

**FLEXPART
description**

A. Stohl et al.

Technical note: The Lagrangian particle dispersion model FLEXPART version 6.2

A. Stohl¹, C. Forster¹, A. Frank², P. Seibert², and G. Wotawa³

¹Norwegian Institute of Air Research, Kjeller, Norway

²Institute of Meteorology, University of Natural Resources and Applied Life Sciences, Vienna, Austria

³Preparatory Commission for the Comprehensive Nuclear Test Ban Treaty Organization, Vienna, Austria

Received: 21 April 2005 – Accepted: 7 June 2005 – Published: 13 July 2005

Correspondence to: A. Stohl (ast@nilu.no)

© 2005 Author(s). This work is licensed under a Creative Commons License.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

Abstract

The Lagrangian particle dispersion model FLEXPART was originally (about 8 years ago) designed for calculating the long-range and mesoscale dispersion of air pollutants from point sources, such as after an accident in a nuclear power plant. In the meantime FLEXPART has evolved into a comprehensive tool for atmospheric transport modeling and analysis. Its application fields were extended from air pollution studies to other topics where atmospheric transport plays a role (e.g., exchange between the stratosphere and troposphere, or the global water cycle). It has evolved into a true community model that is now being used by at least 25 groups from 14 different countries and is seeing both operational and research applications. A user manual has been kept actual over the years and was distributed over an internet page along with the model's source code. However, so far there was no citeable description of FLEXPART. In this note we provide a description of FLEXPART's latest version (6.2).

1. Introduction

Lagrangian particle models compute trajectories of a large number of so-called particles (not necessarily representing real particles, but infinitesimally small air parcels) to describe the transport and diffusion of tracers in the atmosphere. The main advantage of Lagrangian models is that, unlike in Eulerian models, there is no numerical diffusion. Furthermore, in Eulerian models a tracer released from a point source is instantaneously mixed within a grid box, whereas Lagrangian models are independent of a computational grid and have, in principle, infinitesimally small resolution.

The basis for current atmospheric particle models was laid by Thomson (1987), who stated the criteria that must be fulfilled in order for a model to be theoretically correct. A monograph on the theory of stochastic Lagrangian models was published by Rodean (1996) and another good review was written by Wilson and Sawford (1996). The theory of modeling dispersion backward in time with Lagrangian particle models

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

was developed by [Flesch et al. \(1995\)](#) and [Seibert and Frank \(2004\)](#). Reviews of the more practical aspects of particle modeling were provided by [Zannetti \(1992\)](#) and [Uliasz \(1994\)](#).

This note describes FLEXPART, a Lagrangian particle dispersion model that simulates the long-range and mesoscale transport, diffusion, dry and wet deposition, and radioactive decay of tracers released from point, line, area or volume sources. It can also be used in a domain-filling mode where the entire atmosphere is represented by particles of equal mass. FLEXPART can be used forward in time to simulate the dispersion of tracers from their sources, or backward in time to determine potential source contributions for given receptors. The management of input data was largely taken from FLEXTRA, a kinematic trajectory model ([Stohl et al., 1995](#)). FLEXPART's first version was developed during the first author's military service at the nuclear-biological-chemical school of the Austrian Forces, the deposition code was added soon later (version 2), and this version was validated using data from three large tracer experiments ([Stohl et al., 1998](#)). Version 3 saw performance optimizations and the development of a density correction ([Stohl and Thomson, 1999](#)). Further updates included the addition of a convection scheme ([Seibert et al., 2001](#)) (version 4), better backward calculation capabilities ([Seibert and Frank, 2004](#)), and improvements in the input/output handling (version 5). Validation was done during intercontinental air pollution transport studies ([Stohl and Trickl, 1999](#); [Forster et al., 2001](#); [Spichtinger et al., 2001](#); [Stohl et al., 2002, 2003](#)). The most recent version described here is 6.2, which saw corrections to the numerics in the convection scheme, the addition of a domain-filling option, and the possibility to use output nests.

FLEXPART is coded following the Fortran 77 standard and tested with several compilers (gnu, Absoft, Portland Group) under a number of operating systems (Linux, Solaris, etc.). The code is carefully documented and optimized for run-time performance. No attempts have been made to parallelize the code because the model is strictly linear and, therefore, it is most effective to partition problems such that they run on single processors and to combine the results if needed.

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART's source code and a manual are freely available from the internet page <http://zardoz.nilu.no/~andreas/flextra+flexpart.html>. According to a recent user survey, at least 25 groups from 14 countries are currently using FLEXPART. The version of FLEXPART described here is based on model level data of the numerical weather prediction model of the European Centre for Medium-Range Weather Forecasts (ECMWF). Other users have developed FLEXPART versions using input data from a suite of different global (e.g., from the National Centers of Environmental Prediction) and meso-scale (e.g., MM5) models, some of which are available from the FLEXPART website but are not described here.

2. Input data and grid definitions

FLEXPART is an off-line model that uses meteorological fields (analyses or forecasts) in Gridded Binary (GRIB) format from the ECMWF numerical weather prediction model (ECMWF, 1995) on a latitude/longitude grid and on native ECMWF model levels as input. The data can be retrieved from the ECMWF archives using a pre-processor that is also available from the FLEXPART website but not described here. The GRIB decoding software is *not* provided with FLEXPART but is publicly available (see links on the FLEXPART website). The data can be global or only cover a limited area. Furthermore, higher-resolution domains can be nested into a mother domain.

The file `includepar` contains all relevant FLEXPART parameter settings, both physical constants and maximum field dimensions. As the memory required by FLEXPART is determined by the various field dimensions, it is recommended that they are adjusted to actual needs before compilation. The file `includecom` defines all FLEXPART global variables and fields, i.e., those shared between most subroutines.

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

2.1. Input data organisation

A file `pathnames` must exist in the directory where FLEXPART is started. It must contain at least four lines:

1. line: Directory where all the FLEXPART command files are stored.
2. line: Directory to which the model output is written.
3. line: Directory where the GRIB input fields are located.
4. line: Path name of the `AVAILABLE` file (see below).

If nests with higher-resolution input data shall also be used, lines 3 and 4 must be repeated for every nest, thus also specifying the nesting level order. Any number of nesting levels can be used up to a maximum (parameter `maxnests`).

The meteorological input data must be organised such that all data for a domain and a certain date must be contained in a single GRIB file. The `AVAILABLE` file lists all available dates and the corresponding file names. For each nesting level, the input files must be stored in a different directory and the `AVAILABLE` file must contain the same dates as for the mother domain. Given a certain particle position, the last (i.e., innermost) nest is checked first whether it contains the particle or not. If the particle resides in this nest, the meteorological data from this nest is interpolated linearly to the particle position. If not, the next nest is checked, and so forth until the mother domain is reached. There is no nesting in the vertical direction and the poles must not be contained in any nest.

The maximum dimensions of the meteorological fields are specified by the parameters `nxmax`, `nymax`, `nuvzmax`, `nwzmax`, `nzmax` in file `includepar`, for x, y, and three z dimensions, respectively. The three z dimensions are for the original ECMWF data (`nuvzmax`, `nwzmax` for model half levels and model levels, respectively) and transformed data (`nzmax`, see below), respectively. The horizontal dimensions of the nests must be smaller than the parameters `nxmaxn`, `nymaxn`. Grid dimensions and other basic things are checked in routine `gridcheck.f`, and error messages are issued if necessary.

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

**FLEXPART
description**

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

The longitude/latitude range of the mother grid is also used as the computational domain. All internal FLEXPART coordinates run from the western/southern domain boundary with coordinates (0,0) to the eastern/northern boundary with coordinates (nx-1,ny-1), where (nx,ny) are the mother grid dimensions. For global input data, FLEXPART repeats the westernmost grid cells at the easternmost domain “boundary”, in order to facilitate interpolation on all locations of the globe (e.g., if input data run from 0 to 359° with 1° resolution, 0° data are repeated at 360°). A global mother domain can be shifted by nxshift (file includepar) data columns (subroutines shift_field.f and shift_field_0.f) if nested input fields would otherwise overlap the “boundaries”. For instance, a domain stretching from 320° to 30° can be nested into the mother grid of the above example by shifting the mother grid by 30°.

2.2. Vertical model structure and required data

FLEXPART needs five 3-D fields: horizontal and vertical wind components, temperature and specific humidity. Input data must be on ECMWF model (i.e. η) levels which are defined by a hybrid coordinate system. The conversion from η to pressure coordinates is given by $p_k = A_k + B_k p_s$ and the heights of the η surfaces are defined by $\eta_k = A_k / \rho_0 + B_k$, where η_k is the value of η at the k th model level, p_s is the surface pressure and $\rho_0 = 101\,325$ Pa. A_k and B_k are coefficients, chosen such that the levels closest to the ground follow the topography, while the highest levels coincide with pressure surfaces; intermediate levels transition between the two. The vertical wind in hybrid coordinates is calculated mass-consistently from spectral data by the pre-processor. A surface level is defined in addition to the regular η levels. 2 m temperature, 10 m winds and specific humidity from the first regular model level are assigned to this level, to represent “surface” values.

Parameterized random velocities in the atmospheric boundary layer (ABL, see Sects. 3 and 4) are calculated in units of m s^{-1} , and not in η coordinates. Therefore, in order to avoid time-consuming coordinate transformations every time step, all 3-D data are interpolated linearly from the ECMWF model levels to terrain-following Cartesian

z coordinates (subroutine `verttransform.f`).

FLEXPART also needs the two-dimensional fields: surface pressure, total cloud cover, 10 m horizontal wind components, 2 m temperature and dew point temperature, large scale and convective precipitation, sensible heat flux, solar radiation, east/west and north/south surface stress, topography, land-sea-mask and subgrid standard deviation of topography. A landuse inventory (for Europe, the data set of [Velde et al. \(1994\)](#) is used) must be provided in an extra file (`landuse.asc`).

3. Physical parameterization of boundary layer parameters

Accumulated surface sensible heat fluxes and surface stresses are available from ECMWF forecasts. The pre-processor selects the appropriate short-term forecasts from the ECMWF archives and deaccumulates the flux data. The total surface stress is computed from

$$\tau = \sqrt{\tau_1^2 + \tau_2^2}, \quad (1)$$

where τ_1 and τ_2 are the surface stresses in east/west and north/south direction, respectively. Friction velocity is then calculated in subroutine `scalev.f` as

$$u_* = \sqrt{\tau/\rho}, \quad (2)$$

where ρ is the air density ([Wotawa et al., 1996](#)). Friction velocities and heat fluxes calculated using this method are most accurate ([Wotawa and Stohl, 1997](#)). However, if deaccumulated surface stresses and surface sensible heat fluxes are not available, the profile method after [Berkowicz and Prahm \(1982\)](#) (subroutine `pbl_profile.f`) is applied to wind and temperature data at the second model level and at 10 m (for wind) and 2 m (for temperature) (note that previously the first model level was used; as

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

ECMWF has its first model level now close to 10 m, the second level is used instead). The following three equations are solved iteratively:

$$u_* = \frac{\kappa \Delta u}{\ln \frac{z_l}{10} - \Psi_m\left(\frac{z_l}{L}\right) + \Psi_m\left(\frac{10}{L}\right)}, \quad (3)$$

$$\Theta_* = \frac{\kappa \Delta \Theta}{R \left[\ln \frac{z_l}{2} - \Psi_h\left(\frac{z_l}{L}\right) + \Psi_h\left(\frac{2}{L}\right) \right]}, \quad (4)$$

$$L = \frac{\bar{T} u_*^2}{g \kappa \Theta_*}, \quad (5)$$

where κ is the von Kármán constant (0.4), z_l is the height of the second model level, Δu is the difference between wind speed at the second model level and at 10 m, $\Delta \Theta$ is the difference between potential temperature at the second model level and at 2 m, Ψ_m and Ψ_h are the stability correction functions for momentum and heat (Businger et al., 1971; Beljaars and Holtslag, 1991), g is the acceleration due to gravity, Θ_* is the temperature scale and \bar{T} is the average surface layer temperature (taken as T at the first model level). The heat flux is then computed by

$$\overline{(w'\Theta'_v)}_0 = -\rho c_p u_* \Theta_*, \quad (6)$$

where ρc_p is the specific heat capacity of air at constant pressure.

