

**Performance of the  
line-by-line radiative  
transfer model  
(LBLRTM)**

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**Performance of the line-by-line radiative  
transfer model (LBLRTM) for temperature,  
water vapor, and trace gas retrievals:  
recent updates evaluated with IASI case  
studies**

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## Abstract

Modern data assimilation algorithms depend on accurate infrared spectroscopy in order to make use of the information related to temperature, water vapor ( $\text{H}_2\text{O}$ ), and other trace gases provided by satellite observations. Reducing the uncertainties in our knowledge of spectroscopic line parameters and continuum absorption is thus critical to improving the application of satellite data to weather forecasting. Here we present the results of a rigorous validation of spectroscopic updates to an advanced radiative transfer model, the Line-By-Line Radiative Transfer Model (LBLRTM), against a global dataset of 120 near-nadir, over-ocean, nighttime spectra from the Infrared Atmospheric Sounding Instrument (IASI). We compare calculations from the latest version of LBLRTM (v12.1) to those from a previous version (v9.4+) to determine the impact of spectroscopic updates to the model on spectral residuals as well as retrieved temperature and  $\text{H}_2\text{O}$  profiles. We show that the spectroscopy in the  $\text{CO}_2$   $\nu_2$  and  $\nu_3$  bands is significantly improved in LBLRTM v12.1 relative to v9.4+, and that these spectroscopic updates lead to changes of  $\sim 0.5$  K in the retrieved vertical temperature profiles below 10 hPa, with the sign of the change and the variability among cases depending on altitude. We also find that temperature retrievals using each of these two  $\text{CO}_2$  bands are remarkably consistent in LBLRTM v12.1, potentially allowing these bands to be used to retrieve atmospheric temperature simultaneously. The updated  $\text{H}_2\text{O}$  spectroscopy in LBLRTM v12.1 substantially improves the residuals in the P-branch of the  $\text{H}_2\text{O}$   $\nu_2$  band, while the improvements in the R-branch are more modest. The  $\text{H}_2\text{O}$  amounts retrieved with LBLRTM v12.1 are on average 14 % lower between 100 and 200 hPa, 42 % higher near 562 hPa, and 31 % higher near the surface compared to the amounts retrieved with v9.4+ due to a combination of the different retrieved temperature profiles and the updated  $\text{H}_2\text{O}$  spectroscopy. We also find that the use of a fixed ratio of HDO to  $\text{H}_2\text{O}$  in LBLRTM may be responsible for a significant fraction of the remaining bias in the P-branch of the  $\text{H}_2\text{O}$   $\nu_2$  band. There were no changes to  $\text{O}_3$  spectroscopy between the two model versions, and so both versions gives positive residuals of  $\sim 0.3$  K

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in the R-branch of the  $O_3 \nu_3$  band. While the updates to the  $H_2O$  self continuum employed by LBLRTM v12.1 have clearly improved the match with observations near the  $CO_2 \nu_3$  bandhead, we find that these updates have significantly degraded the match with observations in the fundamental band of CO. Finally, significant systematic residuals remain in the  $\nu_4$  band of  $CH_4$ , but the magnitude of the positive bias in the retrieved mixing ratios is reduced in LBLRTM v12.1, suggesting that the updated spectroscopy could improve retrievals of  $CH_4$  from satellite observations.

## 1 Introduction

Modern data assimilation algorithms for numerical weather prediction (NWP) make extensive use of the information related to temperature, water vapor ( $H_2O$ ), and other trace gases provided by satellite observations. The accuracy of the retrieved vertical profiles of temperature,  $H_2O$ , and other trace gases from satellites depends on the accuracy of the radiative transfer model used in the retrieval. Uncertainties in our knowledge of spectroscopic line parameters and continua are the primary limitations on the accuracy of computed absorption in leading edge radiative transfer models, so reducing these uncertainties is critical to improving the application of satellite data to weather forecasting. Radiance closure studies using high-spectral-resolution infrared radiance measurements allow us to assess the systematic differences between the calculated and measured spectral radiances, and provide a means to assess the consistency of the input spectroscopic parameters within different absorption bands of the same trace gas (e.g. Shephard et al., 2009) and for different gaseous absorbers.

Here we present the results of a rigorous validation of spectroscopic updates to an advanced radiative transfer model, the Line-By-Line Radiative Transfer Model (LBLRTM, Clough et al., 1992, 2005), with respect to a dataset of measurements from the Infrared Atmospheric Sounding Instrument (IASI, Chillon et al., 2001), an infrared spectrometer on the MeteoSat MetOp-A satellite launched in 2006. This study uses a global dataset of 120 clear-sky, nighttime, ocean, near-nadir IASI measurements

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during April 2008 culled from the dataset of Matricardi (2009), allowing the evaluation of the model across a large range of atmospheric conditions.

Section 2 describes the current and past versions of LBLRTM used in this study and Sect. 3 provides an overview of the IASI instrument. Section 4 discusses our radiance closure strategy, including information on the a priori profiles, constraint matrices, and microwindows used in the retrievals of temperature, H<sub>2</sub>O, O<sub>3</sub>, CO, and CH<sub>4</sub>. Section 5 analyzes the results of these radiance closure studies and discusses the impact of the spectroscopic changes on retrieved profiles of temperature and H<sub>2</sub>O.

## 2 LBLRTM

The Line-by-Line Radiative Transfer Model (LBLRTM) is an accurate and flexible radiative transfer model that can be used over the full spectral range from the microwave to the ultraviolet, providing the foundation for many radiative transfer applications (Clough et al., 1992, 2005). LBLRTM has a long and successful heritage at the leading edge of the field, and the model is continually updated and validated against high-resolution spectral measurements (e.g. Payne et al., 2008; Shephard et al., 2009; Delamere et al., 2010; Mlawer et al., 2012). LBLRTM calculations in the thermal infrared are recognized as a reference standard for intercomparisons of radiative transfer models, such as the SPARC Chemistry–Climate Model Validation Activity (CCMVal, Forster et al., 2011) and the Continual Intercomparison of Radiation Codes (CIRC, Oreopoulos and Mlawer, 2010; Oreopoulos et al., 2012).

LBLRTM has been widely used for a number of years as the foundation for retrieval algorithms, including those using measurements from ground based instruments such as the Atmospheric Emitted Radiance Interferometer (AERI, Smith et al., 1999) and satellite instruments like IASI (Amato et al., 2002; Tjemkes et al., 2003) and the Tropospheric Emission Spectrometer (TES, Clough et al., 2006). In addition, LBLRTM has been used to derive the absorption coefficients for the fast radiation codes RRTM and RRTMG (Mlawer et al., 1997; Iacono et al., 2008), which are used for broadband flux

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and heating rate calculations in several General Circulation Models (GCMs) for climate and weather prediction. LBLRTM is also used to train fast radiative transfer models used in NWP assimilation systems, such as the Optical Transmittance model OPTRAN (McMillin et al., 1979) and the Optimal Spectral Sampling (OSS) (Moncet et al., 2008) model implemented in the Joint Center for Satellite Data Assimilation (JCSDA) Community Radiative Transfer Model (CRTM), as well as OPTRAN-Compact, which is used operationally at the National Center for Environmental Prediction (NCEP).

The main features of LBLRTM are described in Clough et al. (2005) and are summarized here. The Voigt line shape is used at all atmospheric levels with an algorithm based on a linear combination of approximating functions. Line coupling in LBLRTM is modeled using a first order perturbation approach (Rosenkranz, 1975). In general, errors associated with the computational procedures in LBLRTM are small – around five times less than those associated with the limiting errors in spectral radiance calculations, which are uncertainties in line parameters and line shape. Spectroscopic parameters in the latest version of LBLRTM (v12.1) come from the HITRAN 2008 line database (Rothman et al., 2009) with a few key exceptions discussed in Sect. 2.2. LBLRTM incorporates the continuum model MT\_CKD (Mlawer et al., 2012), which includes self- and foreign-broadened water vapor continua as well as continua for CO<sub>2</sub>, O<sub>2</sub>, N<sub>2</sub>, O<sub>3</sub>, and extinction due to Rayleigh scattering. Temperature dependent cross section data such as those available with the HITRAN database may be used to treat the absorption due to heavy molecules such as halocarbons.

In this work, we use two versions of LBLRTM. The first is a modified version of LBLRTM v9.4, here called v9.4+, which is described in Sect. 2.1. We chose v9.4+ for this study as this model version was released prior to the recent improvements in CO<sub>2</sub> spectroscopy, including the addition of P- and R-branch line coupling for all CO<sub>2</sub> bands (Shephard et al., 2009). We compare the calculations from this older version of the model to those from the latest version, LBLRTM v12.1. The relevant updates made to LBLRTM, the MT\_CKD continuum, and the associated spectral databases between v9.4 and v12.1 are discussed in Sect. 2.2.

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## 2.1 LBLRTM v9.4+

LBLRTM v9.4 was released in January 2005. This version of the model used v1.0 of AER's line parameter database (hereafter AER v1.0), which was based on HITRAN 2000 (including all updates made before September 2001, see Rothman et al., 2003) with the following exceptions. All line parameters for CH<sub>4</sub> between 922.65 and 1678.33 cm<sup>-1</sup> and for the half-widths, temperature dependence, and pressure shifts of CO were updated from HITRAN 2000 based on data supplied by Linda Brown of the Jet Propulsion Laboratory (personal communication). AER v1.0 also incorporated updated line parameters for O<sub>3</sub> (Wagner et al., 2002), formic acid (HCOOH, Perrin et al., 1999) and for the A-bands of the minor isotopologues of O<sub>2</sub> (Camy-Peyret et al., 2000); these three updates were later incorporated into HITRAN 2004 (Rothman et al., 2005).

Q-branch line coupling was included for CO<sub>2</sub> in LBLRTM v9.4, but only for the main isotopologue. The first-order line coupling parameters in LBLRTM v9.4 for the Q-branches of the bands centered at 618, 667, 720, 721, and 791 cm<sup>-1</sup> were based on Hoke et al. (1989), but were recalculated to be consistent with the CO<sub>2</sub> line parameters in HITRAN 2000. The line coupling parameters for the Q-branches of the bands centered at 1932, 2076, 2093, and 2193 cm<sup>-1</sup> were taken from Strow et al. (1994).

LBLRTM v9.4 used the continuum code MT\_CKD v1.2. The formulation of the H<sub>2</sub>O self and foreign continua in MT\_CKD is discussed in detail in Mlawer et al. (2012). The coefficients for the CO<sub>2</sub> foreign continuum were based on Ridgway et al. (1982), with the coefficients between 0 and 1200 cm<sup>-1</sup> increased by a factor of 7 to match AERI observations made at the U.S. Department of Energy (DoE) Atmospheric Radiation Measurement (ARM) program site at the North Slope of Alaska, as well as other observations made during the ARM//FIRE Water vapor Experiment (AFWEX) (Shephard et al., 2003). The continuum coefficients for the collision-induced fundamental bands of O<sub>2</sub> and N<sub>2</sub> were taken from Thibault et al. (1996) and Lafferty et al. (1996), respectively.

