to the integrated ClO produced by simulations employing photolysis rates $J_{\text{ClOOCl}}$ based on i–iv. The measured rise results from the subtraction of ClO measured on the outbound flight leg (grey symbols) from ClO measured on the inbound flight leg (shown as black symbols with the thicker ones representing strict match pairs). Grey and yellow areas show the measurement uncertainties.
Fig. 7. Comparison of HALOX ClO mixing ratio (black points with grey uncertainty areas) with the results of CLaMS simulations using various combinations of kinetic parameters. The “good matches” are marked with big symbols.