ABL heights are calculated according to Vogelezang and Holtslag (1996) using the critical Richardson number concept (subroutine `richardson.f`). The ABL height h_{mix} is set to the height of the first model level l for which the Richardson number

$$Ri_l = \frac{(g/\Theta_{v1})(\Theta_{vl} - \Theta_{v1})(z_l - z_1)}{(u_l - u_1)^2 + (v_l - v_1)^2 + 100u_*^2}, \quad (7)$$

exceeds the critical value of 0.25. Θ_{v1} and Θ_{vl} are the virtual potential temperatures, z_1 and z_l are the heights of, and (u_1, v_1) , and (u_l, v_l) are the wind components at the

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

1st and l th model level, respectively. The formulation of Eq. (7) can be improved for convective situations by replacing Θ_{v1} with

$$\Theta'_{v1} = \Theta_{v1} + 8.5 \frac{\overline{(w'\Theta'_v)_0}}{w_* c_p}, \quad (8)$$

where

$$w_* = \left[\frac{\overline{(w'\Theta'_v)_0} g h_{mix}}{\Theta_{v1} c_p} \right]^{1/3} \quad (9)$$

is the convective velocity scale. The second term on the right hand side of Eq. (8) represents a temperature excess of rising thermals. As w_* is unknown beforehand, h_{mix} and w_* are calculated iteratively.

Spatial and temporal variations of ABL heights on scales not resolved by the ECMWF model play an important role in determining the thickness of the layer over which tracer is effectively mixed. The height of the convective ABL reaches its maximum value (say 1500 m) in the afternoon (say, at 17:00 LST), before a much shallower stable ABL forms. Now, if meteorological data are available only at 12:00 and 18:00 LST and the ABL heights at those times are, say, 1200 m and 200 m, and linear interpolation is used, the ABL height at 17:00 LST is significantly underestimated (370 m instead of 1500 m). If tracer is released at the surface shortly before the breakdown of the convective ABL, this would lead to a serious overestimation of the surface concentrations (a factor of four in the above example). Similar arguments hold for spatial variations of ABL heights due to complex topography and variability in landuse or soil wetness (Hubbe et al., 1997). The thickness of a tracer cloud traveling over such a patchy surface would be determined by the maximum rather than by the average ABL height.

In FLEXPART a somewhat arbitrary parameterization is used to avoid a significant bias in the tracer cloud thickness and the surface tracer concentrations. To account for

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

spatial variations induced by topography, we use an “envelope” ABL height

$$H_{env} = h_{mix} + \min \left[\sigma_z, \max \left(c \frac{V}{N}, 0 \right) \right]. \quad (10)$$

Here, σ_z is the ECMWF model subgrid topography, c is a constant (here: 2.0), V is wind speed, N is the Brunt-Vaisala frequency, and $\frac{V}{N}$ is the local Froude number. H_{env} rather than h_{mix} is used for all subsequent calculations. In addition, H_{env} is not interpolated to the particle position, but the maximum H_{env} of the grid points surrounding a particle's position in space and time is used.

4. Particle transport and diffusion

4.1. Particle trajectory calculations

FLEXPART generally uses the simple “zero acceleration” scheme

$$\mathbf{X}(t + \Delta t) = \mathbf{X}(t) + \mathbf{v}(\mathbf{X}, t)\Delta t, \quad (11)$$

which is accurate to the first order, to integrate the trajectory equation (Stohl, 1998)

$$\frac{d\mathbf{X}}{dt} = \mathbf{v}[\mathbf{X}(t)], \quad (12)$$

with t being time, Δt the time increment, \mathbf{X} the position vector, and $\mathbf{v} = \bar{\mathbf{v}} + \mathbf{v}_t + \mathbf{v}_m$ the wind vector that is composed of the grid scale wind $\bar{\mathbf{v}}$, the turbulent wind fluctuations \mathbf{v}_t and the mesoscale wind fluctuations \mathbf{v}_m .

Since FLEXPART version 5.0, numerical accuracy has been improved by making one iteration of the Petterssen (1940) scheme (which is accurate to the second order) whenever this is possible, but only for the grid-scale winds. It is implemented as a correction applied to the position obtained with the “1zero acceleration” scheme. In three cases it cannot be applied. First, the Petterssen scheme needs winds at a second

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[◀](#)
[▶](#)
[◀](#)
[▶](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Print Version](#)
[Interactive Discussion](#)

FLEXPART
description
A. Stohl et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU

time which may be outside the time interval of the two wind fields kept in memory. Second, if a particle crosses the boundaries of nested domains, and third in the ABL if $ct1 > 0$ (see below).

Particle transport and turbulent dispersion are handled by the subroutine `advance.f` where calls are issued to procedures that interpolate winds and other data to the particle position and the Langevin equations (see below) are solved. The poles are singularities on a latitude/longitude grid. Thus, horizontal winds (variables uu, vv) poleward of latitudes (`switchnorth, switchsouth`) are transformed to a polar stereographic projection (variables $uupol, vvpol$) on which particle advection is calculated. As $uupol, vvpol$ are also stored on the latitude/longitude grid, no additional interpolation is made.

4.2. The Langevin equation

Turbulent motions \mathbf{v}_t for wind components i are parameterized assuming a Markov process based on the Langevin equation (Thomson, 1987)

$$dv_{t_i} = a_i(\mathbf{x}, \mathbf{v}_t, t)dt + b_{ij}(\mathbf{x}, \mathbf{v}_t, t)dW_j, \quad (13)$$

where the drift term a and the diffusion term b are functions of the position, the turbulent velocity and time. dW_j are incremental components of a Wiener process with mean zero and variance dt , which are uncorrelated in time (Legg and Raupach, 1982). Cross-correlations between the different wind components are also not taken into account, since they have little effect for long-range dispersion (Ullasz, 1994).

Gaussian turbulence is assumed in FLEXPART, which is strictly valid only for stable and neutral conditions. Under convective conditions, when turbulence is skewed and larger areas are occupied by downdrafts than by updrafts, this assumption is violated, but for transport distances where particles are rather well mixed throughout the ABL, the error is minor.

FLEXPART
description

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

With the above assumptions, the Langevin equation for the vertical wind component w can be written as

$$dw = -w \frac{dt}{\tau_{L_w}} + \frac{\partial \sigma_w^2}{\partial z} dt + \frac{\sigma_w^2}{\rho} \frac{\partial \rho}{\partial z} dt + \left(\frac{2}{\tau_{L_w}} \right)^{1/2} \sigma_w dW, \quad (14)$$

where w and σ_w are the turbulent vertical wind component and its standard deviation, τ_{L_w} is the Lagrangian timescale for the vertical velocity autocorrelation and ρ is density. The second and the third term on the right hand side are the drift correction (McNider et al., 1988) and the density correction (Stohl and Thomson, 1999), respectively. This Langevin equation is identical to the one described by Legg and Raupach (1982), except for the term from Stohl and Thomson (1999) which accounts for the decrease of air density with height.

Alternatively, the Langevin equation can be re-expressed in terms of w/σ_w instead of w (Wilson et al., 1983):

$$d \left(\frac{w}{\sigma_w} \right) = -\frac{w}{\sigma_w} \frac{dt}{\tau_{L_w}} + \frac{\partial \sigma_w}{\partial z} dt + \frac{\sigma_w}{\rho} \frac{\partial \rho}{\partial z} dt + \left(\frac{2}{\tau_{L_w}} \right)^{1/2} dW. \quad (15)$$

This form was shown by Thomson (1987) to fulfill the well-mixed criterion which states that “if a species of passive marked particles is initially mixed uniformly in position and velocity space in a turbulent flow, it will stay that way” (Rodean, 1996). Although the method proposed by Legg and Raupach (1982) violates this criterion in strongly inhomogeneous turbulence, their formulation was found to be practical, as numerical experiments have shown that it is more robust against an increase in the integration time step. Therefore, Eq. (14) is used with long time steps (see Sect. 4.3); otherwise, Eq. (15) is used. For the horizontal wind components, the Langevin equation is identical to Eq. (14), with no drift and density correction terms.

For the discrete time step implementation of the above Langevin equations (at the example of Eq. (15)), two different methods are used. When $(\Delta t/\tau_{L_w}) < 0.5$,

$$\left(\frac{w}{\sigma_w}\right)_{k+1} = \left(1 - \frac{\Delta t}{\tau_{L_w}}\right) \left(\frac{w}{\sigma_w}\right)_k + \frac{\partial \sigma_w}{\partial z} \Delta t + \frac{\sigma_w}{\rho} \frac{\partial \rho}{\partial z} \Delta t + \left(\frac{2\Delta t}{\tau_{L_w}}\right)^{1/2} \zeta, \quad (16)$$

where ζ is a normally distributed random number with mean zero and unit standard deviation. The subscripts k and $k+1$ refer to subsequent times separated by Δt . When $(\Delta t/\tau_{L_w}) \geq 0.5$,

$$\left(\frac{w}{\sigma_w}\right)_{k+1} = r_w \left(\frac{w}{\sigma_w}\right)_k + \frac{\partial \sigma_w}{\partial z} \tau_{L_w} (1 - r_w) + \frac{\sigma_w}{\rho} \frac{\partial \rho}{\partial z} \tau_{L_w} (1 - r_w) + (1 - r_w^2)^{1/2} \zeta, \quad (17)$$

where $r_w = \exp(-\Delta t/\tau_{L_w})$ is the autocorrelation of the vertical wind. When a particle reaches the surface or the top of the ABL, it is reflected and the sign of the turbulent velocity is changed (Wilson and Flesch, 1993).

4.3. Determination of the time step

FLEXPART can be used in two different modes. The computationally faster one (`ctl < 0` in file `COMMAND`) does not adapt the computation time step to τ_L and FLEXPART uses constant time steps of one synchronisation time interval (`lsync_time`, specified in file `COMMAND`, typically 900 seconds). Usually, autocorrelations are very low in this mode and turbulence is not described well. Nevertheless, for large scale applications FLEXPART works very well with this option (Stohl et al., 1998). If turbulence shall be described more accurately, the time steps must be limited by τ_L . Since the vertical wind is most important, only τ_{L_w} is used for this. The user must specify two constants, `ctl` and `ifine` in file `COMMAND`. The first one determines the time step Δt_i , according to

$$\Delta t_i = \frac{1}{c_{tl}} \min \left(\tau_{L_w}, \frac{h}{2w}, \frac{0.5}{\partial \sigma_w / \partial z} \right). \quad (18)$$

The minimum value of Δt_j is 1 second. Δt_j is used for solving the Langevin equations for the horizontal turbulent wind components.

For solving the Langevin equation for the vertical wind component, a shorter time step $\Delta t_w = \Delta t_j / \text{ifine}$ is used. However, note that there is no interaction between horizontal and vertical wind components on timescales less than Δt_j . This strategy (given sufficiently large values for `ctl` and `ifine`) ensures that the particles stay vertically well-mixed also in very inhomogeneous turbulence, while keeping the computational cost at a minimum.

4.4. Parameterization of the wind fluctuations

For σ_{v_i} and τ_{L_i} Hanna (1982) proposed a parameterization scheme based on the boundary layer parameters h , L , w_* , z_0 and u_* , i.e. ABL height, Monin-Obukhov length, convective velocity scale, roughness length and friction velocity, respectively. It is used in subroutines `hanna.f`, `hanna1.f`, `hanna_short.f` with a modification taken from Ryall et al. (1997) for σ_w , as Hanna's scheme does not always yield smooth profiles of σ_w throughout the whole convective ABL. In the following, subscripts u and v refer to the along-wind and the cross-wind components (transformed to grid coordinates in subroutine `windalign.f`), respectively, and w to the vertical component of the turbulent velocities; f is the Coriolis parameter. The minimum τ_{L_u} , τ_{L_v} and τ_{L_w} used are 10 s, 10 s and 30 s, respectively, in order to avoid excessive computation times for particles close to the surface.