Several improvements to the functionality of LBLRTM were made between v9.4 and v12.1, most notably in the calculations of analytical Jacobians and the number of in-

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strument line shapes included for post-processing the monochromatic spectra. In order to focus on differences between the versions due to spectroscopic parameters rather than these changes, we updated the Fast Fourier Transform (FFT) scan calculations and the analytical Jacobian calculations in LBLRTM v9.4 to match those of LBLRTM v12.1. This code, with spectroscopic parameters equivalent to v9.4 but with the improved features of LBLRTM v12.1, is here called LBLRTM v9.4+.

## 2.2 LBLRTM v12.1

LBLRTM v12.1 was released in November 2011. LBLRTM v12.1 uses v3.1 of the AER line parameter database (hereafter AER v3.1), which is based on the HITRAN 2008 line parameters (Rothman et al., 2009) with exceptions within the range of the IASI instrument for H<sub>2</sub>O, CO<sub>2</sub>, and CH<sub>4</sub>, which are discussed below. The spectroscopic changes between AER v1.0 and AER v3.1 are negligible for CO and O<sub>3</sub>.

The H<sub>2</sub>O line positions and intensities for the range 10 to 2500 cm<sup>-1</sup> in AER v3.1 are from Coudert et al. (2008). The Coudert et al. (2008) line list, as implemented in AER v3.1, included not only parameters for lines that they had measured in the laboratory (for wavenumbers up to 1750 cm<sup>-1</sup>) but also calculated values for lines that had not been measured, in order to provide a linelist that would cover the entire range of the H<sub>2</sub>O ν<sub>2</sub> band in the thermal infrared. Note that while the Coudert et al. measured values were included in the HITRAN 2008 compilation, the HITRAN team had made the decision not to include the calculated values. The impacts of this difference between AER v3.1 and HITRAN 2008 are discussed in Sect. 5.2. The air-broadened half-widths, temperature dependences, and pressure shifts for H<sub>2</sub>O between 350 and 667 cm<sup>-1</sup> are described in Delamere et al. (2010).

The CO<sub>2</sub> line parameters in AER v3.1 were built by starting with the CO<sub>2</sub> line mixing database of Lamouroux et al. (2010). This database takes most of its line positions, intensities, and lower state energies from the HITRAN 2008 database, but the values for air-broadened half-widths and their temperature dependence parameters are adjusted from the HITRAN 2008 values to be consistent throughout the bands, and

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the air-induced pressure shifts (not given for a majority of transitions in HITRAN 2008) were added. For AER v3.1, the CO<sub>2</sub> line intensities and positions in the 597–2500 cm<sup>-1</sup> spectral range were then modified to be consistent with the Carbon Dioxide Spectral Database (Tashkun et al., 1998, 2003) as implemented for MIPAS retrievals (Flaud et al., 2003). These line parameters were used to calculate first-order line coupling parameters for all lines of CO<sub>2</sub> isotopologues 1 through 7 using the method and relaxations matrices of Lamouroux et al. (2010).

First-order line coupling parameters for the ν<sub>4</sub> and ν<sub>3</sub> bands of CH<sub>4</sub> were calculated using HITRAN 2008 line parameters with the method and relaxation matrices of Tran et al. (2006).

LBLRTM v12.1 uses v2.5.2 of the MT\_CKD continuum code, as described in detail by Mlawer et al. (2012). The self-broadened continuum for H<sub>2</sub>O between 2000 and 3200 cm<sup>-1</sup> is based on IASI, Atmospheric Infrared Sounder (AIRS) and AERI measurements between 2385 and 2600 cm<sup>-1</sup>, with the functional form for 2000–3200 cm<sup>-1</sup> shaped by the near-IR studies of Bicknell et al. (2006) and Fulghum and Tilleman (1991). When P- and R-branch line coupling for CO<sub>2</sub> was added to LBLRTM, the foreign broadened continuum of CO<sub>2</sub> was completely recalculated under the impact approximation. In MT\_CKD v2.5.2, empirical scaling factors, derived using IASI and AERI observations, were applied to these CO<sub>2</sub> continuum coefficients between 2000 and 3000 cm<sup>-1</sup>, with an additional correction for the temperature dependence of CO<sub>2</sub> absorption applied between 2386 and 2434 cm<sup>-1</sup> (Mlawer et al., 2012). The coefficients for the collision-induced fundamental bands of O<sub>2</sub> and N<sub>2</sub> in MT\_CKD 2.5.2 are the same as for MT\_CKD v1.2 (see Sect. 2.1).

### 3 IASI

The Infrared Atmospheric Sounding Interferometer (IASI) instrument, which is on the European meteorological polar orbiting meteorological satellite MetOp-A, was designed for (1) operational meteorological soundings (e.g. of temperature and H<sub>2</sub>O) with

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the goal of improving medium range weather forecasts and (2) studying atmospheric chemistry, with the goal of measuring and monitoring trace gases such as  $O_3$ ,  $CO$ , and  $CH_4$  on a global scale. IASI is an accurately calibrated Fourier Transform Spectrometer operating in the  $645\text{--}2760\text{ cm}^{-1}$  ( $15.5\text{--}3.6\text{ }\mu\text{m}$ ) spectral range with  $0.5\text{ cm}^{-1}$  (apodized) resolution (Challon et al., 2001). Instrument characteristics are as described in Shephard et al. (2009). The IASI instrument line shape is modeled as a truncated Gaussian with a  $1/e$  point at  $0.25\text{ cm}^{-1}$ . The calibration and validation of the IASI Level 1 radiances is described in Blumstein et al. (2007). The assessment of the in-flight performance has shown remarkably good radiometric performance (noise characteristics are similar to those measured on the ground) and excellent absolute calibration (better than  $0.1\text{ K}$ ). The spectral calibration was also shown to be remarkable (mean relative error on spectral calibration of  $3.1 \pm 3.3 \times 10^{-7}$ , Blumstein et al., 2007).

Here we use a set of 120 near-nadir spectra measured by IASI in April 2008, which are a subset of the profiles analyzed by Matricardi (2009). Only clear sky, ocean, nighttime cases were selected in order to minimize uncertainties associated with cloud, surface emissivity and non local thermodynamic equilibrium (non-LTE) effects. This left 130 profiles, 9 of which failed to converge for one or more of the retrievals discussed in Sect. 4 below, and 1 of which showed clearly biased residuals in the  $H_2O\text{ }v_2$  and  $CO_2\text{ }v_3$  band after convergence, leaving us with a set of 120 cases for analysis. These 120 profiles cover a wide range of surface temperatures and precipitable water vapor amounts. Table 1 shows the number of spectra with precipitable water vapor (PWV) between 0–2, 2–4, and 4–6 cm, along with the corresponding ranges of surface temperature ( $T_{\text{srf}}$ ) and latitude.

## 4 Radiance closure method

The validation of the molecular spectroscopy in a radiative transfer model like LBLRTM with observed spectral radiances from satellite and ground-based sensors, as in Shephard et al. (2009) and this study, requires a careful consideration of the dominant uncer-

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tainties in the input parameters of the model. For the cloudless, nighttime, over-ocean cases considered in this paper, the dominant sources of uncertainty in the model inputs are: (1) errors in the specification of the atmospheric state, i.e. the vertical profiles of temperature and trace gases for each case; and (2) errors in the spectroscopic line parameters and continua. In general, errors in the specification of the atmospheric state are unavoidable – either in situ profile measurements are unavailable, or the in situ measurements have significant uncertainties, or the in situ measurements did not sample the atmosphere at the exact same space and time as the satellite observation. Thus, in order to examine the spectroscopic parameters that are the focus of this work, we first minimize the errors in the specification of the atmospheric state by using each of the two versions of LBLRTM discussed in Sect. 2 to retrieve best-fit specifications of the atmospheric state for all cases in the data set.

We use an optimal estimation retrieval approach to minimize the difference between the observed IASI spectral radiances and corresponding LBLRTM calculations subject to the constraint that the estimated atmospheric state must be consistent with an a priori probability distribution for that state (Bowman et al., 2006; Clough et al., 1995; Rodgers, 2000). Specifically, we retrieve the following parameters: the surface or “skin” temperature  $T_{\text{srf}}$ ; the vertical profile of temperature  $T_{\text{atm}}$ ; and the vertical profiles of water vapor ( $\text{H}_2\text{O}$ ), ozone ( $\text{O}_3$ ), carbon monoxide ( $\text{CO}$ ), and methane ( $\text{CH}_4$ ).

Our retrieval procedure is shown schematically in Fig. 1. We start with an a priori specification of the atmospheric state for each IASI spectrum. This a priori specification was constructed as follows.  $T_{\text{srf}}$  is taken from the output of version 33R1 of the ECMWF model as discussed in Matricardi (2009). The a priori profile for  $T_{\text{atm}}$  below 10 hPa and above 0.1 hPa is taken from the same ECMWF model output. However, the work of Masiello et al. (2011) showed that ECMWF temperature profiles overestimate the true atmospheric temperature by up to 12 K between 10 hPa and 0.1 hPa, leading to significant errors in the radiances near the  $667\text{ cm}^{-1}$  Q-branch of  $\text{CO}_2$ . Thus, in this region we apply a correction to the ECMWF temperature profile. The correction peaks

at  $-12\text{ K}$  at  $0.3\text{ hPa}$ , and is linearly interpolated in the natural logarithm of pressure between this point and the  $0\text{ K}$  correction values at  $10\text{ hPa}$  and  $0.1\text{ hPa}$ .

The a priori profile for  $\text{H}_2\text{O}$  is also taken from the ECMWF model output as discussed in Matricardi (2009). For  $\text{O}_3$ , these ECMWF model profiles were scaled to match the total column of  $\text{O}_3$  as observed by the Ozone Monitoring Instrument (OMI). For  $\text{CO}_2$ ,  $\text{N}_2\text{O}$ ,  $\text{CH}_4$ , and  $\text{CO}$ , the a priori profiles are from the NASA Aura TES monthly climatology for April 2008, which is based on the MOZART global chemical transport model (Brasseur et al., 1998). A priori profiles for  $\text{HCOOH}$  and  $\text{C}_2\text{H}_4$  are from the climatologies developed for the NASA Upper Atmospheric Research Satellite (UARS, A. Goldman, personal communication, 2012). For  $\text{CCl}_4$ , CFC-11, CFC-12, and  $\text{CHClF}_2$  the (vertically well-mixed) a priori profiles are the UARS profiles scaled to match the April 2008 global average mixing ratios as reported by the NOAA/ESRL Global Monitoring Division (91.3 ppt, 245.4 ppt, 535.7 ppt, and 190.1 ppt, respectively; available at <ftp://ftp.cmdl.noaa.gov/hats/cfcs/>). The  $\text{CF}_4$  a priori profile is the UARS profile scaled to match the latest available observations (74 ppt in 1997; Khalil et al., 2003; Forster et al., 2007). A priori profiles for all other trace gases considered in this study (i.e.  $\text{SO}_2$ ,  $\text{NH}_3$ ,  $\text{HNO}_3$ ,  $\text{OCS}$ ,  $\text{HCN}$ , and  $\text{C}_2\text{H}_2$ ) were taken from the US Standard Atmosphere (NOAA, 1976). For all the retrievals discussed below, the a priori profile is also the initial guess profile for the retrievals. A priori spectral surface emissivity was estimated from the Wu and Smith (1997) model (zero wind speed, zero viewing angle) as described in van Delst and Wu (2000).