Unstable conditions:

$$\frac{\sigma_u}{u_*} = \frac{\sigma_v}{u_*} = \left(12 + \frac{h}{2|L|} \right)^{1/3} \quad (19)$$

$$\tau_{L_u} = \tau_{L_v} = 0.15 \frac{h}{\sigma_u} \quad (20)$$

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[◀](#)
[▶](#)
[◀](#)
[▶](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Print Version](#)
[Interactive Discussion](#)

EGU

$$\frac{\sigma_w}{w_*} = \left[1.2 \left(1 - 0.9 \frac{z}{h} \right) \left(\frac{z}{h} \right)^{2/3} + \left(1.8 - 1.4 \frac{z}{h} \right) u_*^2 \right]^{1/2} \quad (21)$$

For $z/h < 0.1$ and $z - z_0 > -L$:

$$\tau_{L_w} = 0.1 \frac{z}{\sigma_w [0.55 - 0.38(z - z_0)/L]} \quad (22)$$

For $z/h < 0.1$ and $z - z_0 < -L$:

$$5 \quad \tau_{L_w} = 0.59 \frac{z}{\sigma_w} \quad (23)$$

For $z/h > 0.1$:

$$\tau_{L_w} = 0.15 \frac{h}{\sigma_w} \left[1 - \exp\left(\frac{-5z}{h}\right) \right] \quad (24)$$

Neutral conditions:

$$\frac{\sigma_u}{u_*} = 2.0 \exp(-3fz/u_*) \quad (25)$$

$$10 \quad \frac{\sigma_v}{u_*} = \frac{\sigma_w}{u_*} = 1.3 \exp(-2fz/u_*) \quad (26)$$

$$\tau_{L_u} = \tau_{L_v} = \tau_{L_w} = \frac{0.5z/\sigma_w}{1 + 15fz/u_*} \quad (27)$$

Stable conditions:

$$\frac{\sigma_u}{u_*} = 2.0 \left(1 - \frac{z}{h} \right) \quad (28)$$

$$\frac{\sigma_v}{u_*} = \frac{\sigma_w}{u_*} = 1.3 \left(1 - \frac{z}{h} \right) \quad (29)$$

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART
description

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

$$\tau_{L_u} = 0.15 \frac{h}{\sigma_u} \left(\frac{z}{h} \right)^{0.5} \quad (30)$$

$$\tau_{L_v} = 0.07 \frac{h}{\sigma_v} \left(\frac{z}{h} \right)^{0.5} \quad (31)$$

$$\tau_{L_w} = 0.1 \frac{h}{\sigma_w} \left(\frac{z}{h} \right)^{0.5} \quad (32)$$

Lacking suitable turbulence parameterizations above the ABL ($z > h$), a constant vertical diffusivity $D_z = 0.1 \text{ m}^2 \text{ s}^{-1}$ is used in the stratosphere, following recent work of Legras et al. (2003), whereas a horizontal diffusivity $D_h = 50 \text{ m}^2 \text{ s}^{-1}$ is used in the free troposphere. Stratosphere and troposphere are distinguished based on a threshold of 2 pvu (potential vorticity units). Diffusivities are converted into velocity scales using $\sigma_{v_i} = \sqrt{D_i / dt}$.

4.5. Mesoscale velocity fluctuations

Mesoscale motions are neither resolved by the ECMWF data nor covered by the turbulence parameterization. This unresolved spectral interval needs to be taken into account at least in an approximate way, since mesoscale motions can significantly accelerate the growth of a dispersing plume (Gupta et al., 1997). For this, we use a similar method as Maryon (1998), namely to solve an independent Langevin equation for the mesoscale wind velocity fluctuations (“meandering” in Maryon’s terms). Assuming that the variance of the wind at the grid scale provides some information on its subgrid variance, the wind velocity standard deviation used for the mesoscale Langevin equation is set to `turbmesoscale` (set in file `includepar`) times the standard deviation of the grid points surrounding the particle’s position. The corresponding time scale is taken as half the interval at which wind fields are available, assuming that the linear interpolation between the grid points can recover half the subgrid variability, not an

unlikely assumption (Stohl et al., 1995). This empirical approach does not describe actual mesoscale phenomena, but it is similar to the ensemble methods used to assess trajectory accuracy (Kahl, 1996; Baumann and Stohl, 1997; Stohl, 1998).

4.6. Moist Convection

5 An important transport mechanism are the updrafts in convective clouds. They occur in conjunction with downdrafts within the clouds and compensating subsidence in the cloud-free surroundings. These convective transports are grid-scale in the vertical, but sub-grid scale in the horizontal, and are not represented by the ECMWF vertical velocity.

10 To represent convective transport in a particle dispersion model, it is necessary to redistribute particles in the entire vertical column. For FLEXPART we chose the convective parameterization scheme by Emanuel and Živković-Rothman (1999), as it relies on the grid-scale temperature and humidity fields and calculates a displacement matrix providing the necessary mass flux information for the particle redistribution. The convective parameterization is switched on using `lconvection` in file `COMMAND`. It's computation time scales to the square of the number of vertical model levels and may account for up to 70% of FLEXPART's computation time using current 60-level ECMWF data.

15 The convection is computed within the subroutines `convmix.f`, `calcmatrix.f`, `convect43c.f`, and `redist.f`. It is called every FLEXPART `lsynctime` time step (typically 900 s) with time-interpolated temperature and specific humidity profiles from the ECMWF data. Note that the original ECMWF model levels, not the Cartesian coordinates, are used in the convection scheme. For efficiency reasons, particles are sorted according to their horizontal grid positions (`sort2.f`) before calling the convection scheme once per grid column.

25 In the Emanuel scheme (`convect43c.f`), convection is triggered whenever

$$T_{vp}^{LCL+1} \geq T_v^{LCL+1} + T_{thres} \quad (33)$$

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

with T_{vp}^{LCL+1} the virtual temperature of a surface air parcel lifted to the level above the lifting condensation level LCL , T_v^{LCL+1} the virtual temperature of the environment there, and $T_{thres}=0.9\text{ K}$ a threshold temperature value. Based on the buoyancy sorting principle (Emanuel, 1991; Telford, 1975), a matrix MA of the saturated upward and downward mass fluxes within clouds is calculated by accounting for entrainment and detrainment:

$$MA^{i,j} = \frac{M^i (|\sigma^{i,j+1} - \sigma^{i,j}| + |\sigma^{i,j} - \sigma^{i,j-1}|)}{(1 - \sigma^{i,j}) \sum_{j=LCL}^{LNB} [|\sigma^{i,j+1} - \sigma^{i,j}| + |\sigma^{i,j} - \sigma^{i,j-1}|]} \quad (34)$$

Here $MA^{i,j}$ are the mass fractions displaced from level i to level j , M^i the mass fraction displaced from the surface to level i , LNB the level of neutral buoyancy of a surface air parcel and $0 < \sigma^{i,j} < 1$ the mixing fraction between level i and level j . The fraction $\sigma^{i,j}$ is determined by the environmental potential temperature θ^j , the liquid potential temperature $\theta_{lp}^{i,j}$ of air displaced adiabatically from i to j , and the liquid potential temperature $\theta_{lp}^{i,j}$ of an air parcel first lifted adiabatically to level i and further to level j :

$$\sigma^{i,j} = \frac{\theta^j - \theta_{lp}^{i,j}}{\theta_l^{i,j} - \theta_{lp}^{i,j}} \quad (35)$$

By summing up over all levels j , we then calculate the saturated up- and downdrafts at each level i from Eq. (34) and assume that these fluxes are balanced by a subsidence mass flux in the environment.

The particles in each convectively active box are then redistributed (`redist.f`) according to the matrix MA . If the mass of an ECMWF model layer i is m^i and the mass flux from layer i to layer j accumulated over one time step is $\Delta MA^{i,j}$, then the probability of a particle to be moved from layer i to layer j is $\Delta MA^{i,j} / m^i$. Whether a

**FLEXPART
description**

A. Stohl et al.

[Title Page](#)
[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU

given particle is displaced or not is determined by drawing a random number between $[0,1]$, which also determines the position of the particle within the destination layer j . After the convective redistribution of the particles, the compensating subsidence mass fluxes are converted to a vertical velocity acting on those particles in the grid box that are not displaced by convective drafts. By calculating a subsidence velocity rather than displacing particles randomly between layers the scheme's numerical diffusion in the cloud-free environment is eliminated. The scheme was tested and fulfills the well-mixed criterion, i.e., if a tracer is well mixed in the whole atmospheric column, it remains so after the convection.

4.7. Particle splitting

During the initial phase of dispersion from a point source in the atmosphere, particles normally form a compact cloud. Relatively few particles suffice to simulate this initial phase correctly. After some time, however, the particle cloud gets distorted and particles spread over a much larger area. More particles are now needed. FLEXPART allows the user to specify a time constant Δt_s (file `COMMAND`). Particles are split into two (each of which receives half of the mass of the original particle) after travel times of Δt_s , $2\Delta t_s$, $4\Delta t_s$, $8\Delta t_s$, and so on (subroutine `timemanager.f`).

5. Forward and backward modeling

Normally, when FLEXPART is run forward in time (`ldirect=1` in file `COMMAND`), particles are released from one or a number of sources and concentrations are determined downwind on a grid. However, FLEXPART can also run backward in time (`ldirect=-1`), which is more efficient than forward modeling for calculating source-receptor relationships if the number of receptors is smaller than the number of (potential) sources. In the backward mode, particles are released from a receptor location (e.g., a measurement site) and a four-dimensional (3 space dimensions plus time) response function

(sensitivity) to emission input is calculated.

Since this version (6.2) of FLEXPART, the calculation of the source-receptor relationships is generalized for both forward and backward runs, allowing much greater flexibility regarding the input and output units than before. `ind_source` and `ind_receptor` in file `COMMAND` switch between mass and mass mixing ratio units at the source and at the receptor, respectively. Note that source always stands for the physical source and not the location of the particle release, which is done at the source in forward mode but at the receptor in backward mode. Table 1 gives an overview of the units used in forward and backward modeling for different settings of the above switches. A “normal” forward simulation which specifies the release in mass units (i.e., kg) and also samples the output in mass units (i.e., a concentration in ng m^{-3}) requires both switches to be set to 1.

In the backward mode, any value not equal zero can be entered as the release “mass” in file `RELEASES` because the output is normalized by this value. The calculated response function is related to the particles’ residence time in the output grid cells. The unit of the output response function varies, depending on how the switches are set. If `ind_source=1` and `ind_receptor=1`, the response function has the unit s. If this response function is folded (i.e., multiplied) with a 3-D field of emission mass fluxes into the output grid boxes (in $\text{kg m}^{-3}\text{s}^{-1}$), a concentration at the receptor (kg m^{-3}) is obtained. If `ind_source=1` and `ind_receptor=2`, the response function has the unit $\text{s m}^3 \text{kg}^{-1}$ and if it is folded with the emission mass flux (again in $\text{kg m}^{-3}\text{s}^{-1}$), a mass mixing ratio at the receptor is obtained. The units of the response function for `ind_source=2` can be understood in analogy.

In the case of loss processes (dry or wet deposition, decay) the response function is “corrected” for these loss processes. See Seibert (2001) and, particularly, Seibert and Frank (2004) for a description of these generalized in- and output options and the implementation of backward modeling in FLEXPART. Seibert and Frank (2004) also describe the theory of backward modeling and give some examples, and Stohl et al. (2003) presents an application.

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

There are differences between the dimensions of the output fields needed to run FLEXPART forward and backward in time. For backward simulations, the output for all the receptors (particle release locations) must be kept separate, contrary to the forward runs where releases from several source locations can be combined. Therefore, for the backward runs the output fields must contain a dimension `maxpoint`, the maximum number of release points, that is not needed in forward runs. In order to avoid creating a further dimension for the output fields (thus increasing FLEXPART's memory demands), `maxspec`, the maximum number of chemical species used in forward runs, is replaced with `maxpoint` in the backward runs as a dimension of the output fields. This has the disadvantage that only one species (with certain properties regarding the removal processes, see below) can be calculated in a backward run. To switch between forward and backward runs, the parameter `maxpointspec` is used. It must be set (in `includecom`) to `maxspec` for forward runs and to `maxpoint` for backward runs. FLEXPART must be recompiled upon changing this.

6. Plume trajectories

In a recent paper, [Stohl et al. \(2002\)](#) proposed a method to condense the complex and large FLEXPART output using a cluster analysis ([Dorling et al., 1992](#)). The idea behind this is to cluster, at every output time, the positions of all particles originating from a release point, and write out only clustered particle positions, along with additional information (e.g., fraction of particles in the ABL and in the stratosphere). This creates information that is almost as compact as traditional trajectories but accounts for turbulence and convection. This option can be activated by setting `iout` to 4 or 5 in file `COMMAND`. The number of clusters can be set with the parameter `ncluster` in file `includepar`. The clustering is handled and output is produced by subroutine `plumetraj.f`.

7. Removal processes

FLEXPART takes into account radioactive (or other) decay, wet deposition, and dry deposition by reducing a particle's mass. However, as atmospheric transport is the same for all chemical species, a single particle can represent several (up to `maxspec`) chemical species, each affected differently by the removal processes.

7.1. Radioactive decay

Radioactive decay is accounted for by reducing the particle mass according to

$$m(t + \Delta t) = m(t) \exp(-\Delta t / \beta) , \quad (36)$$

where m is particle mass, and the time constant $\beta = T_{1/2} / \ln(1/2)$ is determined from the half life $T_{1/2}$ specified in file `SPECIES`. Deposited pollutant mass decays at the same rate.

7.2. Wet deposition

Wet deposition removes aerosols and gases from the atmosphere. In principle, in-cloud and below-cloud scavenging must be separated (Asman, 1995). However, as data on cloud base height and depth are not available, in-cloud and below-cloud scavenging are treated jointly in FLEXPART. Using scavenging coefficients, wet deposition takes the form of an exponential decay process (McMahon, 1979)

$$m(t + \Delta t) = m(t) \exp(-\Lambda \Delta t) , \quad (37)$$

where m and Λ are the particle mass and the scavenging coefficient, respectively. The scavenging coefficient Λ increases with precipitation rate according to

$$\Lambda = A l^B , \quad (38)$$

where l is the precipitation rate in mm/hour, A [s^{-1}] is the scavenging coefficient at $l=1$ mm/h and B gives the dependency on precipitation rate. Both A and B must be

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART
description

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

specified in file `SPECIES`. FLEXPART uses the same scavenging coefficients for snow and rain.

As wet deposition depends nonlinearly on precipitation rate, subgrid variability of precipitation must be accounted for (Hertel et al., 1995). The area fraction which experiences precipitation given a certain grid-scale precipitation rate is calculated by

$$F = \max \left[0.05, CC \frac{I_l f r_l(I_l) + I_c f r_c(I_c)}{I_l + I_c} \right], \quad (39)$$

where CC is the total cloud cover, I_l and I_c are the large scale and convective precipitation rates, respectively, and $f r_l$ and $f r_c$ are correction factors that depend on I_l and I_c (see Table 2). The subgrid scale precipitation rate is then $I_s = (I_l + I_c) / F$.

7.3. Dry deposition

Dry deposition is described in FLEXPART by a deposition velocity

$$v_d(z) = -F_C / C(z), \quad (40)$$

where F_C and C are the flux and the concentration of a species at height z within the constant flux layer. A constant deposition velocity v_d can be set (file `SPECIES`).

Alternatively, if the physical and chemical properties of a substance are known (file `SPECIES`), more complex parameterizations for gases and particles are also available.

7.3.1. Dry deposition of gases

The deposition velocity of a gas is calculated with the *resistance method* (Wesely and Hicks, 1977) in subroutine `getvdep.f` according to

$$|v_d(z)| = [r_a(z) + r_b + r_c]^{-1}, \quad (41)$$

where r_a is the aerodynamic resistance between z and the surface, r_b is the quasilaminar sublayer resistance, and r_c is the bulk surface resistance.

The aerodynamic resistance r_a is calculated in function `raerod.f` using the flux-profile relationship based on Monin-Obukhov similarity theory (Stull, 1988)

$$r_a(z) = \frac{1}{\kappa u_*} [\ln(z/z_0) - \Psi_h(z/L) + \Psi_h(z_0/L)] . \quad (42)$$

Following Erisman et al. (1994), the quasilaminar sublayer resistance is

$$r_b = \frac{2}{\kappa u_*} \left(\frac{Sc}{Pr} \right)^{2/3} , \quad (43)$$

where Sc and Pr are the Schmidt and Prandtl numbers, respectively. Pr is 0.72 and $Sc = \nu/D_i$, with ν being the kinematic viscosity of air and D_i being the molecular diffusivity of species i in air. The slight dependency of ν on air temperature is formulated in accordance with Pruppacher and Klett (1978). r_b is calculated in function `getrb.f`.

The surface resistance is calculated in function `getrc.f` following Wesely (1989) as

$$r_c = \left[1/(r_s + r_m) + 1/r_{lu} + 1/(r_{dc} + r_{cl}) + 1/(r_{ac} + r_{gs}) \right]^{-1} , \quad (44)$$

where r_s , r_m and r_{lu} represent the bulk values for leaf stomatal, leaf mesophyll and leaf cuticle surface resistances (altogether the upper canopy resistance), r_{dc} represents the gas-phase transfer affected by buoyant convection in canopies, r_{cl} the resistance of leaves, twig, bark and other exposed surfaces in the lower canopy, r_{ac} the resistance for transfer that depends only on canopy height and density, and r_{gs} the resistance for the soil, leaf litter, etc., at the ground. Each of these resistances is parameterized according to the species' chemical reactivity and solubility, the landuse type, and the meteorological conditions. The landuse inventory (Velde et al., 1994) provides the area fractions of eight landuse classes for which roughness lengths z_0 are estimated, on a grid with 10' resolution (Table 3). Charnock's relationship (Stull, 1988) $z_0 = 0.016u_*^2/g$ is used to calculate z_0 for the classes "Ocean" and "Inland water", because of its dependence on wave height. Deposition velocities are calculated for all landuse classes and

**FLEXPART
description**

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

weighted with their respective areas. Outside of Europe, the landuse classes are determined only from the ECMWF land-sea-mask, attributing the landuse classes “Ocean” to the sea surfaces and “Grasslands” to the land surfaces.

7.3.2. Dry deposition of particulate matter

5 The deposition of particulates is calculated in subroutine `partdep.f` according to

$$v_d(z) = [r_a(z) + r_b + r_a(z)r_b v_g]^{-1} + v_g, \quad (45)$$

where v_g is the gravitational settling velocity calculated from (Slinn, 1982)

$$v_g = \frac{g \rho_p d_p^2 C_{cun}}{18 \mu}, \quad (46)$$

10 where ρ_p and d_p are the particle density and diameter, μ the dynamic viscosity of air ($0.000018 \text{ kg m}^{-1} \text{ s}^{-1}$) and C_{cun} the Cunningham slip-flow correction. The quasilaminar sublayer resistance is calculated from the same relationship as for gases, with an additional impaction term. For further details see Slinn (1982).

Settling and dry deposition velocities are strongly dependent on particulate size. FLEXPART assumes a logarithmic normal size distribution of the particulate mass.

15 The user must specify the mean particulate diameter $\overline{d_p}$ and a measure of the variation around $\overline{d_p}$, σ_p . Then, the settling and deposition velocities are calculated for several particle diameters and are weighted with their respective particulate mass fractions.

20 Gravitational settling is important not only for the computation of the dry deposition velocity, but also affects the particle’s trajectory. As a FLEXPART particle can normally represent several species, gravitational settling can only be taken into account correctly (i.e., influence particle trajectories) in single-species simulations.

**FLEXPART
description**

A. Stohl et al.

[Title Page](#)
[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU

7.3.3. Loss of particle mass due to dry deposition

The deposition velocity is calculated for a reference height (parameter `href` in file `includepar`) of 15 m. For all particles below $2h_{ref}$, the mass lost by deposition is calculated by

$$\Delta m(t) = m(t) \left[1 - \exp \left(\frac{-v_d(h_{ref})\Delta t}{2h_{ref}} \right) \right]. \quad (47)$$

8. Calculation of concentrations, uncertainties, age spectra, and mass fluxes

Output quantities C_{T_c} at time T_c (output interval `loutstep` is set in file `COMMAND`) are calculated as time-averages over period $[T_c - \Delta T_c/2, T_c + \Delta T_c/2]$. ΔT_c must be specified (`loutaver`) in file `COMMAND`. To calculate the time-averages, concentrations C_{T_s} at times T_s within $[T_c - \Delta T_c/2, T_c + \Delta T_c/2]$ are sampled at shorter intervals ΔT_s (`loutsample` in file `COMMAND`) and are then divided by the number $N = \frac{\Delta T_c}{\Delta T_s}$ of samples taken:

$$C_{T_c} = \frac{1}{N} \sum_{i=1}^N C_{T_s}. \quad (48)$$

Both ΔT_c and ΔT_s must be multiples of the FLEXPART synchronisation interval (`lsynctime` in file `COMMAND`). The shorter the sampling interval ΔT_s , the more samples are taken and the more accurate are thus the time-averaged concentrations.

8.1. Concentrations, mixing ratios, and emission response functions

The user can choose (`iout` in file `COMMAND`, which must be set to 1 for backward runs) whether concentrations, volume mixing ratios or both shall be produced. We shall use the term “concentration” and particle mass here, but note that the actual units are

determined by the settings of `ind_source` and `ind_receptor`, according to Table 1. The concentration in a grid cell is calculated in subroutine `conccalc.f` by sampling the tracer mass fractions of all particles within the grid cell and dividing by the grid cell volume

$$C_{T_s} = \frac{1}{V} \sum_{i=1}^N (m_i f_i), \quad (49)$$

with V being the grid cell volume, m_i particle mass, N the total number of particles, and f_i the fraction of the mass of particle i attributed to the respective grid cell. This mass fraction is calculated by a uniform kernel with bandwidths $(\Delta x, \Delta y)$, where Δx and Δy are the grid distances on the longitude-latitude output grid. Figure 1 illustrates this:

The particle is located at the center of the shaded rectangle with side lengths $(\Delta x, \Delta y)$. Generally, the shaded area stretches over four grid cells, each of which receives a fraction of the particle's mass equal to the fraction of the shaded area falling within this cell. The uniform kernel is not used during the first 3 h after a particle's release (when the mass is attributed only to the grid cell it resides in), in order to avoid smoothing close to the source.

Wet and dry deposition fields are calculated on the same output grid (subroutines `wetdepokernel.f` and `drydepokernel.f`) and are written to all output grid files. The deposited matter is accumulated over the course of a model run, i.e. it generally increases with model time. However, radioactive decay is calculated also for the deposited matter.

8.2. Uncertainties

The uncertainty of the output is estimated by carrying `nclassunc` classes of particles in the model simulation, and determining the concentration separately for each class (subroutine `conccalc.f`). The standard deviation, calculated from `nclassunc` concentration estimates and divided by $\sqrt{nclassunc}$, is the standard deviation of the mean

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

concentration (subroutine `concoutput.f`), which is also written to the output files for every grid cell. Note that the memory needed for some auxiliary fields increases with `nclassunc` and the number of age classes (see below). It may, thus, be necessary to reduce `nclassunc` for runs with large output grids and age spectra calculations or in the backward mode.

8.3. Age spectra

The age spectra option is switched on using `lagespectra` in file `COMMAND`, with the age classes specified in seconds in file `AGECLASSES`. Concentrations are split into contributions from particles of different age, defined as the time passed since their release. Particles are terminated once they are older than the oldest age class and their storage space is made available to new particles. Therefore, the age spectra option can be used also with a single age class for defining a maximum particle age.

8.4. Parabolic kernel

In addition to the simple uniform kernel method, a computationally demanding parabolic kernel as described in (Uliasz, 1994) can be used to calculate surface concentrations for a limited number of receptor points (age spectra are not available in this case):

$$C_{T_s}(x, y, z = 0) = \sum_{i=1}^N \left[\frac{2m_i K(r_x, r_y, r_z)}{h_{x_i} h_{y_i} h_{z_i}} \right], \quad (50)$$

where h_{x_i} , h_{y_i} and h_{z_i} are the kernel bandwidths which determine the degree of smoothing, $r_x = (X_i - x) / h_{x_i}$, $r_y = (Y_i - y) / h_{y_i}$, $r_z = Z_i / h_{z_i}$ with X_i , Y_i and Z_i being the position of particle i . The kernel bandwidths are a function of the particles' age.

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

8.5. Mass fluxes

Mass flux calculations can be switched on using `iflux` in file `COMMAND`. Mass fluxes are calculated separately for eastward, westward, northward, southward, upward and downward directions and contain both grid-scale and subgrid-scale motions. Mass fluxes are determined for the centerlines of the output grid cells, e.g. vertical fluxes are calculated for motions across the half level of each output cell.

9. Domain-filling option

9.1. General

If `mdomainfill=1` in file `COMMAND` particles are not released at specific locations. Instead, the longitudes and latitudes specified for the first release in the `RELEASES` file are used to set up a global or limited model domain. The particles (number is also taken from `RELEASES`) are then distributed in the model domain proportionally to air density (subroutine `init_domainfill.f`). Each particle receives the same mass, altogether accounting for the total atmospheric mass. Subsequently, particles move freely in the atmosphere.

If a limited domain is chosen, mass fluxes are determined in small grid boxes at the boundary of this domain (boundaries must be at least one grid box away from the boundaries of the meteorological input data). In the grid cells with air flowing into the model domain, mass fluxes are accumulated over time and whenever the accumulated mass exceeds the mass of a particle, a new particle (or more, if required) is released at a randomly chosen position at the boundary of the box (subroutine `boundcond_domainfill.f`). At the outflowing boundaries particles are terminated. Note that, due to the change of mass of the atmosphere in the model domain and due to numerical effects, the number of particles used is not exactly constant throughout the simulation.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

9.2. Stratospheric ozone tracer

If `mdomainfill=2`, the domain-filling option is used to simulate a stratospheric ozone tracer. Upon particle creation, the potential vorticity (PV) at its position is determined by interpolation from the ECMWF data. Particles initially located in the troposphere (PV < `pvcrit` potential vorticity units (pvu), default 2 pvu) are not used. In contrast, stratospheric particles (PV > `pvcrit`) are given a mass according to:

$$M_{O_3} = M_{air} P C 48/29, \quad (51)$$

where M_{air} is the mass of air a particle represents, P is PV in pvu, $C=60 \times 10^{-9}$ pvu⁻¹ is the ozone/PV relationship (Stohl et al., 2000) (parameter `ozonescale`), and the factor 48/29 converts from volume to mass mixing ratio. Particles are then allowed to advect through the stratosphere and into the troposphere according to the winds.