Starting with these a priori profiles, we performed the following retrievals for each IASI spectrum and version of LBLRTM. First, we performed two combined retrievals of  $T_{\text{srf}}$  and  $T_{\text{atm}}$ , as shown in Fig. 1. These retrievals use different spectral regions to provide information on  $T_{\text{atm}}$  – in the first, the  $\nu_2$  band of  $\text{CO}_2$  is used, while in the second the  $\nu_3$  band of  $\text{CO}_2$  is used. (Example averaging kernels for these and the other retrievals in this work are included in the Supplement, which also shows the pressure levels included in the state vector for each retrieval.) The profiles of  $T_{\text{atm}}$  retrieved using these two different spectral ranges are compared in Sect. 5.1 below in order to evaluate

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the consistency of the spectroscopy for the two CO<sub>2</sub> bands in both versions of LBLRTM. The set of spectral ranges or “microwindows” primarily sensitive to  $T_{\text{srf}}$  are the same in both of these retrievals – these ranges (and the spectral ranges used to retrieve all other parameters) are given in Table 2. The associated covariance matrix for  $T_{\text{srf}}$  and  $T_{\text{atm}}$  was constructed as follows. First,  $T_{\text{srf}}$  was assumed to have a standard deviation of 1 K. For  $T_{\text{atm}}$ , the diagonal of the covariance matrix was set by assuming that the a priori probability distribution of  $T_{\text{atm}}$  had a standard deviation of 2 K between the surface at 200 hPa, 4 K between 200 hPa and 50 hPa, 7 K between 50 and 10 hPa, and 10 K above 10 hPa. This gradually increasing uncertainty in the a priori of  $T_{\text{atm}}$  with altitude was chosen to avoid overconstraining the upper atmospheric temperatures, as the values in our corrected a priori profile above 10 hPa must be considered highly uncertain. We then assumed a Gaussian correlation length of 1 km between the retrieved levels in order to account for the correlations between the different retrieval levels (the off-diagonal elements of the covariance matrix).

Next, we used the a posteriori values for  $T_{\text{srf}}$  and  $T_{\text{atm}}$  from the CO<sub>2</sub>  $\nu_2$  band temperature retrieval as inputs into two retrievals of H<sub>2</sub>O (see Fig. 1; note that for all gases, the parameter retrieved was the natural log of the volume mixing ratio). In the first H<sub>2</sub>O retrieval, hereafter referred to as the “P- and R-branch retrieval,” all of the microwindows given for H<sub>2</sub>O in Table 2 are used. In the second retrieval, hereafter the “P-branch retrieval”, the range covering the R-branch of the  $\nu_2$  band of H<sub>2</sub>O (i.e. 1640–2020 cm<sup>-1</sup>) was excluded in order to investigate the impact on water vapor retrievals when only the P-branch is observed, as is the case for the Atmospheric Infrared Sounder (AIRS) and the Cross-track Infrared Sounder (CrIS) satellite instruments (see Sect. 5.2 below). The a priori covariance of H<sub>2</sub>O were assumed to have an uncertainty of 20 % at all levels with off-diagonal correlation lengths of 1 km.

The a posteriori profiles of  $T_{\text{srf}}$ ,  $T_{\text{atm}}$ , and H<sub>2</sub>O from the P- and R-branch retrieval of H<sub>2</sub>O were then used as input to a more highly constrained retrieval of  $T_{\text{srf}}$  and  $T_{\text{atm}}$ . This additional retrieval step is required to minimize the impacts of errors in the atmospheric state on the final spectral residuals, as the temperatures retrieved using the a priori H<sub>2</sub>O

profile can be somewhat different than the temperatures retrieved using the a posteriori H<sub>2</sub>O profile. In this step, both  $T_{\text{atm}}$  and  $T_{\text{srf}}$  were assumed to have standard deviations of 1 K at all heights, and the off-diagonal elements for  $T_{\text{atm}}$  a priori covariance matrix were again calculated using a 1 km correlation length.

5 These improved a posteriori profiles of  $T_{\text{srf}}$  and  $T_{\text{atm}}$ , along with the H<sub>2</sub>O profiles from the P- and R-branch retrieval, were then used as input to sequential retrievals of O<sub>3</sub>, CO, and CH<sub>4</sub> to produce our final estimate of the atmospheric state and thus our final set of spectral residuals (see Fig. 1; note that all residual plots in this paper are from this final set of residuals after all parameters have been retrieved). The a priori  
10 covariance matrices of O<sub>3</sub>, CO, and CH<sub>4</sub> were generated assuming uncertainties of 20 %, 20 %, and 5 % at all levels and off-diagonal correlation lengths of 1 km, 2 km, and 1 km, respectively.

## 5 Results and discussion

Figure 2 shows the mean residuals for the 120 scans across the IASI spectral range using LBLRTM v12.1 and LBLRTM v9.4+, along with an example observed spectrum for a profile with 1.5 cm PWV. The figures also display the mean and root-mean-square (RMS) of the residuals for each model across the IASI spectral range. Note that the RMS is the RMS of the mean residuals plotted in the figure, that is:

$$\text{RMS} = \sqrt{\frac{\sum_{j=1, N_{\text{scans}}} (\text{Obs}_{i,j} - \text{Model}_{i,j})^2}{N_{\text{scans}}}} \quad (1)$$

20 where  $\text{Obs}_{i,j}$  is the IASI-observed radiance (or brightness temperature) in spectral channel  $i$  for scan  $j$ ,  $\text{Model}_{i,j}$  is the corresponding LBLRTM-simulated radiance or brightness temperature, and  $N_{\text{scans}}$  and  $N_{\text{channels}}$  are the number of scans and spectral channels included in the average, respectively.

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Figure 2 shows that the spectroscopy in LBLRTM v12.1 is generally improved from that in LBLRTM v9.4+; however, significant systematic spectral residuals remain in LBLRTM v12.1. In the following sections, we discuss the changes in the spectral residuals in several spectral regions within the IASI spectral range, along with the associated impact on the retrievals on atmospheric profiles of temperature, H<sub>2</sub>O, O<sub>3</sub>, CO, and CH<sub>4</sub>.

## 5.1 CO<sub>2</sub> $\nu_2$ (640–800 cm<sup>-1</sup>) and $\nu_3$ (2200–2550 cm<sup>-1</sup>) bands

Figure 3 shows the mean of the final brightness temperature residuals in the CO<sub>2</sub>  $\nu_2$  band for all 120 IASI scans considered here for both LBLRTM v12.1 (Fig. 3b) and v9.4+ (Fig. 3c). The CO<sub>2</sub>  $\nu_2$  band atmospheric temperature retrieval microwindow is shown in red in both panels, along with the mean and RMS of the residuals within this microwindow.

The addition of P- and R-branch line coupling to LBLRTM (and the associated recalculation of the CO<sub>2</sub> continuum in MT\_CKD) clearly improved the spectroscopy on either side of the CO<sub>2</sub> Q-branch at 720 cm<sup>-1</sup>. The RMS of the residuals in the  $\nu_2$  microwindow shows substantial improvement in LBLRTM v12.1 (0.12 K) versus v9.4+ (0.21 K). The major remaining residual features in LBLRTM v12.1 are negative residuals of  $\sim 0.5$  K in the 667 and 720 cm<sup>-1</sup> Q-branches and positive offset of  $\sim 0.2$  K between 755 and 770 cm<sup>-1</sup>.

Figures 4 and 5 show the average residuals binned by PWV for LBLRTM v12.1 and LBLRTM v9.4+, respectively. The residual feature at the 667 cm<sup>-1</sup> Q-branch appears to be independent of PWV, while the residuals in the 720 cm<sup>-1</sup> Q-branch and between 755 and 770 cm<sup>-1</sup> both appear to increase with increasing amounts of water vapor. The RMS of the residuals in the  $\nu_2$  microwindow also appears to increase with PWV in both models. However, since PWV is well correlated with atmospheric temperature and latitude (see Table 1), these apparent dependencies on PWV may actually be due to dependencies on atmospheric temperature and tropopause height, with increasing residuals in the warmer, wetter, tropical regions.

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The residuals in the  $667\text{ cm}^{-1}$  Q-branch of  $\text{CO}_2$  are greatly improved from the study of Shephard et al. (2009), where the residuals between LBLRTM calculations and IASI measurements in this region were  $\sim 1.7\text{ K}$ . This is because, following the study of Masiello et al. (2009), we now include the  $667\text{ cm}^{-1}$  Q-branch in our temperature retrievals and allow the temperature near the stratopause ( $\sim 1\text{ hPa}$ ) to adjust. Our mean radiance residual of  $-4.2 \times 10^{-8}\text{ W cm}^2\text{ cm}^{-1}\text{ sr}^{-1}$  ( $-0.45\text{ K}$  in brightness temperature) is very similar to the remaining residual shown in Fig. 12 of Masiello et al. (2011). This improvement in the residuals is consistent with the hypothesis of Shephard et al. (2009) that the residual they observed at the  $667\text{ cm}^{-1}$  Q-branch was most likely due to errors in the temperature profile in the upper stratosphere and mesosphere. However, the remaining residual is still too large to be accounted for by randomly-distributed instrument noise: for the 120 scans considered here, the error in the mean residual due to noise should be only  $6.5 \times 10^{-9}\text{ W cm}^2\text{ cm}^{-1}\text{ sr}^{-1}$ . Therefore, there is still a systematic error in the spectroscopy in this region, possibly due to errors in the line parameters (e.g. line positions) or due to non-Voigt line shapes (e.g. speed dependence and/or Dicke narrowing).

Figure 6 shows the differences between the temperature profiles retrieved using the updated  $\text{CO}_2$  spectroscopy in LBLRTM v12.1 versus those retrieved using LBLRTM v9.4+. (Note that in Fig. 6 and all similar plots in this paper, the error bars represent the variability, expressed as standard deviation, among the 120 cases analyzed, and do not represent the estimated error for a single retrieved profile.) Here both models used the  $\nu_2$  band of  $\text{CO}_2$  to retrieve temperature. The differences between the two temperature profiles show an oscillatory structure versus altitude, consistent with the reduced oscillation in the residuals when LBLRTM v12.1 is used (see Fig. 3). Examination of the Jacobian from the temperature retrievals (not shown) suggests that the retrieved temperatures at altitudes below 500 hPa are most sensitive to the P- and R-branches on the left and right of the  $720\text{ cm}^{-1}$  Q-branch (approximately  $700\text{--}718\text{ cm}^{-1}$  and  $722\text{--}750\text{ cm}^{-1}$ ), while the retrieved temperatures at altitudes above 500 hPa are most sensitive to the spectroscopy of the  $720\text{ cm}^{-1}$  Q-branch itself ( $\sim 718\text{--}722\text{ cm}^{-1}$ ).