10. Model output

Tracer concentrations and/or mixing ratios (for forward runs), or emission sensitivity response functions (for backward runs) are calculated on a 3-D longitude-latitude grid, defined in file `OUTGRID`, whose domain and resolution can differ from the grid on which meteorological input data are given. Two-dimensional wet and dry deposition fields are calculated over the same spatial domain, and tracer mass fluxes can also be determined on the 3-D grid. Except for the mass fluxes, output can also be produced on one nested output grid with higher horizontal but the same vertical resolution, defined in file `OUTGRID_NEST`. For certain locations, specified in file `RECEPTORS`, concentrations can also be calculated independently from a grid (see below). The time interval (variable `loutstep`) at which output is produced is read in from file `COMMAND`. For every output time, files are created, whose file name ends with the date and time in the format `yyyymmddhhmmss`. A list of all these output times is written to the formatted file `dates`. The dates indicate the ending time of an output sampling interval (see

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

section 8).

10.1. Gridded output

There are several output options in FLEXPART, which can all be selected in file COMMAND. Gridded output fields can be concentrations (files `grid_conc_date`), volume mixing ratios (files `grid_pptv_date`), emission response sensitivity in backward simulations (files `grid_time_date`), or fluxes (files `grid_flux_date`, unit $10^{-12} \text{ kg m}^{-2} \text{ s}^{-1}$ for forward runs). Files `grid_conc_date` are created only in forward runs, whereas files `grid_time_date` are only created in backward runs. Note that the units of the files `grid_conc_date` and `grid_time_date` depend on the settings of the switches `ind_source` and `ind_receptor`, following Table 1. In particular, the units of `grid_conc_date` can also be mass mixing ratios. For forward runs, additional files `grid_pptv_date` can be created, which contain volume mixing ratios for gases. Output files `grid_conc*`, `grid_pptv*`, and `grid_time*` also contain wet and dry deposition fields (unit $10^{-12} \text{ kg m}^{-2}$ in forward mode), and all files contain, for each grid cell, corresponding uncertainties. All these file types share a common header, file `header` produced by subroutine `writeheader.f`, where important information on the model run (start of simulation, grid domain, number and position of vertical levels, age classes, release points, etc.) is stored. In all postprocessing programs, the header must be read in before the actual data files. File names for the output nests follow the same nomenclature as described above, but with `_nest` added (e.g., `header_nest`, or `grid_conc_nest_date`). The output files are written with subroutines `concoutput.f` and `fluxoutput.f`.

FLEXPART output files, except for `dates`, are all binary and often contain many grid cells with zero concentrations (or mixing ratios, fluxes, etc.). Writing out only those cells with non-zero values can produce smaller output than a full grid dump. But in this case the grid indices (note that all three are combined into a single integer number) must also be written out and this produces bigger output than a full grid dump if most grid cells contain non-zero concentrations. Therefore, at every output time and for every

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

output field the more efficient method is determined and used.

10.2. Receptor point output

For a list of points at the surface, concentrations or mixing ratios in forward simulations can be determined with a grid-independent method. This information is written to files
5 `receptor_conc` and `receptor_pptv`, respectively, for all dates of a simulation.

10.3. Particle dump and warm start option

Particle information (3-D position, release time, release point, and release masses for all species) can be written out to files (subroutine `partoutput.f`) either continuously (binary files `partposit_date`), or “only at the end” of a simulation (file
10 `partposit_end`). In both cases output is written every output interval but file `partposit_end` is overwritten upon each new output. If FLEXPART must be terminated, it can be continued later on by reading in files `header` and `partposit_end` produced by the previous run (subroutine `readpartpositions.f`). Such a warm start is done if variable `ipin` is set to 1 in file `COMMAND`.

15 If option `mquasilag` is chosen in file `COMMAND`, particle dumps every output interval are produced in a very compact format by converting the positions to an `integer*2` format (subroutine `partoutput_short.f`). As some accuracy is lost in the conversion, this output is not used for the warm start option. Another difference to the normal particle dump is that every particle gets a unique number, thus allowing postprocessing
20 routines to identify continuous particle trajectories.

10.4. Clustered plume trajectories

Condensed particle output using the clustering algorithm described in Sect. 6 is written to the formatted file `trajectories.txt`. Information on the release points (coordinates, release start and end, number of particles) is written by subroutine
25 `openouttraj.f` to the beginning of file `trajectories.txt`. Subsequently,

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

**FLEXPART
description**

A. Stohl et al.

 Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

plumetraj.f writes out a time sequence of the clustering results for each release point: release point number, time in seconds elapsed since the middle of the release interval, plume centroid position coordinates, various overall statistics (e.g., fraction of particles residing in the ABL and troposphere), and then for each cluster the cluster centroid position, the fraction of particles belonging to the cluster, and the root-mean-square distance of cluster member particles from the cluster centroid.

11. Final remark

In this note, we have described the particle dispersion model FLEXPART version 6.2 with the dual purpose of creating a citeable reference for FLEXPART and providing an actual user manual. In the appendix the various FLEXPART input files are briefly explained and examples are given. As FLEXPART develops this text will be kept actual and will be accessible from the internet site <http://zardoz.nilu.no/~andreas/flextra+flexpart.html>.

Appendix A FLEXPART sample input files

A1 The pathnames file

A file `pathnames` must exist in the directory where FLEXPART is started. It states the pathnames (absolute or relative) of input and output files:

```
/home/as/FLEXPART50/options/
/volc/as/contrace/modelresults/forward/
/volc/windcontrace/
/volc/windcontrace/AVAILABLE
/volc/nested/
/volc/nested/AVAILABLE
=====
```

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

Line 1: path where control files "COMMAND" and "RELEASES" are available
 Line 2: name of directory where output files are generated
 Line 3: path where meteorological fields are available (mother grid)
 5 Line 4: full filename of "AVAILABLE"-file (mother grid)

Subsequent lines:

Line 2n+3: path where meteorological fields are available (nested grid n)
 Line 2n+4: full filename of "AVAILABLE"-file (nested grid n)

10 Line below last pathname must be:

=====

The grids must be arranged such as that the coarse-scale nests
 15 come before the fine-scale nests. Multiple nests of the same
 nesting level are allowed. In that case, the order is arbitrary.

A2 Files in directory windfields

The directory where the meteorological input data are stored, here called windfields
 (/volc/windcontrace/ in the above example pathnames file), contains grib-code
 20 files containing the ECMWF data. All meteorological fields must have the same struc-
 ture, i.e. the same computational domain and the same resolution. An example listing
 of this directory is given below.

```
AVAILABLE      EN01102806      EN01102815
EN01102800      EN01102809      EN01102818
25 EN01102803      EN01102812      EN01102821
```

The file names of the grib-code files and their validation dates and times (in UTC)
 must be listed in the file AVAILABLE. While it is practical to have this file reside in
 30 the same directory as the wind fields, this is no necessity and it can also be located
 elsewhere, as its file name is also given in the pathnames file.

DATE YYYYMMDD	TIME HHMISS	FILENAME	SPECIFICATIONS
20011028	000000	EN01102800	ON DISC
20011028	030000	EN01102803	ON DISC
5 20011028	060000	EN01102806	ON DISC
20011028	090000	EN01102809	ON DISC
20011028	120000	EN01102812	ON DISC
20011028	150000	EN01102815	ON DISC
20011028	180000	EN01102818	ON DISC
10 20011028	210000	EN01102821	ON DISC

Nested wind fields must be stored in one or more different directory/ies, as specified in the `pathnames` file.

A3 Files in directory options

15 The files in `directory options` are used to specify the model run. An example listing of `options` is given below.

AGECLASSES	EMISSION_VARIATION_023.dat	landuse.asc	RELEASES.alternative
COMMAND	EMISSION_VARIATION_025.dat	OUTGRID	RELEASES.reference
COMMAND.alternative	EMISSION_VARIATION_026.dat	OUTGRID_NEST	SPECIES
20 COMMAND.reference	EMISSION_VARIATION_027.dat	RECEPTORS	surfdata.t
EMISSION_VARIATION_008.dat	EMISSION_VARIATION_028.dat	RELEASES	surfdepo.t

A3.1 File COMMAND

25 The most important file is the `COMMAND` file which specifies (1) the simulation direction (either forward or backward), (2) the start and (3) the end time of the simulation, (4) the frequency T_c of the model output, (5) the averaging time ΔT_c of model output, and (6) the intervals ΔT_s at which concentrations are sampled, (7) the time constant for particle splitting Δt_s , (8) the synchronisation interval of FLEXPART, (9) the factor c_{tl} by which

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART description

A. Stohl et al.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[◀](#)
[▶](#)
[◀](#)
[▶](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Print Version](#)
[Interactive Discussion](#)

EGU

the time steps must be smaller than the Lagrangian time scale, and (10) the refinement factor for the time step used for solving the Langevin equation of the vertical component of the turbulent wind. If (9) (c_{tl}) is negative, the Langevin equations are solved with constant time steps according to the synchronisation interval. In that case, the value of (10) is arbitrary. The synchronisation interval is the minimum time interval used by the model for all activities (such as concentration calculations, wet deposition calculations, interpolation of data, mesoscale wind fluctuations or output of data) other than the simulation of turbulent transport and dry deposition (if ($c_{tl}>0$)). Further switches determine (11) whether concentrations, mixing ratios, residence times or plume trajectories (or combinations thereof) are to be calculated, (12) the option of particle position dump either at the end of or continuously during the simulation, (13) on/off of subgrid terrain effect parameterization, (14) on/off of deep convection parameterization, (15) on/off calculation of age spectra, (16) continuation of simulation from previous particle dump, (17) on/off for mass flux calculations and output, (18) on/off for the domain-filling option of FLEXPART, (19) an indicator that determines whether mass or mass mixing ratio units are to be used at the source, (20) an indicator that determines whether mass or mass mixing ratio units are to be used at the receptor, (21) on/off of additional compact dump of the positions of numbered particles, (22) on/off for the use of nested output fields.

Two versions of `COMMAND` may be used, which both can be read in by FLEXPART: the first contains formatted input (i.e., a mask to be filled for the various input options that must be filled in), the second contains largely unformatted input and is recommended for the more experienced FLEXPART user. The following example is for formatted input.

```
*****
*
*      Input file for the Lagrangian particle dispersion model FLEXPART
*
*      Please select your options
*
*****
```

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

```

1.  ___          3X, I2
    1
    LDIRECT      1 FOR FORWARD SIMULATION, -1 FOR BACKWARD SIMULATION

5  2.  _____ 3X, I8, 1X, I6
    20040626 000000
    YYYYMMDD HHMISS BEGINNING DATE OF SIMULATION

3.  _____ 3X, I8, 1X, I6
10  20040816 120000
    YYYYMMDD HHMISS ENDING DATE OF SIMULATION

4.  _____ 3X, I5
    7200
15  SSSSS      OUTPUT EVERY SSSSS SECONDS

5.  _____ 3X, I5
    7200
20  SSSSS      TIME AVERAGE OF OUTPUT (IN SSSSS SECONDS)

6.  _____ 3X, I5
    900
    SSSSS      SAMPLING RATE OF OUTPUT (IN SSSSS SECONDS)

25 7.  _____ 3X, I9
    999999999
    SSSSSSSSS TIME CONSTANT FOR PARTICLE SPLITTING (IN SECONDS)

8.  _____ 3X, I5
30  900
    SSSSS      SYNCHRONISATION INTERVAL OF FLEXPART (IN SECONDS)

9.  ---.---    4X, F6.4
    -5.0
35  CTL        FACTOR, BY WHICH TIME STEP MUST BE SMALLER THAN TL

10. ---        4X, I3
    4
40  IFINE      DECREASE OF TIME STEP FOR VERTICAL MOTION BY FACTOR IFINE

11. -          4X, I1
    3

```

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

```

IOUT          1 CONCENTRATION (RESIDENCE TIME FOR BACKWARD RUNS) OUTPUT,
              2 MIXING RATIO OUTPUT, 3 BOTH,4 PLUME TRAJECT., 5=1+4

12.  -        4X, I1
     2
IPOINT       PARTICLE DUMP: 0 NO, 1 EVERY OUTPUT INTERVAL, 2 ONLY AT END

13.  -        4X, I1
     1
LSUBGRID     SUBGRID TERRAIN EFFECT PARAMETERIZATION: 1 YES, 0 NO

14.  -        4X, I1
     1
LCONVECTION  CONVECTION: 1 YES, 0 NO

15.  -        4X, I1
     0
LAGESPECTRA AGE SPECTRA: 1 YES, 0 NO

20.  -        4X, I1
     0
IPIN         CONTINUE SIMULATION WITH DUMPED PARTICLE DATA: 1 YES, 0 NO

17.  -        4X, I1
     0
IFLUX       CALCULATE FLUXES: 1 YES, 0 NO

18.  -        4X, I1
     2
MDOMAINFILL DOMAIN-FILLING TRAJECTORY OPTION: 1 YES, 0 NO, 2 STRAT. O3 TRACER

19.  -        4X, I1
     1
IND_SOURCE  1=MASS UNIT , 2=MASS MIXING RATIO UNIT

20.  -        4X, I1
     1
IND_RECEPTOR 1=MASS UNIT , 2=MASS MIXING RATIO UNIT

40.  -        4X, I1
     0
MQASILAG    QUASILAGRANGIAN MODE TO TRACK INDIVIDUAL PARTICLES: 1 YES, 0 NO
    
```


**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

0=no output, 1 output every output interval, 2 only at end of the simulation

13. Switch on/off subgridscale terrain parameterization (increase of mixing heights due to subgridscale orographic variations)

14. Switch on/off the convection parameterization

15. Switch on/off the calculation of age spectra: if yes, the file AGECLASSES must be available

16. If IPIN=1, a file "partposit_end" from a previous run must be available in the output directory. Particle positions are read in and previous simulation is continued. If IPIN=0, no particles from a previous run are used

17. If IFLUX is set to 1, fluxes of each species through each of the output boxes are calculated. Six fluxes, corresponding to northward, southward, eastward, westward, upward and downward are calculated for each grid cell of the output grid. The control surfaces are placed in the middle of each output grid cell. If IFLUX is set to 0, no fluxes are determined.