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With the updated spectroscopy the retrieved temperatures are on average  $0.6 \pm 0.8$  K higher at 562 hPa, consistent with the  $\sim -0.5$  K shift in the mean residuals in the wings surrounding the  $720 \text{ cm}^{-1}$  Q-branch as a result of the addition of P- and R-branch line coupling to LBLRTM. The retrieved temperatures are on average  $0.5 \pm 0.4$  K lower in the upper troposphere near 300 hPa,  $0.8 \pm 0.7$  K higher in the UTLS region between 100 and 200 hPa,  $0.4 \pm 0.4$  K lower between 40–100 hPa, and  $0.3 \pm 0.5$  K lower between 10–30 hPa.

Figure 7 shows the mean of the final residuals for both model versions in the  $\text{CO}_2 \nu_3$  band for the 120 IASI spectra. The retrieval microwindow for the  $\nu_3$  atmospheric temperature retrieval is shown in green. However, it is important to recall that the residuals in the  $\nu_3$  microwindow are not included in the main line of our retrieval procedure. Thus, these residuals, as well as those in Figs. 8, 9, show the average residuals in this region after temperature has been retrieved using only the  $\text{CO}_2 \nu_2$  band, and thus these  $\text{CO}_2 \nu_3$  residuals can be used to assess the consistency of the spectroscopy between the bands.

Figure 7 shows that the spectroscopy in the  $\text{CO}_2 \nu_3$  band has been greatly improved in LBLRTM v12.1, especially in the region past the bandhead ( $2385\text{--}2500 \text{ cm}^{-1}$ ). However, a small systematic residual near the bandhead remains in LBLRTM v12.1. In addition, the large negative residuals between  $2200\text{--}2270 \text{ cm}^{-1}$  in both Fig. 7b, c suggest that the  $\text{N}_2\text{O}$  optical depth in this region is still largely underestimated in LBLRTM v12.1. While our a priori  $\text{N}_2\text{O}$  profile should be fairly accurate in the troposphere (where  $\text{N}_2\text{O}$  is a well-mixed gas), this spectral region is sensitive to  $\text{N}_2\text{O}$  in the stratosphere, which can have a significant day-to-day variability (Randel et al., 1994). Thus, the observed underestimate of optical depth could be due to errors in the  $\text{N}_2\text{O}$  spectroscopy or to errors in the stratospheric portion of our a priori  $\text{N}_2\text{O}$  profile.

Figures 8 and 9 show the mean residuals in the  $\text{CO}_2 \nu_3$  band for v12.1 and v9.4+, respectively, binned by PWV. The residuals in LBLRTM v12.1 have little dependence on water vapor, in contrast to the results for LBLRTM v9.4+, showing that the increase in the  $\text{H}_2\text{O}$  self continuum between  $2385\text{--}2500 \text{ cm}^{-1}$  in MT\_CKD v2.5 reproduces well the

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observed absorption in this region (Mlawer et al., 2012). The self continuum absorption in MT\_CKD v2.5 in this spectral region is somewhat less than corresponding values determined in recent laboratory (Baranov et al., 2011; Ptashnik et al., 2011) and field studies (Strow et al., 2006). The results shown in Fig. 8c would permit a slight increase in the water vapor continuum optical depth past the  $\nu_3$  bandhead, but this modification would negatively affect the residuals for cases with 2–4 cm of PWV (Fig. 8b). Other recent work (e.g. Baranov et al., 2011, 2012) has suggested that the collision-induced  $N_2$  continuum in this region may be underestimated at high PWV, as collisions between  $H_2O$  and  $N_2$  are more effective than collisions between  $N_2$  and air at inducing this absorption. We have not investigated the impact of this possible modification to the  $N_2$  continuum in MT\_CKD on the cases in the three PWV categories shown in Fig. 8.

A second test of the consistency of the spectroscopy in the  $CO_2$   $\nu_2$  and  $\nu_3$  bands is to evaluate the consistency of the atmospheric temperature profiles retrieved using each band. Generally, the  $\nu_3$  retrievals have a lower number of degrees of freedom for signal (DOFS) than the  $\nu_2$  retrievals (e.g. 7.1 vs. 10.3 for a moderate PWV IASI spectrum). Because of this, the retrievals should not be compared directly, as some of the differences between the retrievals will be due to the lower resolution of the  $\nu_3$  retrievals. Instead, we smooth the  $\nu_2$  retrievals by applying the averaging kernel and retrieved profile of the  $\nu_3$  retrievals, following the procedure of Rodgers and Connor (2003). The smoothed  $\nu_2$  temperature profile ( $\hat{\mathbf{x}}_{\nu_2, \text{smooth}}$ ) is calculated as:

$$\hat{\mathbf{x}}_{\nu_2, \text{smooth}} = \hat{\mathbf{x}}_{\nu_3} + \mathbf{A}_{\nu_3} \left( \hat{\mathbf{x}}_{\nu_2} - \hat{\mathbf{x}}_{\nu_3} \right) \quad (2)$$

where  $\hat{\mathbf{x}}_{\nu_2}$  is the retrieved  $\nu_2$  temperature profile,  $\hat{\mathbf{x}}_{\nu_3}$  is the retrieved  $\nu_3$  temperature profile, and  $\mathbf{A}_{\nu_3}$  is the averaging kernel of the  $\nu_3$  temperature retrieval. The differences between these smoothed  $\nu_2$  temperature retrievals and the  $\nu_3$  retrievals can thus illustrate inconsistencies in the modeling of these two spectral regions.

Figure 10a shows the mean and standard deviation of the differences between the  $\nu_3$  and  $\nu_2$  temperature retrievals before smoothing for both model versions, while Fig. 10b shows the same information after applying the smoothing as in Eq. (2). Figure 11 shows

the smoothed differences binned by PWV. All of these plots show that the mean differences between the temperature profiles are substantially reduced in LBLRTM v12.1, especially in the stratosphere. In addition, the differences show less dependence on PWV in LBLRTM v12.1, again demonstrating the improved CO<sub>2</sub> spectroscopy.

## 5.2 H<sub>2</sub>O ν<sub>2</sub> band (1350–2050 cm<sup>-1</sup>)

Figure 12 shows the mean of the final brightness temperature residuals in the H<sub>2</sub>O ν<sub>2</sub> band for both LBLRTM v12.1 (Fig. 12b) and v9.4+ (Fig. 12c). The P-branch and R-branch H<sub>2</sub>O retrieval windows (1375.0–1560.0 cm<sup>-1</sup> and 1640.0–2020.0 cm<sup>-1</sup>, respectively) are shown in red in both panels, which also show the mean and root mean square of the residuals in each window. The updated H<sub>2</sub>O spectroscopy in LBLRTM v12.1 substantially reduces the RMS of the residuals in the H<sub>2</sub>O P-branch (to 0.27 K from 0.34 K) and R-branch (to 0.31 K from 0.34 K). While the R-branch improvement is more modest, the updated spectroscopy has reduced several features that were present in the LBLRTM v9.4+ R-branch residuals (e.g. the positive spikes near 1920 cm<sup>-1</sup>). Figures 13 and 14 show the mean residuals in the H<sub>2</sub>O ν<sub>2</sub> band for v12.1 and v9.4+, respectively, binned by PWV. In both models, the P-branch RMS is largest for the cases with 2–4 cm PWV, while the R-branch residuals appear to steadily increase with PWV.

In discussing the remaining residuals in the P-branch, one issue to consider is the error induced by mis-specification of the isotopic ratios. In both versions of LBLRTM shown here, the abundances of minor isotopologues of all molecules are set to fixed ratios relative to the major isotopologue for that molecule. These reference ratios are consistent with the reference ratios used in the HITRAN database. However, atmospheric profiles of these ratios may vary substantially from these default reference values. In particular, atmospheric water vapor is always depleted in minor isotopologues relative to the reference standard, and the isotopic ratios show substantial variations in space (both vertically and horizontally) and time due to condensation/evaporation processes in the atmosphere (e.g. Worden et al., 2006, 2007, 2012; Nassar et al., 2007;

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Payne et al., 2007). This effect is most strongly pronounced for HDO. As HDO has substantial spectral features in the P-branch of the  $\text{H}_2\text{O } \nu_2$  band, the fixed HDO/ $\text{H}_2\text{O}$  ratio in LBLRTM v12.1 could be responsible for some of the remaining residual features seen in Fig. 12b.

For example, Fig. 15 shows the P-branch residuals for a high water vapor spectrum (5.4 cm PWV) immediately after the P- and R-branch retrieval of  $\text{H}_2\text{O}$ , plotted against the optical depth of HDO (calculated assuming the HITRAN reference ratio). The original LBLRTM v12.1 results are plotted as black circles. We can see that there is a clear upward trend in the residuals as HDO optical depth increases, suggesting that LBLRTM v12.1 is on average overestimating the optical depth of HDO for this set of cases. This is consistent with what we know about the real atmosphere, which is nearly always depleted in HDO relative to the HITRAN reference ratio. The red triangles show the same model results when the HDO/ $\text{H}_2\text{O}$  ratio is set to the uniform climatology used in the Aura TES retrievals and the  $\text{H}_2\text{O}$  retrieval is repeated. In order to model this vertically-varying HDO/ $\text{H}_2\text{O}$  ratio in LBLRTM, we temporarily added HDO as a separate molecule and removed the HDO lines from the  $\text{H}_2\text{O}$  line parameters. This removed 29 % of the remaining bias in the P-branch in this case, suggesting that errors in the HDO profile are responsible for a significant fraction of the remaining bias in the P-branch of the  $\text{H}_2\text{O } \nu_2$  band. Note that the same exercise could be performed for  $\text{H}_2^{18}\text{O}$  and  $\text{H}_2^{17}\text{O}$ . The deviation of isotopic ratios from the reference standard is  $\sim 5$  times smaller for these isotopologues than for HDO, but their atmospheric abundances are considerably larger.