18. If MDOMAINFILL is set to 1, the first box specified in file RELEASES is used as the domain where domain-filling trajectory calculations are to be done. Particles are initialized uniformly distributed (according to the air mass distribution) in that domain at the beginning of the simulation, and are created at the boundaries throughout the simulation period.

19. IND_SOURCE switches between different units for concentrations at the source
NOTE that in backward simulations the release of computational particles takes place at the "receptor" and the sampling of particles at the "source".
1=mass units (for bwd-runs = concentration)
2=mass mixing ratio units

20. IND_RECEPTOR switches between different units for concentrations at the receptor
1=mass units (concentrations)
2=mass mixing ratio units

21. MQASILAG indicates whether particles shall be numbered consecutively (1) or with their release location number (0). The first option allows tracking of individual particles using the partposit output files

22. NESTED_OUTPUT decides whether model output shall be made also for a nested output field (normally with higher resolution)

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

A3.2 File OUTGRID

The file OUTGRID specifies the output grid. The maximum allowed number of output levels is set by parameter maxzgrid in file includepar. The maximum dimensions in x and y by parameters maxxgrid and maxygrid.

```
*****
*
*      Input file for the Lagrangian particle dispersion model FLEXPART
*      Please specify your output grid
*
*****
1.  -----
    -10.0000      4X,F11.4
    OUTLONLEFT   GEOGRAFICAL LONGITUDE OF LOWER LEFT CORNER OF OUTPUT GRID
                  (left boundary of the first grid cell -- not its centre)
15
2.  -----
    40.0000      4X,F11.4
    OUTLATLOWER  GEOGRAFICAL LATITUDE OF LOWER LEFT CORNER OF OUTPUT GRID
                  (lower boundary of the first grid cell -- not its centre)
20
3.  -----
    101          4X,I5
    NUMXGRID     NUMBER OF GRID POINTS IN X DIRECTION (= No. of cells + 1)
25
4.  -----
    47           4X,I5
    NUMYGRID     NUMBER OF GRID POINTS IN Y DIRECTION (= No. of cells + 1)
30
5.  -----
    0.5000      4X,F10.3
    DXOUTLON    GRID DISTANCE IN X DIRECTION
35
6.  -----
    0.5000      4X,F10.3
    DYOUTLAT    GRID DISTANCE IN Y DIRECTION
7.  -----
    100.0       4X, F7.1
```

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

```

LEVEL 1          HEIGHT OF LEVEL (UPPER BOUNDARY)
8.  -----.-    4X, F7.1
    300.0
5   LEVEL 2          HEIGHT OF LEVEL (UPPER BOUNDARY)
9.  -----.-    4X, F7.1
    600.0
10  LEVEL 3          HEIGHT OF LEVEL (UPPER BOUNDARY)
10. -----.-    4X, F7.1
    1000.0
    LEVEL 4          HEIGHT OF LEVEL (UPPER BOUNDARY)
15 11. -----.-    4X, F7.1
    2000.0
    LEVEL 5          HEIGHT OF LEVEL (UPPER BOUNDARY)
20 12. -----.-    4X, F7.1
    3000.0
    LEVEL 6          HEIGHT OF LEVEL (UPPER BOUNDARY)

```

In order to define the grid for a nested output field, the file OUTGRID_NEST must exist. It has the same format as file OUTGRID, but does not contain the vertical level information:

```

*****
25 *
*   Input file for the Lagrangian particle dispersion model FLEXPART   *
*   Please specify your output grid                                   *
*                                                                 *
*****
30
1.  -----.-    4X,F11.4
    -125.0000    GEOGRAFICAL LONGITUDE OF LOWER LEFT CORNER OF OUTPUT GRID
    OUTLONLEFT   (left boundary of the first grid cell -- not its centre)
35 2.  -----.-    4X,F11.4
    25.0000     GEOGRAFICAL LATITUDE OF LOWER LEFT CORNER OF OUTPUT GRID
    OUTLATLOWER  (lower boundary of the first grid cell -- not its centre)
3.  -----    4X,I5

```

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

```

1          NUMBER OF GRID POINTS IN X DIRECTION (= No. of cells + 1)
NUMXGRID

4.  -----  4X,I5
5          1          NUMBER OF GRID POINTS IN Y DIRECTION (= No. of cells + 1)
NUMYGRID

5.  -----.-----  4X,F12.5
      0.33333          GRID DISTANCE IN X DIRECTION
DXOUTLON

10
6.  -----.-----  4X,F12.5
      0.25000          GRID DISTANCE IN Y DIRECTION
DYOUTLAT

```

A3.3 File RECEPTORS

RECEPTORS specifies the receptor locations for which the parabolic kernel method shall be applied to calculate air concentrations. The maximum number of receptor sites is set by parameter maxreceptor in file includepar.

```

*****
*
20 *      Input file for the Lagrangian particle dispersion model FLEXPART      *
*      Please specify your receptor points                                  *
*      For the receptor points, ground level concentrations are calculated   *
*
*****
25 1.  -----  4X,A16
      F15          NAME OF RECEPTOR POINT
      RECEPTORNAME

2.  -----.-----  4X,F11.4
      6.1333          GEOGRAFICAL LONGITUDE
30  XRECEPTOR

3.  -----.-----  4X,F11.4
      49.0833          GEOGRAFICAL LATITUDE
35  YRECEPTOR

=====
1.  -----  4X,A16

```

```

NL01          NAME OF RECEPTOR POINT
RECEPTORNAME
5  2.  -----
      5.7833      4X,F11.4
      XRECEPTOR  GEOGRAFICAL LONGITUDE

```

```

3.  -----
      50.9167     4X,F11.4
      YRECEPTOR GEOGRAFICAL LATITUDE

```

A3.4 File RELEASES

RELEASES defines the release specifications. In the first input line, the number *N* of emitted species is defined (1 in the example below). At all locations, the same species must be released. The next *N* input lines give a cross-reference to file SPECIES, where the physical and chemical properties of the released species are given (also the temporal variations of emissions is defined for each species). Then follows a list of release sites (maximum specified by parameter `maxpoint` in file `includepar`), for each of which the release characteristics must be entered: the beginning and the ending time of the release, geographical coordinates of the lower left and upper right corners of the release location, type of vertical coordinate (above ground level, or above sea level), lower level and upper level of source box, the number of particles to be used, and the total mass emitted. Note that the mass entry must be repeated *N* times, one mass per species released. Finally, a name is assigned to each release point.

The particles are released from random locations within a four-dimensional box extending from the lower to the upper level above a rectangle (on a lat/lon grid) defined by the geographical coordinates, and between the release's start and end. With some identical coordinates, line or point sources can be specified, too.

As for COMMAND, the RELEASES file can be provided formatted or unformatted. The example below shows the formatted version.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

```

*****
*
*
*
*   Input file for the Lagrangian particle dispersion model FLEXPART
5 *   Please select your options
*
*
*
*****
10 ++++++
    1
    ___
        i3   Total number of species emitted

    24
    ___
        i3   Index of species in file SPECIES
=====
20 20011028  150007
    ___    ___
        i8,i1,x,i6 Beginning date and time of release

    20011028  150046
    ___    ___
        i8,i1,x,i6 Ending date and time of release

    9.4048
25  ___ . ___
        f9.4 Longitude [DEG] of lower left corner

    48.5060
    ___ . ___
        f9.4 Latitude [DEG] of lower left corner

    9.5067
30  ___ . ___
        f9.4 Longitude [DEG] of upper right corner

    48.5158
    ___ . ___
        f9.4 Latitude [DEG] of upper right corner

35  2
    ___
        i9   1 for m above ground, 2 for m above sea level,
            3 for pressure in hPa

40  6933.60
    ___ . ___
        f10.3 Lower z-level (in m agl or m asl)

```

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

```

6950.40
_____.____
                    f10.3 Upper z-level (in m agl or m asl)
5      20000
_____.____
                    i9      Total number of particles to be released

1.0000E00
_._.____E_
                    e9.4 Total mass emitted
10
FLIGHT_11242
_____ character*40 comment
+++++
20011028  150047
_____.____
                    i8,1x,i6 Beginning date and time of release

20011028  150107
_____.____
                    i8,1x,i6 Ending date and time of release

20      9.3038
_____.____
                    f9.4 Longitude [DEG] of lower left corner

        48.5158
_____.____
                    f9.4 Latitude [DEG] of lower left corner
25
        9.4048
_____.____
                    f9.4 Longitude [DEG] of upper right corner

        48.5906
_____.____
                    f9.4 Latitude [DEG] of upper right corner
30
        2
_____.____
                    i9      1 for m above ground, 2 for m above sea level,
                           3 for pressure in hPa

35      6833.50
_____.____
                    f10.3 Lower z-level (in m agl or m asl)

        6950.40
_____.____
                    f10.3 Upper z-level (in m agl or m asl)
40
        20000
_____

```

```

_____ i9 Total number of particles to be released
1.0000E00
_._.____E_ e9.4 Total mass emitted
5 FLIGHT_11185
_____ character*40 comment
+++++

```

10 A3.5 EMISSION_VARIATION files

Since FLEXPART version 6.0, emission factors can be defined that change the temporal variation of particle releases. This is useful, for instance, to simulate the typical daily and weekly cycle of anthropogenic emissions. The emission factors are given in files EMISSION_VARIATION_nnn.dat, where nnn is the species number defined in file RELEASES. If file EMISSION_VARIATION_nnn.dat does not exist, emission rates for species nnn are taken as constant. Release rates can vary with the hour of the day and with the day of the week, according to the local time at the release location. Emission factors must be 1 on average. 24 hourly as well as 7 daily values must be specified. Furthermore, different disaggregation factors must be given for area sources and for point sources. FLEXPART distinguishes between the two using the lower altitude of the release box: area sources are assumed to start below 0.5m above the ground, whereas point sources are assumed to be higher. Please note that when this option is used, it is not so easy to determine the maximum number of particles present at a particular time of the model run. It might then be necessary to increase the parameter maxpart to a higher value than what would otherwise be needed. The following is an example for an EMISSION_VARIATION_nnn.dat file.

```

30 hr_start nox_area nox_point
0 0.578 0.845 0-1 local time
1 0.491 0.806 1-2 local time
2 0.428 0.786
3 0.329 0.779
4 0.384 0.793

```

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract Introduction

Conclusions References

Tables Figures

◀ ▶

◀ ▶

Back Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART
description

A. Stohl et al.

Title Page

Abstract Introduction

Conclusions References

Tables Figures

◀ ▶

◀ ▶

Back Close

Full Screen / Esc

Print Version

Interactive Discussion

5	0.485	0.832	
6	0.763	0.895	
7	1.103	0.977	
8	1.084	1.031	
5	9	1.047	1.071
10	10	1.096	1.105
11	11	1.196	1.118
12	12	1.298	1.131
13	13	1.357	1.136
10	14	1.447	1.143
15	15	1.565	1.141
16	16	1.636	1.133
17	17	1.662	1.118
18	18	1.401	1.097
15	19	1.168	1.091
20	20	1.031	1.079
21	21	0.926	1.036
22	22	0.816	0.966
23	23	0.709	0.892
			23-24 local time
20	week_day	nox_area	nox_point
	1	1.060	1.000
	2	1.060	1.000
	3	1.060	1.000
	4	1.060	1.000
25	5	1.060	1.000
	6	0.900	1.000
	7	0.800	1.000
			Monday
			Tuesday
			Wednesday
			Thursday
			Friday
			Saturday
			Sunday

A3.6 File AGECLASSES

AGECLASSES provides the times for the age class calculation. In the first data line, the number *n* of age classes is set, and ages are listed in the following *n* lines. The entries specify the end times (in seconds) of the respective intervals to be used, the first one starting at zero seconds. Particles are dropped from the simulation once they exceed the maximum age. Even if no age classes are needed, this option (with the number of age classes set to 1) can be useful to determine the age at which particles are removed from the simulation.