We also examined how the differences in  $\text{H}_2\text{O}$  line parameters between HITRAN 2008 and the AER v3.1 parameters used in LBLRTM v12.1 impacted the spectral residuals that remain after the corresponding P- and R-branch  $\text{H}_2\text{O}$  retrievals. As discussed in Sect. 2.2, HITRAN 2008 did not adopt the calculated line positions and intensities from Coudert et al. (2008), which are the primary source of these line parameters in AER v3.1 in the  $1750\text{--}2500\text{ cm}^{-1}$  range. Figure 16a shows the residuals for a moderate PWV case after the  $\text{H}_2\text{O}$  retrieval when the HITRAN 2008  $\text{H}_2\text{O}$  line parameters

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are used in LBLRTM v12.1. Figure 16b shows the same for a standard run of LBLRTM v12.1 using the AER v3.1 H<sub>2</sub>O line parameters, and Fig. 16c shows the difference between the two model runs. The two sets of line parameters agree fairly well for wavenumbers below 1750 cm<sup>-1</sup>, but above this wavenumber the HITRAN 2008 parameters give generally larger brightness temperatures than AER v3.1. This region of substantial disagreement between HITRAN 2008 and AER v3.1 (1750 cm<sup>-1</sup>–2020 cm<sup>-1</sup>) is highlighted in red in Fig. 16, and the mean and RMS of the residuals in this range are given as well. We can see in Fig. 16 that the HITRAN 2008 residuals in Fig. 16a leave a substantially larger bias in this region than those from LBLRTM v12.1 using AER v3.1 line parameters (–0.28 K versus –0.12 K) while the RMS is unchanged, suggesting that the HITRAN 2008 H<sub>2</sub>O parameters have a discontinuity between the parts of the H<sub>2</sub>O ν<sub>2</sub> band above and below 1750 cm<sup>-1</sup> due to the exclusion of the calculated Coudert et al. (2008) parameters. We extended this analysis to two other example cases, corresponding to low (0.3 cm) and high (5.4 cm) values of PWV. Table 3 shows the mean and RMS of the residuals for these spectra using both HITRAN 2008 and AER v3.1 H<sub>2</sub>O line parameters. We examine three spectral ranges: the P-branch microwindow (1375–1560 cm<sup>-1</sup>), the R-branch microwindow (1640–2020 cm<sup>-1</sup>) and the “region of disagreement” in the R-branch (1750–2020 cm<sup>-1</sup>). We can see that while the mean and RMS of the residuals are virtually identical in the P-branch for all three spectra, in the R-branch the HITRAN 2008 line parameters give a substantially larger mean residual for the 1.5 cm PWV spectrum and the 5.4 cm PWV spectrum, as well as giving larger RMS values for all three spectra examined here. In the “region of disagreement” the absolute value of the mean residual for each of the three spectra is below 0.12 K when AER v3.1 is used, while the mean residual is about –0.28 K for the 1.5 cm PWV and –0.31 K for the 5.4 cm PWV spectra when HITRAN 2008 is used. This suggests that the exclusion of the calculated Coudert et al. (2008) parameters in HITRAN 2008 leads to an unphysical discontinuity in the spectra near 1750 cm<sup>-1</sup>, with negative impacts on the spectral residuals in the R-branch of the H<sub>2</sub>O ν<sub>2</sub> band. In contrast, using the full set of the Coudert et al. (2008) line positions and intensities

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between 10–2500  $\text{cm}^{-1}$ , as in LBLRTM v12.1 and AER v3.1, results in more consistent spectroscopy within the R-branch.

It is also possible that the improvement in the  $\text{H}_2\text{O } \nu_2$  band residuals in LBLRTM v12.1 is not due to improved water vapor spectroscopy, but instead is due to the improvements in the  $\text{CO}_2$  spectroscopy, which then gives a better estimate of the temperature profile and thus results in better modeling of the  $\text{H}_2\text{O } \nu_2$  band. To examine this possibility, we ran two additional  $\text{H}_2\text{O}$  retrievals for our three representative cases: one that used the temperature profile retrieved by LBLRTM v9.4+ but used LBLRTM v12.1 (with its updated  $\text{H}_2\text{O}$  spectroscopy) as the forward model, and another that used the temperature profile retrieved by LBLRTM v12.1 but used LBLRTM v9.4+ (with its older  $\text{H}_2\text{O}$  spectroscopy) as the forward model. We performed these additional retrievals using both the P- and R-branches of the  $\text{H}_2\text{O } \nu_2$  band. This gave us three cases with four different water vapor retrievals each, corresponding to the 4 possible combinations of retrieved temperature profile and  $\text{H}_2\text{O}$  spectroscopy.

Table 4 shows the mean and RMS of the residuals in the  $\text{H}_2\text{O}$  P- and R-branch retrieval microwindows for the three cases and four combinations discussed above. Note that this comparison is looking at the residuals immediately after the  $\text{H}_2\text{O}$  retrieval, rather than at the end of the entire retrieval procedure as in Figs. 12–14. This means that the results in Table 4 do not include the impact of the second, constrained temperature retrieval that was performed after the  $\text{H}_2\text{O}$  retrieval. Table 4 shows that both the improved temperature profiles and the updated  $\text{H}_2\text{O}$  spectroscopy in LBLRTM v12.1 contribute to the improvements of the P-branch residuals for each case examined here. The results are more ambiguous in the R-branch. First, the R-branch RMS is the same or slightly higher for LBLRTM v12.1 than it was in LBLRTM v9.4+. This is consistent with the residuals averaged over all 120 scans immediately after the  $\text{H}_2\text{O}$  retrieval (not shown), in which the mean R-branch RMS increases slightly from 0.28 K in v9.4+ to 0.29 K in v12.1. This suggests that the improvement in the RMS for the R-branch that we saw in Figs. 12–14 are mainly due to our second, constrained temperature retrieval, rather than improvement in the  $\text{H}_2\text{O}$  spectroscopy in this region. This

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is consistent with the results in Table 4, where the RMS in the R-branch is higher when the new H<sub>2</sub>O spectroscopy is used, regardless of which temperature profile is used. Thus in the P-branch of the H<sub>2</sub>O ν<sub>2</sub> band, the decrease in the RMS of the mean residuals is likely due to a combination of the improvements in H<sub>2</sub>O spectroscopy and the improvements in the CO<sub>2</sub> spectroscopy, while in the R-branch the improvements are mainly due to improvements in the CO<sub>2</sub> spectroscopy and our second, constrained temperature retrieval.

Both the CO<sub>2</sub> and the H<sub>2</sub>O spectroscopic improvements in LBLRTM v12.1 can also have a large impact on the retrieved profiles of H<sub>2</sub>O. The H<sub>2</sub>O mixing ratios retrieved by LBLRTM v12.1 and LBLRTM v9.4+ can differ by a factor of 2 or more. Figure 17 shows the mean and standard deviation of the ratio of the LBLRTM v12.1 retrieved H<sub>2</sub>O profiles (P- and R-branch) to the v9.4+ retrieved profiles for the 120 IASI spectra. We can see that the new spectroscopy generally reduces the retrieved H<sub>2</sub>O mixing ratio between 100 and 200 hPa by 14 %, with a variability (expressed as a standard deviation) among the cases of ±8 %. This is consistent with the results of Shephard et al. (2009), who first showed that the addition of the Coudert et al. (2008) line positions and intensities to HITRAN 2004 resulted in a 10 % reduction in the upper tropospheric H<sub>2</sub>O mixing ratio. At 562 hPa the new spectroscopy increases the retrieved H<sub>2</sub>O mixing ratio by 42 % on average, but the variability is very large (±58 %). Near the surface, this variability is quite large (over a factor of 2), so even though the mean difference is 31 % there can be dramatic differences for any given case.

In order to examine whether CO<sub>2</sub> or H<sub>2</sub>O spectroscopic changes between LBLRTM v12.1 and v9.4+ are primarily responsible for the changes in the retrieved H<sub>2</sub>O seen in Fig. 17, we analyzed the retrieved H<sub>2</sub>O profiles for our three representative cases and the 4 possible combinations of retrieved temperature profile and H<sub>2</sub>O spectroscopy (see Table 4 and the discussion above). Figure 18 shows the (a) temperature differences and (b) retrieved H<sub>2</sub>O differences for our example case with 1.5 PWV. We plot three H<sub>2</sub>O profiles, which are the profile retrieved using (i) LBLRTM v12.1 to retrieve both temperature and H<sub>2</sub>O (black circles), (ii) LBLRTM v9.4+ to retrieve temperature

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and LBLRTM v12.1 to retrieve H<sub>2</sub>O (red triangles), and (iii) LBLRTM v12.1 to retrieve temperature and LBLRTM v9.4+ to retrieve H<sub>2</sub>O (green diamonds). All three profiles are normalized by the H<sub>2</sub>O profile retrieved when LBLRTM v9.4+ is used to retrieve both temperature and H<sub>2</sub>O.

5 For the moderate water vapor case (Fig. 18), the temperature change is relatively large (> 1.5 K) near the surface, and this temperature change seems to be responsible for most of the lower troposphere water vapor changes. However, the reduction in H<sub>2</sub>O between 100 hPa and 300 hPa is primarily caused by the new H<sub>2</sub>O spectroscopy. Based on Shephard et al. (2009), this upper atmospheric change is more likely due to the inclusion of the Coudert et al. (2008) positions and intensities in AER v3.1, rather than  
10 the H<sub>2</sub>O width updates that were included in HITRAN 2008. In the low water vapor case (Fig. 19), the temperature differences are low, so the changes to H<sub>2</sub>O spectroscopy are primarily responsible for the changes in the retrieved H<sub>2</sub>O profile in this case, with the width changes likely more important in the lower troposphere. However, the high water  
15 vapor case (not shown) shows similar results to the moderate case, where temperature changes dominate in the lower troposphere. Thus, the observed mean changes in the retrieved H<sub>2</sub>O profile in Fig. 17 are likely due to a combination of the H<sub>2</sub>O spectroscopy changes and the temperature profile changes (themselves caused by CO<sub>2</sub> spectroscopy changes), with the relative importance of each depending on the size of  
20 the temperature profile changes.

Finally, we find that the R-branch can provide additional information on near-surface H<sub>2</sub>O that is not available from the P-branch due to the interference of CH<sub>4</sub> and N<sub>2</sub>O. This is shown in Fig. 20, which shows the DOFS versus PWV for the P- and R-branch retrieval (black circles) and the P-branch only retrieval (red triangles) of H<sub>2</sub>O using  
25 LBLRTM v12.1. We can see that the P- and R-branch generally has higher values for DOFS, and that the difference is largest for the moist cases. Examination of the averaging kernels for these moist cases (not shown) demonstrates that the additional information is primarily in the near-surface layers. This suggests that an instrument that uses both the P- and R-branches like IASI will have more sensitivity to near-surface

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water vapor in moist atmospheres than an instrument like AIRS or CrIS, which only measures radiances within the P-branch. Thus including the R-branch in the water vapor retrieval can substantially alter the retrieved H<sub>2</sub>O mixing ratio near the surface. Figure 21 shows the mean and standard deviation of the ratio of the LBLRTM v12.1 retrieved H<sub>2</sub>O profiles that only included the P-branch to those that included both the P- and R-branches as discussed in Sect. 4. Both the mean and variability of the differences is small above 500 hPa, but the P-branch only retrievals show on average 33 % more water vapor than the P- and R-branch retrievals near the surface, with a variability of  $\pm 51$  %. These differences are not necessarily due to inconsistencies in the spectroscopy between the branches, but rather reflect that the region of the P-branch most sensitive to near-surface water vapor were removed from the retrieval to avoid interferences from the CH<sub>4</sub>  $\nu_4$  band (centered near 1306 cm<sup>-1</sup>) and the N<sub>2</sub>O  $\nu_1$  band (centered near 1285 cm<sup>-1</sup>).