FLEXPART
description

A. Stohl et al.

[Title Page](#)
[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[|◀](#)[▶|](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU

```

*****
*
*Lagrangian particle dispersion model FLEXPART *
5 *       Please select your options       *
*
*This file determines the ageclasses to be used*
*
*Ages are given in seconds. The first class *
10 *starts at age zero and goes up to the first *
*age specified. The last age gives the maximum *
*time a particle is carried in the simulation. *
*
*****
15  6           Integer           Number of age classes
43200          Integer           Age class 1
86400          Integer           Age class 2
129600
172800
20 259200
345600

```

A3.7 Files SPECIES and surfdata.t

SPECIES is a list of species and their physico-chemical properties from which the user can select some for the simulation. Entries are the half life (due to radioactive or chemical decay), wet deposition information (A and B are the factors defined by equation 38), dry deposition information for gases ($D = D_{H_2O} / D_i$, D_{H_2O} is the diffusivity of water vapor and D_i is the diffusivity of the species, H is the effective Henry's constant, and $f0$ varies between 0 and 1 and gives the reactivity of a species relative to that of ozone. For non-reactive species $f0$ is 0, for slightly reactive it is 0.1 and for highly reactive it is 1.),

**FLEXPART
description**

A. Stohl et al.

dry deposition information for particulates (*rho* specifies the density of the substance, *dquer* its mean diameter \bar{d}_p , and *dsig* the measure of variation σ_p). Radioactive decay is switched off by specifying a negative half life, wet deposition is switched off by specifying negative *A*, dry deposition of gases is switched off by negative *D*, dry deposition of particles is switched off by negative *rho*. If no detailed information for deposition velocity calculation is available, a constant deposition velocity *vd* (cm s^{-1}) can be used. Finally, *molweight* gives the molecular weight of the species, which is needed for mixing ratio output.

```

10 *****
*
*      Input file for the Lagrangian particle dispersion model FLEXPART
*      Definition file of chemical species/radionuclides
*
15 *****
      Radioactivity      Wet depo      Dry depo (gases)      Dry depo (particles)      Dry depo
      SPECIES HALF LIFE [s]      A      B      D      H      f0      rho      dquer      dsig      vd      molweight
1  TRACER      -999.9      -9.9E-09
2  O3      -999.9      -9.9E-09      1.5      1.0e-02      1.0      -9.9E09
20 3  NO      -999.9      8.0E-06      0.62      1.2      2.0e-03      0.0      -9.9E09
4  NO2      -999.9      1.0E-05      0.62      1.6      1.0e-02      0.1      -9.9E09
5  HNO3      -999.9      8.0E-04      0.62      1.9      1.0e+14      0.0      -9.9E09
6  HNO2      -999.9      -9.9E-09      1.6      1.0e+05      0.1      -9.9E09
25 7  H2O2      -999.9      1.0E-04      0.62      1.4      1.0e+05      1.0      -9.9E09
8  SO2      -999.9      -9.9E-09      0.62      2.0      1.0e+05      0.0      -9.9E09
9  HCHO      -999.9      -9.9E-09      1.3      6.0e+03      0.0      -9.9E09
10 PAN      -999.9      -9.9E-09      2.6      3.6e+00      0.1      -9.9E09
11 NH3      -999.9      -9.9E-09      1.1      2.0e+14      0.0      -9.9E09
12 SO4-aero      -999.9      1.0E-04      0.80      -9.9
30 13 NO3-aero      -999.9      1.0E-04      0.80      -9.9
14 I2-131      691200.0      8.0E-05      0.62      2.7      1.0e+05      0.1      -9.9E09
15 I-131      691200.0      1.0E-04      0.80      -9.9
16 Cs-137      -999.9      1.0E-04      0.80      -9.9
17 Y-91      5037120.0      1.0E-04      0.80      -9.9
35 18 Ru-106      31536000.0      1.0E-04      0.80      -9.9
19 Kr-85      -999.9      -9.9E-09      -9.9
20 Sr-90      -999.9      1.0E-04      0.80      -9.9
21 Xe-133      198720.0      -9.9E-09      -9.9
22 CO      -999.9      -9.9E-09      -9.9
40 23 NO2TRACER      -999.9      -9.9E-09      -9.9
24 AIRTRACER      -999.9      -9.9E-09      -9.9

```

Title Page

Abstract Introduction

Conclusions References

Tables Figures

◀ ▶

◀ ▶

Back Close

Full Screen / Esc

Print Version

Interactive Discussion

landuse.asc contains the landuse inventory, and surfdata.t, shown below, gives the roughness lengths for each landuse class:

8 landuse categories are related with Leaf Area Index and roughness length

landuse	LAI	z0(m)	Albedo	comment
1	0.50	0.10	0.18	Grassland for agricultural use
2	2.00	0.15	0.10	Arable land
3	3.00	0.30	0.18	Permanent crops
4	7.00	0.60	0.15	Forest
5	0.00	0.10	0.12	Inland water
6	0.20	0.70	0.20	Urban areas
7	1.00	0.10	0.15	Other
8	0.00	0.10	0.12	Ocean

A3.8 File surfdepo.t

surfdepo.t gives the resistances needed for the parameterization of dry deposition of gases for the eight landuse classes and five seasonal categories. This file must not be changed by the user.

```

=====
INPUT RESISTANCES (s/m) FOR THE COMPUTATION OF SURFACE RESISTANCES TO
DRY DEPOSITION
=====
AFTER WESELY, 1989
=====
1 to 8: Landuse types
=====
Values are tabulated for 5 seasonal categories:
1 Midsummer with lush vegetation
2 Autumn with unharvested cropland
3 Late autumn after frost, no snow
4 Winter, snow on ground and subfreezing

```

FLEXPART description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

	1	2	3	4	5	6	7	8		
5	ri	120.	60.	65.	90.	9999.	9999.	150.	9999.	1
	rlu	2000.	2000.	2000.	2000.	9999.	9999.	4000.	9999.	
	rac	100.	200.	365.	2000.	0.	100.	200.	0.	
	rgss	350.	150.	230.	500.	0.	400.	400.	0.	
	rgso	200.	150.	170.	200.	2000.	300.	200.	2000.	
10	rcls	2000.	2000.	2000.	2000.	9999.	9999.	4000.	9999.	
	rclo	1000.	1000.	1000.	1000.	9999.	9999.	1000.	9999.	
	ri	9999.	9999.	9999.	500.	9999.	9999.	9999.	9999.	2
	rlu	9000.	9000.	9000.	5500.	9999.	9999.	9000.	9999.	
15	rac	100.	150.	270.	1710.	0.	100.	140.	0.	
	rgss	350.	200.	285.	500.	0.	400.	400.	0.	
	rgso	200.	150.	170.	200.	2000.	300.	200.	2000.	
	rcls	9000.	9000.	9000.	3270.	9999.	9999.	9000.	9999.	
	rclo	400.	400.	400.	570.	9999.	9999.	400.	9999.	
20	ri	9999.	9999.	9999.	500.	9999.	9999.	9999.	9999.	3
	rlu	9000.	9999.	9470.	5500.	9999.	9999.	9000.	9999.	
	rac	100.	10.	20.	1330.	0.	100.	120.	0.	
	rgss	350.	150.	230.	500.	0.	400.	400.	0.	
25	rgso	200.	150.	170.	200.	2000.	300.	200.	2000.	
	rcls	9000.	9999.	9470.	4500.	9999.	9999.	9000.	9999.	
	rclo	400.	1000.	570.	570.	9999.	9999.	600.	9999.	
	ri	9999.	9999.	9999.	800.	9999.	9999.	9999.	9999.	4
30	rlu	9999.	9999.	9999.	1200.	9999.	9999.	9000.	9999.	
	rac	10.	10.	20.	1330.	0.	100.	50.	0.	
	rgss	100.	100.	100.	100.	0.	100.	50.	0.	
	rgso	3500.	3500.	3500.	3500.	2000.	600.	3500.	2000.	
	rcls	9999.	9999.	9470.	390.	9999.	9999.	9000.	9999.	
35	rclo	1000.	1000.	570.	632.	9999.	9999.	800.	9999.	
	ri	240.	120.	130.	180.	9999.	9999.	300.	9999.	5
	rlu	4000.	4000.	4000.	2670.	9999.	9999.	8000.	9999.	
	rac	80.	50.	100.	1500.	0.	100.	120.	0.	
40	rgss	350.	150.	230.	500.	0.	500.	400.	0.	
	rgso	200.	150.	170.	200.	2000.	300.	200.	2000.	
	rcls	4000.	4000.	4000.	2670.	9999.	9999.	8000.	9999.	

FLEXPART description

A. Stohl et al.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[◀](#)
[▶](#)
[◀](#)
[▶](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Print Version](#)
[Interactive Discussion](#)

References

- Asman, W. A. H.: Parameterization of below-cloud scavenging of highly soluble gases under convective conditions, *Atmos. Environ.*, 29, 1359–1368, 1995. [4760](#)
- Baumann, K. and Stohl, A.: Validation of a long-range trajectory model using gas balloon tracks from the Gordon Bennett Cup 95, *J. Appl. Meteor.*, 36, 711–720, 1997. [4755](#)
- Beljaars, A. C. M. and Holtlag, A. A. M.: Flux parameterization over land surfaces for atmospheric models, *J. Appl. Meteor.*, 30, 327–341, 1991. [4746](#)
- Berkowicz, R. and Prahm, L. P.: Evaluation of the profile method for estimation of surface fluxes of momentum and heat, *Atmos. Environ.*, 16, 2809–2819, 1982. [4745](#)
- Businger, J. A., Wyngaard, J. C., Izumi, Y., and Bradley, E. F.: Flux-profile relationships in the atmospheric surface layer, *J. Atmos. Sci.*, 28, 181–189, 1971. [4746](#)
- Dorling, S. R., Davies, T. D., and Pierce, C. E.: Cluster analysis: a technique for estimating the synoptic meteorological controls on air and precipitation chemistry – method and applications, *Atmos. Environ.*, 26A, 2575–2581, 1992. [4759](#)
- ECMWF: User Guide to ECMWF Products, 2.1. Meteorological Bulletin M3.2, Reading, UK, 1995. [4742](#)
- Emanuel, K. A.: A scheme for representing cumulus convection in large-scale models, *J. Atmos. Sci.*, 48, 2313–2335, 1991. [4756](#)
- Emanuel, K. A. and Živković-Rothman, M.: Development and evaluation of a convection scheme for use in climate models, *J. Atmos. Sci.*, 56, 1766–1782, 1999. [4755](#)
- Erisman, J. W., Van Pul, A., and Wyers, P.: Parametrization of surface resistance for the quantification of atmospheric deposition of acidifying pollutants and ozone, *Atmos. Environ.*, 28, 2595–2607, 1994. [4762](#)
- Flesch, T. K., Wilson, J. D., and Lee, E.: Backward-time Lagrangian stochastic dispersion models and their application to estimate gaseous emissions, *J. Appl. Meteorol.*, 34, 1320–1333, 1995. [4741](#)
- Forster, C., Wandinger, U., Wotawa, G., James, P., Mattis, I., Althausen, D., Simmonds, P., O'Doherty, S., Kleefeld, C., Jennings, S. G., Schneider, J., Trickl, T., Kreipl, S., Jäger, H.,

FLEXPART
description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

and Stohl, A.: Transport of boreal forest fire emissions from Canada to Europe, *J. Geophys. Res.*, 106, 22 887–22 906, 2001. [4741](#)

Gupta, S., McNider, R. T., Trainer, M., Zamora, R. J., Knupp, K., and Singh, M. P.: Nocturnal wind structure and plume growth rates due to inertial oscillations, *J. Appl. Meteor.*, 36, 1050–1063, 1997. [4754](#)