### 5.3 O<sub>3</sub> $\nu_3$ band (950–1150 cm<sup>-1</sup>)

Figure 22 shows the mean of the final brightness temperature residuals in the O<sub>3</sub>  $\nu_3$  band for the 120 IASI spectra for both LBLRTM v12.1 (Fig. 22b) and v9.4+ (Fig. 22c). The O<sub>3</sub> retrieval microwindow (see Table 1) is highlighted in red. The residuals show little change between LBLRTM v9.4+ and v12.1: both show positive residuals of about 0.3K in the R-branch of the band that have little dependence on PWV. As the O<sub>3</sub> spectroscopy in this region was not substantially changed between LBLRTM v9.4+ and v12.1, these small differences are likely due to the improved temperature retrieval in LBLRTM v12.1 (see Sect. 5.1). The retrieved ozone profiles (not shown) also show little change between the two model versions: the mean differences between the profiles are generally less than 3 % below 10 hPa.

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## 5.4 CO fundamental band (2050–2250 cm<sup>-1</sup>)

Figure 23 shows the mean of the final brightness temperature residuals in the fundamental vibrational band of CO. The CO retrieval microwindows (see Table 1) are illustrated with red dots, and the mean and RMS residuals for the CO microwindows are printed in red. We can see that the residuals between 2060–2170 cm<sup>-1</sup> have increased in LBLRTM v12.1, both within and outside the CO microwindows. As CO spectroscopy was unchanged between LBLRTM v9.4+ and v12.1, the differences between the two model versions in this spectral region is primarily due to the changes in water vapor spectroscopy in this region. Figure 24a shows the optical depth from the major absorbers in this region for a profile with 1.5 cm PWV, while Fig. 24b shows the changes in optical depth between the two versions of LBLRTM. We can see that the optical depth of the H<sub>2</sub>O lines has substantially increased in this region, and that the H<sub>2</sub>O self continuum was increased in this region in MT\_CKD v2.5. However, since the CO microwindows were selected to avoid strong water lines, the increased positive residuals within the CO microwindows are likely due to the increase of the H<sub>2</sub>O self continuum in MT\_CKD v2.5. This change to the continuum was motivated by improving agreement with measurements near the CO<sub>2</sub> ν<sub>3</sub> bandhead (see Sect. 5.1), but it clearly has the unintended consequence of degrading the match with observations between 2060–2170 cm<sup>-1</sup>. It also affects the retrieved profiles of CO (not shown): the differences can be 40 ppb or larger for cases with PWV between 4–6 cm. Thus, correcting the MT\_CKD 2.5 H<sub>2</sub>O self-continuum in this spectral region is critical to the accurate retrieval of CO concentrations. It is important to note that the recently developed CAVIAR water vapor self continuum (Ptashnik et al., 2011) is even greater than MT\_CKD\_2.5 in this region.

## 5.5 CH<sub>4</sub> ν<sub>4</sub> band (1250–1350 cm<sup>-1</sup>)

Figure 25 shows the mean of the final average brightness temperature residuals in the CH<sub>4</sub> ν<sub>4</sub> band for both LBLRTM v12.1 (Fig. 25b) and v9.4+ (Fig. 25c). The HITRAN 2008 updates to the CH<sub>4</sub> line parameters, along with the line mixing calculations of Tran

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et al. (2006), have substantially changed the shape of the residuals, leaving three clear residual peaks at about 1295, 1297.5 and 1300  $\text{cm}^{-1}$  in LBLRTM v12.1 that have only a slight dependence on PWV. However, it is not clear if the residuals have substantially improved: although the mean residual is now closer to 0 in LBLRTM v12.1, the RMS of the residuals is larger, making the effect of the spectroscopy change ambiguous. Figure 26, which shows the optical depth for  $\text{CH}_4$ ,  $\text{H}_2\text{O}$  (which includes HDO), and  $\text{N}_2\text{O}$  in this spectral region, demonstrates that the three residual peaks in Fig. 25b are associated with  $\text{CH}_4$  lines, and thus likely reflect errors in  $\text{CH}_4$  spectroscopy. We also examined the impact of including  $\text{CH}_4$  line coupling on the spectral residuals and retrieved  $\text{CH}_4$  profiles for our three representative spectra (see Sect. 5.2). The addition of line coupling to HITRAN 2008 changed the mean of the spectral residuals in the  $\text{CH}_4$  microwindow by less than 0.004 K and the RMS by less than 0.003 K, suggesting that the addition of  $\text{CH}_4$  line coupling had little impact on the differences between LBLRTM v12.1 and LBLRTM v9.4+ in the  $\text{CH}_4 \nu_4$  band.

The spectroscopic issues in the  $\text{CH}_4 \nu_4$  band will impact the retrieved  $\text{CH}_4$  profiles. Retrieval of  $\text{CH}_4$  from satellite observations is recognized as a difficult problem due to the presence of interfering species (e.g.  $\text{H}_2\text{O}$ , HDO,  $\text{N}_2\text{O}$ ; Worden et al., 2012), errors in the  $\text{CH}_4$  spectroscopy (such as those discussed above), and the fact that, as  $\text{CH}_4$  is reasonably well-mixed in the troposphere, we are interested in small changes ( $\sim 50$  ppbv) on a relatively large background mixing ratio ( $\sim 1800$  ppbv). Most operational satellite retrievals of  $\text{CH}_4$  using the  $\nu_4$  band show positive biases relative to in situ observations (e.g. Razavi et al., 2009; Wecht et al., 2012), which are corrected for by various methods, such as correcting the  $\text{CH}_4$  retrieved profile by assuming that deviations of the retrieved  $\text{N}_2\text{O}$  profile from the a priori are a result of systematic errors that also impact the retrieved  $\text{CH}_4$  (Razavi et al., 2009; Worden et al., 2012). As the goal of our study is to validate the spectroscopy in LBLRTM, rather than to design an accurate retrieval for  $\text{CH}_4$ , our retrieval approach was fairly simplistic (see Sect. 4). We performed a single-species retrieval and used only a small spectral region in an attempt to avoid interference from  $\text{H}_2\text{O}$  and  $\text{N}_2\text{O}$ . We recognize that we retrieved unphysically

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large mixing ratios of CH<sub>4</sub> in the free troposphere in both model versions. However, the average retrieved CH<sub>4</sub> mixing ratio below 100 hPa was reduced by 44 ppbv when the updated CH<sub>4</sub> spectroscopy in LBLRTM v12.1 was used, suggesting that the updated spectroscopy might improve operational retrievals of CH<sub>4</sub> from satellites by reducing the current positive biases in these retrievals.

## 6 Conclusions

We have performed an extensive validation of the thermal infrared spectroscopy (between 645–2760 cm<sup>-1</sup>) in LBLRTM v12.1 using a global dataset of 120 clear-sky, nighttime spectra measured over the ocean with the IASI instrument, and have compared the performance of LBLRTM v12.1 with a previous version of LBLRTM (v9.4+) to determine if the spectroscopy in various spectral regions has improved over time.

We find that the CO<sub>2</sub> spectroscopy in the  $\nu_2$  and  $\nu_3$  bands is significantly improved in LBLRTM v12.1 relative to v9.4+. The spectroscopy of the two bands is remarkably consistent in LBLRTM v12.1, as determined both by spectral residuals and by comparing the atmospheric temperature profiles retrieved with each band. The improvement in the spectroscopy in these bands is mainly due to (1) the addition of P- and R-branch line coupling for CO<sub>2</sub> based on Lamouroux et al. (2010), (2) the addition of CDS line positions and intensities (Tashkun et al., 2003; Flaud et al., 2003), and (3) improvements in the CO<sub>2</sub> foreign and H<sub>2</sub>O self continua near the  $\nu_3$  bandhead. Including the 667 cm<sup>-1</sup> Q-branch of CO<sub>2</sub> in the retrieval of atmospheric temperature substantially improves the spectral residuals in this region relative to the results of Shephard et al. (2009), and the remaining average spectral residuals of about -0.5 K are consistent with the results of Masiello et al. (2009). Other remaining residual features in these bands include a negative residual of -0.5 K in the 720 cm<sup>-1</sup> Q-branch of CO<sub>2</sub>, a positive offset of about 0.2 K between 755 and 770 cm<sup>-1</sup>, and a small residual feature near the  $\nu_3$  bandhead.

We have also examined the impact of the updated CO<sub>2</sub> spectroscopy on the temperature profiles retrieved using the  $\nu_2$  band of CO<sub>2</sub>. The retrieved temperatures with the

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updated spectroscopy are on average  $0.6 \pm 0.8$  K higher at 562 hPa,  $0.5 \pm 0.4$  K lower in the upper troposphere near 300 hPa,  $0.8 \pm 0.7$  K higher in the UTLS region between 100 and 200 hPa,  $0.4 \pm 0.4$  K lower between 40–100 hPa and  $0.3 \pm 0.5$  K lower between 10–30 hPa. These changes are consistent with the improvements in the spectroscopy between  $700\text{--}750\text{ cm}^{-1}$ , with most of the UTLS and stratospheric improvements likely due to spectroscopic improvements in the region near the  $720\text{ cm}^{-1}$  Q-branch.

The updated  $\text{H}_2\text{O}$  spectroscopy in LBLRTM v12.1 substantially reduces both the mean and the RMS of the residuals in the P-branch of the  $\text{H}_2\text{O } \nu_2$  band due to both the improved temperature retrieval and the improved  $\text{H}_2\text{O}$  spectroscopy in this region. The improvements in the R-branch are more modest and appear to be primarily due to the improved temperature retrievals in LBLRTM v12.1 rather than improvements in the  $\text{H}_2\text{O}$  spectroscopy, although the updated spectroscopy has reduced some systematic residual features in the R-branch that were present in the LBLRTM v9.4+ residuals. We find that the use of a fixed ratio of HDO to  $\text{H}_2\text{O}$  in LBLRTM may be responsible for a significant fraction of the remaining bias in the P-branch of the  $\text{H}_2\text{O } \nu_2$  band. We also find that including the calculated Coudert et al. (2008) line positions and intensities, as in LBLRTM v12.1, is necessary for proper modeling of the R-branch of the  $\text{H}_2\text{O } \nu_2$  band, as excluding the calculated values (as in HITRAN 2008) leads to discontinuities in the residuals within the R-branch.