Hanna, S. R.: Applications in air pollution modeling, in: *Atmospheric Turbulence and Air Pollution Modelling*, edited by: Nieuwstadt, F. T. M. and van Dop, H., D. Reidel Publishing Company, Dordrecht, Holland, 1982. [4752](#)

Hertel, O., Christensen, J., Runge, E. H., Asman, W. A. H., Berkowicz, R., Hovmand, M. F., and Hov, O.: Development and testing of a new variable scale air pollution model – ACDEP. *Atmos. Environ.*, 29, 1267–1290, 1995. [4761](#)

Hubbe, J. M., Doran, J. C., Liljegren, J. C., and Shaw, W. J.: Observations of spatial variations of boundary layer structure over the southern Great Plains cloud and radiation testbed, *J. Appl. Meteor.*, 36, 1221–1231, 1997. [4747](#)

Kahl, J. D. W.: On the prediction of trajectory model error, *Atmos. Environ.*, 30, 2945–2957, 1996. [4755](#)

Legg, B. J. and Raupach, M. R.: Markov-chain simulation of particle dispersion in inhomogeneous flows: the mean drift velocity induced by a gradient in Eulerian velocity variance, *Bound.-Layer Met.*, 24, 3–13, 1982. [4749](#), [4750](#)

Legras, B., Joseph, B., and Lefevre, F.: Vertical diffusivity in the lower stratosphere from Lagrangian back-trajectory reconstructions of ozone profiles, *J. Geophys. Res.*, 108, 4562, doi:10.1029/2002JD003045, 2003. [4754](#)

Maryon, R. H.: Determining cross-wind variance for low frequency wind meander, *Atmos. Environ.*, 32, 115–121, 1998. [4754](#)

McMahon, T. A. and Denison, P. J.: Empirical atmospheric deposition parameters – a survey, *Atmos. Environ.*, 13, 571–585, 1979. [4760](#)

McNider, R. T., Moran, M. D., and Pielke, R. A.: Influence of diurnal and inertial boundary layer oscillations on long-range dispersion, *Atmos. Environ.*, 22, 2445–2462, 1988. [4750](#)

Petterssen, S.: *Weather Analysis and Forecasting*, McGraw-Hill Book Company, New York, 221–223, 1940. [4748](#)

Pruppacher, H. R. and Klett, J. D.: *Microphysics of clouds and precipitation*, D. Reidel Publishing Company, Dordrecht, 714, 1978. [4762](#)

Rodean, H.: *Stochastic Lagrangian models of turbulent diffusion*, Meteorological Monographs,

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

26 (48), American Meteorological Society, Boston, USA, 1996. [4740](#), [4750](#)

Ryall, D. B. and Maryon, R. H.: Validation of the UK Met Office's NAME model against the ETEX dataset, in: Model Verification and Emergency Response, edited by: Nodop, K., ETEX Symposium on Long-Range Atmospheric Transport, European Commission EUR 17346, 151–154, 1997. [4752](#)

Seibert, P.: Inverse modelling with a Lagrangian particle dispersion model: application to point releases over limited time intervals, in: Air Pollution Modeling and its Application, edited by: Gryning, S. E. and Schiermeier, F. A., XIV. Proc. of ITM Boulder. New York: Plenum Press, 381–389, 2001. [4758](#)

Seibert, P. and Frank, A.: Source-receptor matrix calculation with a Lagrangian particle dispersion model in backward mode, *Atmos. Chem. Phys.*, 4, 51–63, 2004, [SRef-ID: 1680-7324/acp/2004-4-51](#). [4741](#), [4758](#)

Seibert, P., Krüger, B., and Frank, A.: Parametrisation of convective mixing in a Lagrangian particle dispersion model, Proceedings of the 5th GLOREAM Workshop, Wengen, Switzerland, September 24–26, 2001. [4741](#)

Slinn, W. G. N.: Predictions for particle deposition to vegetative canopies, *Atmos. Environ.*, 16, 1785–1794, 1982. [4763](#)

Spichtinger, N., Wenig, M., James, P., Wagner, T., Platt, U., and Stohl, A.: Satellite detection of a continental-scale plume of nitrogen oxides from boreal forest fires, *Geophys. Res. Lett.*, 28, 4579–4582, 2001. [4741](#)

Stohl, A.: Computation, accuracy and applications of trajectories – a review and bibliography, *Atmos. Environ.*, 32, 947–966, 1998. [4748](#), [4755](#)

Stohl, A., Eckhardt, S., Forster, C., James, P., Spichtinger, N., and Seibert, P.: A replacement for simple back trajectory calculations in the interpretation of atmospheric trace substance measurements, *Atmos. Environ.*, 36, 4635–4648, 2002. [4741](#), [4759](#)

Stohl, A., Forster, C., Eckhardt, S., Spichtinger, N., Huntrieser, H., Heland, J., Schlager, H., Wilhelm, S., Arnold, F., and Cooper, O.: A backward modeling study of intercontinental pollution transport using aircraft measurements, *J. Geophys. Res.*, 108, 4370, doi:10.1029/2002JD002862, 2003. [4741](#), [4758](#)

Stohl, A., Hittenberger, M., and Wotawa, G.: Validation of the Lagrangian particle dispersion model FLEXPART against large scale tracer experiment data, *Atmos. Environ.*, 32, 4245–4264, 1998. [4741](#), [4751](#)

Stohl, A., Spichtinger-Rakowsky, N., Bonasoni, P., Feldmann, H., Memmesheimer, M., Scheel,

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

- H. E., Trickl, T., Hübener, S. H., Ringer, W., and Mandl, M.: The influence of stratospheric intrusions on alpine ozone concentrations, *Atmos. Environ.*, 34, 1323–1354, 2000. [4768](#)
- Stohl, A. and Thomson, D. J.: A density correction for Lagrangian particle dispersion models, *Bound.-Layer Met.*, 90, 155–167, 1999. [4741](#), [4750](#)
- 5 Stohl, A. and Trickl, T.: A textbook example of long-range transport: Simultaneous observation of ozone maxima of stratospheric and North American origin in the free troposphere over Europe, *J. Geophys. Res.*, 104, 30 445–30 462, 1999. [4741](#)
- Stohl, A., Wotawa, G., Seibert, P., and Kromp-Kolb, H.: Interpolation errors in wind fields as a function of spatial and temporal resolution and their impact on different types of kinematic trajectories, *J. Appl. Meteor.*, 34, 2149–2165, 1995. [4741](#), [4755](#)
- 10 Stull, R. B.: *An Introduction to Boundary Layer Meteorology*, Kluwer Academic Publishers, Dordrecht, 1988. [4762](#)
- Telford, J. W.: Turbulence, entrainment and mixing in cloud dynamics, *Pure Appl. Geophys.*, 113, 1067–1084, 1975. [4756](#)
- 15 Thomson D. J.: Criteria for the selection of stochastic models of particle trajectories in turbulent flows, *J. Fluid Mech.*, 180, 529–556, 1987. [4740](#), [4749](#), [4750](#)
- Uliasz, M.: Lagrangian particle dispersion modeling in mesoscale applications, in: *Environmental Modeling*, Vol. II, edited by: Zannetti, P., Computational Mechanics Publications, Southampton, UK, 1994. [4741](#), [4749](#), [4766](#)
- 20 van de Velde, R. J., Faber, W. S., Katwijk van, V. F., Kuylensstierna, J. C. I., Scholten., H. J., Thewessen, T. J. M., Verspuij, M., and Zevenbergen, M.: The Preparation of a European Land Use Database, National Institute of Public Health and Environmental Protection, Report nr 712401001, Bilthoven, The Netherlands, 1994. [4745](#), [4762](#)
- Vogelezang, D. H. P. and Holtslag, A. A. M.: Evaluation and model impacts of alternative boundary-layer height formulations, *Bound.-Layer Met.*, 81, 245–269, 1996. [4746](#)
- 25 Wesely, M. L. and Hicks, B. B.: Some factors that affect the deposition rates of sulfur dioxide and similar gases on vegetation, *J. Air Poll. Contr. Assoc.*, 27, 1110–1116, 1977. [4761](#)
- Wesely, M. L.: Parameterization of surface resistances to gaseous dry deposition in regional-scale numerical models, *Atmos. Environ.*, 23, 1293–1304, 1989. [4762](#)
- 30 Wilson, D. J. and Flesch, T. K.: Flow boundaries in random-flight dispersion models: enforcing the well-mixed condition, *J. Appl. Meteor.*, 32, 1695–1707, 1993. [4751](#)
- Wilson, J. D., Legg, B. J., and Thomson, D. J.: Calculation of particle trajectories in the presence of a gradient in turbulent-velocity scale, *Bound.-Layer Met.*, 27, 163–169, 1983. [4750](#)

Wilson, J. D. and Sawford, B. L.: Review of Lagrangian stochastic models for trajectories in the turbulent atmosphere, *Bound.-Layer Met.*, 78, 191–210, 1996. [4740](#)

Wotawa, G., Stohl, A., and Kromp-Kolb, H.: Parameterization of the planetary boundary layer over Europe – a data comparison between the observation based OML preprocessor and ECMWF model data, *Contr. Atmos. Phys.*, 69, 273–284, 1996. [4745](#)

Wotawa, G. and Stohl, A.: Boundary layer heights and surface fluxes of momentum and heat derived from ECMWF data for use in pollutant dispersion models – problems with data accuracy, in: *The Determination of the Mixing Height – Current Progress and Problems*, edited by: Gryning, S.-E., Beyrich, F., and Batchvarova, E., EURASAP Workshop Proceedings, 1–3 October 1997, Risø National Laboratory, Denmark, 1997. [4745](#)

Zannetti, P.: Particle Modeling and Its Application for Simulating Air Pollution Phenomena, in: *Environmental Modelling*, edited by: Melli, P. and Zannetti, P., Computational Mechanics Publications, Southampton, UK, 1992. [4741](#)

**FLEXPART
description**

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

FLEXPART
description

A. Stohl et al.

Table 1. Physical units of the input (in file RELEASES) and output data for forward (files grid_conc_date) and backward (files grid_time_date) runs for the various settings of the unit switches ind_source and ind_receptor (in both switches 1 refers to mass units, 2 to mass mixing ratio units).

Direction	ind_source	ind_receptor	input unit	output unit
Forward	1	1	kg	ng m ⁻³
Forward	1	2	kg	ppt by mass
Forward	2	1	1	ng m ⁻³
Forward	2	2	1	ppt by mass
Backward	1	1	1	s
Backward	1	2	1	s m ³ kg ⁻¹
Backward	2	1	1	s kg m ⁻³
Backward	2	2	1	s

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Print Version](#)
[Interactive Discussion](#)

EGU

FLEXPART
 description

A. Stohl et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Print Version

Interactive Discussion

EGU

Table 2. Correction factors used for the calculation of the area fraction that experiences precipitation. Precipitation rates are in mm/h.

Factor	I_l and I_c				
	$I \leq 1$	$1 < I \leq 3$	$3 < I \leq 8$	$8 < I \leq 20$	$20 < I$
fr_l	0.50	0.65	0.80	0.90	0.95
fr_c	0.40	0.55	0.70	0.80	0.90

**FLEXPART
description**

A. Stohl et al.

Table 3. List of the landuse classes and roughness lengths used by FLEXPART.

Grassland	0.10
Arable land	0.15
Permanent crops	0.30
Forest	0.60
Inland water	Charnock
Urban areas	0.70
Other	0.10
Ocean	Charnock

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[I◀](#)[▶I](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU

FLEXPART
description

A. Stohl et al.

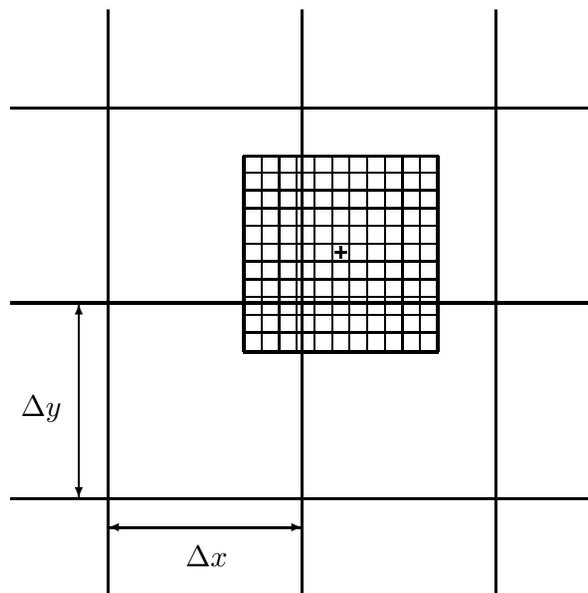


Fig. 1. Illustration of the uniform kernel used to calculate gridded concentration and deposition fields. The particle position is marked by “+”.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Print Version](#)[Interactive Discussion](#)

EGU