The improved  $\text{H}_2\text{O}$  spectroscopy also has significant impacts on the retrieved  $\text{H}_2\text{O}$  profiles. The retrieved  $\text{H}_2\text{O}$  with the new spectroscopy is on average  $14\% \pm 8\%$  lower between 100 and 200 hPa,  $42\% \pm 58\%$  higher near 562 hPa, and  $31\% \pm 100\%$  higher near the surface. These changes are due to a combination of  $\text{H}_2\text{O}$  spectroscopy changes and the temperature profile changes (themselves caused by  $\text{CO}_2$  spectroscopy changes). The upper atmospheric changes are consistent with the results of Shephard et al. (2009), who showed that the addition of the Coudert et al. (2008) line positions and intensities to HITRAN 2004 resulted in a 10% reduction in the upper tropospheric  $\text{H}_2\text{O}$  mixing ratio. We also find that including the R-branch of the  $\text{H}_2\text{O } \nu_2$  band can provide more information on near-surface  $\text{H}_2\text{O}$  for moist cases, with the P-

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branch retrievals showing on average  $33\% \pm 51\%$  more  $\text{H}_2\text{O}$  than the P- and R-branch retrievals near the surface.

We have identified the remaining systematic residuals in the main absorption bands of  $\text{O}_3$ ,  $\text{CO}$ , and  $\text{CH}_4$  within the IASI spectral range. Neither the spectroscopy nor the retrievals of  $\text{O}_3$  have changed significantly between LBLRTM v9.4+ and v12.1, and thus significant average residuals of +0.3 K remain in the R-branch of the  $\nu_3$  band of  $\text{O}_3$ . While the updates to the  $\text{H}_2\text{O}$  self continuum in MT\_CKD v2.5 have clearly improved the match with observations near the  $\text{CO}_2 \nu_3$  bandhead, these updates have degraded the match with observations in the fundamental band of  $\text{CO}$  (between  $2060\text{--}2170\text{ cm}^{-1}$ ) and can create significant errors in the retrieved  $\text{CO}$  profile for profiles with high levels of water vapor. An improved version of the MT\_CKD self continuum is needed to preserve the improved  $\text{H}_2\text{O}$  self continuum values in the  $\text{CO}_2 \nu_3$  band while returning the  $\text{H}_2\text{O}$  self continuum in the  $2060\text{--}2170\text{ cm}^{-1}$  range to the values used in MT\_CKD 2.4. Finally, significant systematic residuals remain in the  $\nu_4$  band of  $\text{CH}_4$ , even with the updated spectroscopy, which leads to unphysically large retrieved mixing ratios of  $\text{CH}_4$  using our rather simple retrieval procedure. However, the average  $\text{CH}_4$  mixing ratios below 100 hPa are reduced by 44 ppbv when LBLRTM v12.1 is used instead of LBLRTM v9.4+, suggesting that this updated  $\text{CH}_4$  spectroscopy might improve more rigorous retrievals of  $\text{CH}_4$  from satellite observations, which currently can show positive biases relative to in situ observations. This suggests the need for further work on evaluating the updated  $\text{CH}_4$  spectroscopy for use in operational  $\text{CH}_4$  retrievals.

**Supplementary material related to this article is available online at:**  
[http://www.atmos-chem-phys-discuss.net/13/79/2013/  
acpd-13-79-2013-supplement.pdf](http://www.atmos-chem-phys-discuss.net/13/79/2013/acpd-13-79-2013-supplement.pdf).

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**Table 1.** Number of IASI spectra in each range of precipitable water vapor (PWV), and the associated surface temperature ( $T_{\text{srf}}$ ) and latitude range.

PWV (cm)	Number of Spectra	$T_{\text{srf}}$ Range (K)	Latitude Range
0–2	33	271–300	59° S to 21° S, 28° N to 62° N
2–4	57	293–303	35° S to 28° N
4–6	30	298–303	25° S to 21° N

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**Table 2.** List of spectral regions used in the retrievals.

Retrieved Parameter	Spectral regions used in retrievals [ $\text{cm}^{-1}$ ]
$T_{\text{srf}}$	817.0–823.0, 831.0–834.5, 843.0–848.0, 960.0–965.0, 1088.0–1090.0, 1144.0–1146.0, 1231.5–1232.50, 1330.25–1330.75, 2001.5–2005.0, 2011.5–2013.25, 2030.0–2032.0, 2499.0–2501.0, 2604.0–2606.0
$T_{\text{atm}}$ ( $\text{CO}_2 \nu_2$ band)	650.0–780.0
$T_{\text{atm}}$ ( $\text{CO}_2 \nu_3$ band)	2270.0–2499.0
$\text{H}_2\text{O}$	1164.5–1166.25, 1173.0–1175.5, 1186.0–1188.0, 1197.0–1199.0, 1210.75–1213.25, 1224.25–1226.25, 1242.5–1245.25, 1257.75–1261.75, 1375.0–1560.0, 1640.0–2020.0
$\text{O}_3$	990.0–1070.0
$\text{CO}$	2072.75–2074.0, 2094.0–2095.50, 2098.0–2099.75, 2102.25–2104.25, 2110.25–2112.50, 2118.75–2120.50, 2127.0–2135.0, 2149.5–2151.75, 2153.50–2155.50, 2157.25–2159.50, 2164.75–2177.0, 2179.0–2180.50, 2182.5–2184.0, 2186.0–2187.5, 2189.25–2190.75, 2192.75–2194.25
$\text{CH}_4$	1292.0–1305.0

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**Table 3.** Water vapor retrieval brightness temperature residuals (Mean and RMS in K, RMS in bold) for LBLRTM v12.1 when the H<sub>2</sub>O line parameters from HITRAN 2008 and those from AER v3.1 are used.

	0.3 cm PWV		1.5 cm PWV		5.4 cm PWV	
	AER v3.1	HITRAN 2008	AER v3.1	HITRAN 2008	AER v3.1	HITRAN 2008
P-branch (1375–1560 cm <sup>-1</sup> )	0.11 ( <b>0.35</b> )	0.13 ( <b>0.36</b> )	0.06 ( <b>0.39</b> )	0.07 ( <b>0.40</b> )	0.05 ( <b>0.36</b> )	0.05 ( <b>0.37</b> )
R-branch (1640–2020 cm <sup>-1</sup> )	0.09 ( <b>0.64</b> )	-0.02 ( <b>0.66</b> )	-0.06 ( <b>0.61</b> )	-0.18 ( <b>0.66</b> )	-0.05 ( <b>0.75</b> )	-0.19 ( <b>0.79</b> )
R-branch (1750–2020 cm <sup>-1</sup> )	0.12 ( <b>0.63</b> )	-0.03 ( <b>0.65</b> )	-0.12 ( <b>0.58</b> )	-0.28 ( <b>0.59</b> )	-0.13 ( <b>0.75</b> )	-0.31 ( <b>0.80</b> )

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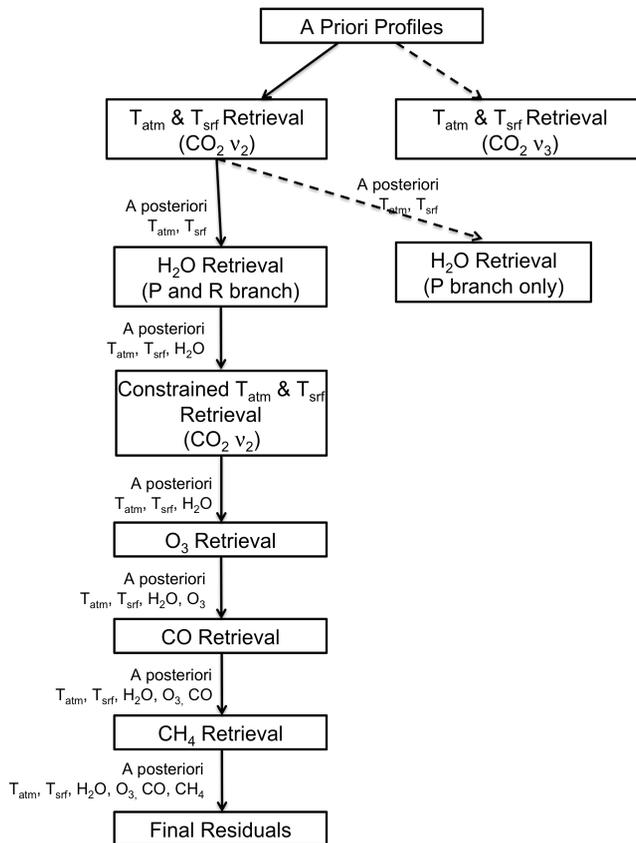
**Table 4.** Water vapor retrieval brightness temperature residuals (Mean and RMS in K, RMS in bold) for different combinations of retrieved temperature profiles and H<sub>2</sub>O spectroscopy.

	0.3 cm PWV		1.5 cm PWV		5.4 cm PWV	
	P-branch	R-branch	P-branch	R-branch	P-branch	R-branch
New T, New H <sub>2</sub> O	0.11 ( <b>0.35</b> )	0.09 ( <b>0.64</b> )	0.06 ( <b>0.39</b> )	−0.06 ( <b>0.61</b> )	0.05 ( <b>0.36</b> )	−0.05 ( <b>0.75</b> )
Old T, New H <sub>2</sub> O	0.16 ( <b>0.37</b> )	0.11 ( <b>0.64</b> )	0.07 ( <b>0.39</b> )	−0.06 ( <b>0.59</b> )	0.05 ( <b>0.37</b> )	−0.03 ( <b>0.76</b> )
New T, Old H <sub>2</sub> O	0.09 ( <b>0.37</b> )	0.04 ( <b>0.64</b> )	0.06 ( <b>0.42</b> )	−0.08 ( <b>0.58</b> )	0.04 ( <b>0.38</b> )	−0.05 ( <b>0.72</b> )
Old T, Old H <sub>2</sub> O	0.13 ( <b>0.39</b> )	0.05 ( <b>0.64</b> )	0.07 ( <b>0.41</b> )	−0.08 ( <b>0.57</b> )	0.04 ( <b>0.38</b> )	−0.05 ( <b>0.73</b> )

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**Fig. 1.** Schematic of the retrieval procedure used in this study. Solid arrows show the main line from the a priori profiles described in Sect. 4 to the final residuals for the two models (LBLRTM v12.1 and LBLRTM v9.4+). The dashed lines show additional retrievals done to assess the consistency of spectroscopy between the  $\nu_2$  and  $\nu_3$  bands of  $\text{CO}_2$  and the P- and R-branches of  $\text{H}_2\text{O}$ .

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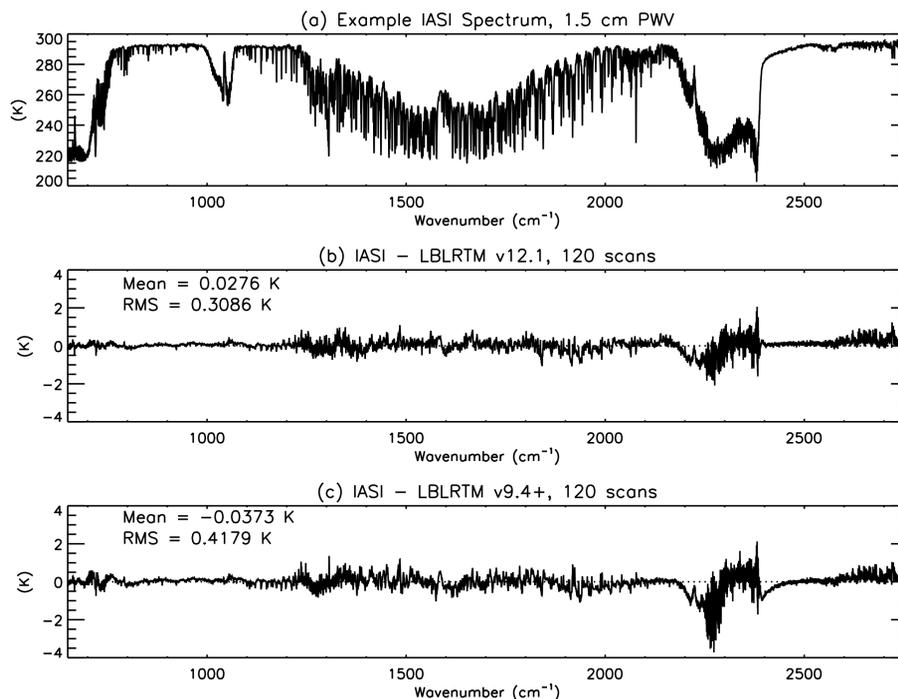
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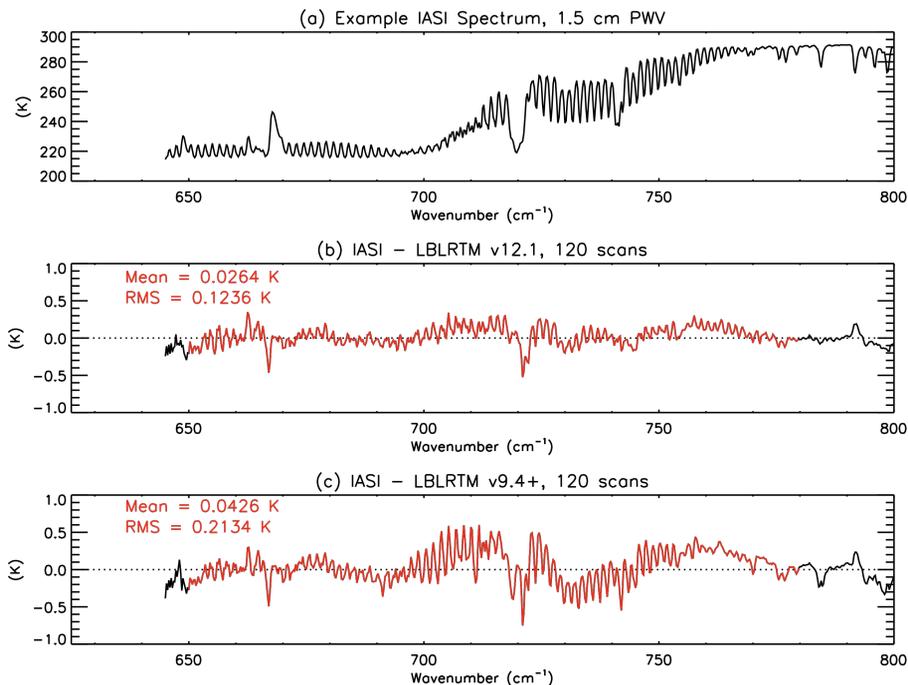


**Fig. 2.** (a) IASI observed brightness temperature spectrum for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+.

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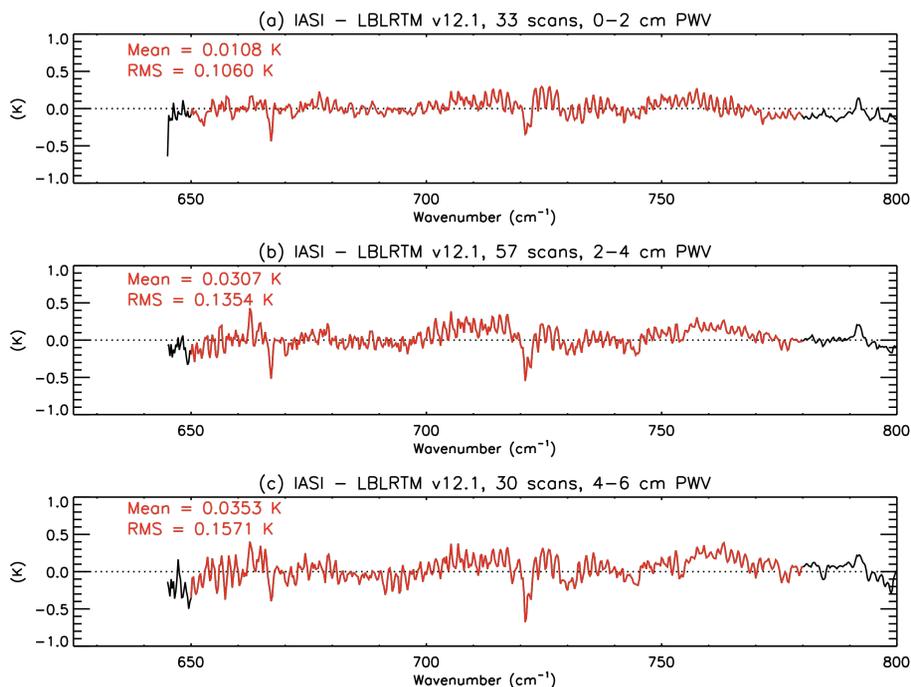


**Fig. 3.** (a) IASI observed brightness temperature spectrum in the  $\nu_2$  band of  $\text{CO}_2$  for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The  $\nu_2$  retrieval microwindow (which was used to retrieve  $T_{\text{atm}}$  and  $T_{\text{srf}}$ ) is highlighted in red.

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**Fig. 4.** Mean brightness temperature residuals in the  $\text{CO}_2$   $\nu_2$  band using LBLRTM v12.1 as in Fig. 3b, but only for profiles with (a) 0 to 2 cm, (b) 2 to 4 cm, and (c) 4 to 6 cm of precipitable water vapor (PWV).

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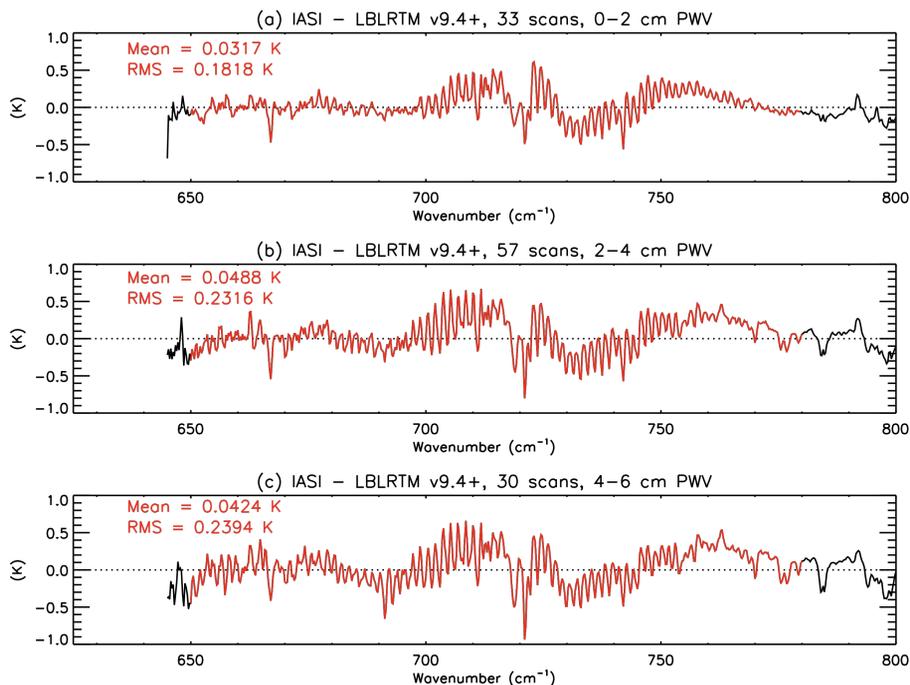
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**Fig. 5.** Mean brightness temperature residuals in the  $\text{CO}_2$   $\nu_2$  band using LBLRTM v9.4+ as in Fig. 3c, but only for profiles with **(a)** 0 to 2 cm, **(b)** 2 to 4 cm, and **(c)** 4 to 6 cm of precipitable water vapor (PWV).

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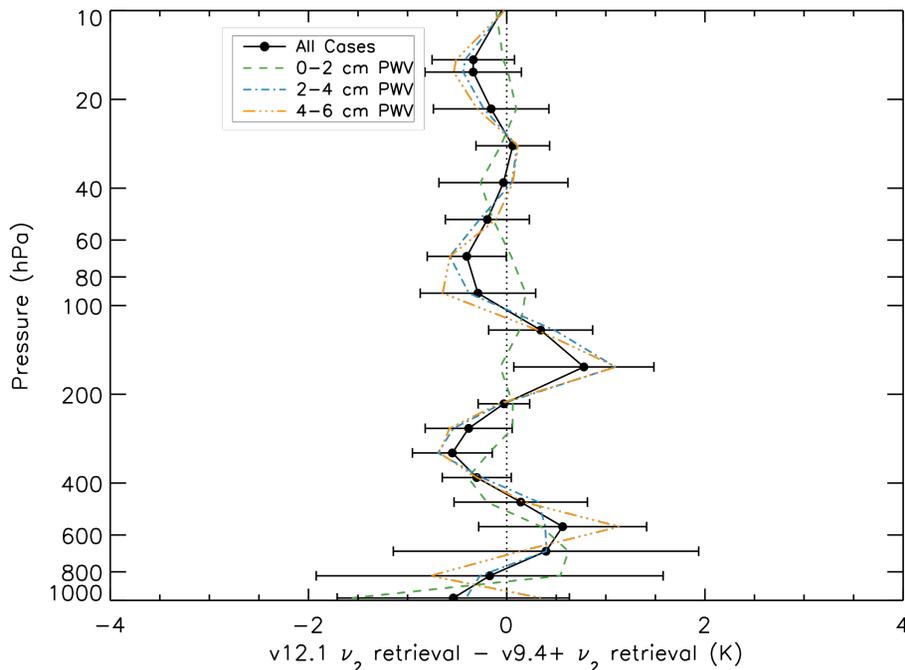
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**Fig. 6.** Mean (solid line) and standard deviation (error bars) of the differences between the 120 temperature profiles retrieved with LBLRTM v12.1 and v9.4+. Dashed lines show the mean differences for the spectra with 0–2 cm (green), 2–4 cm (blue), and 4–6 cm (orange) PWV. Note that these temperature retrievals used the  $\text{CO}_2 \nu_2$  band but not the  $\text{CO}_2 \nu_3$  band.

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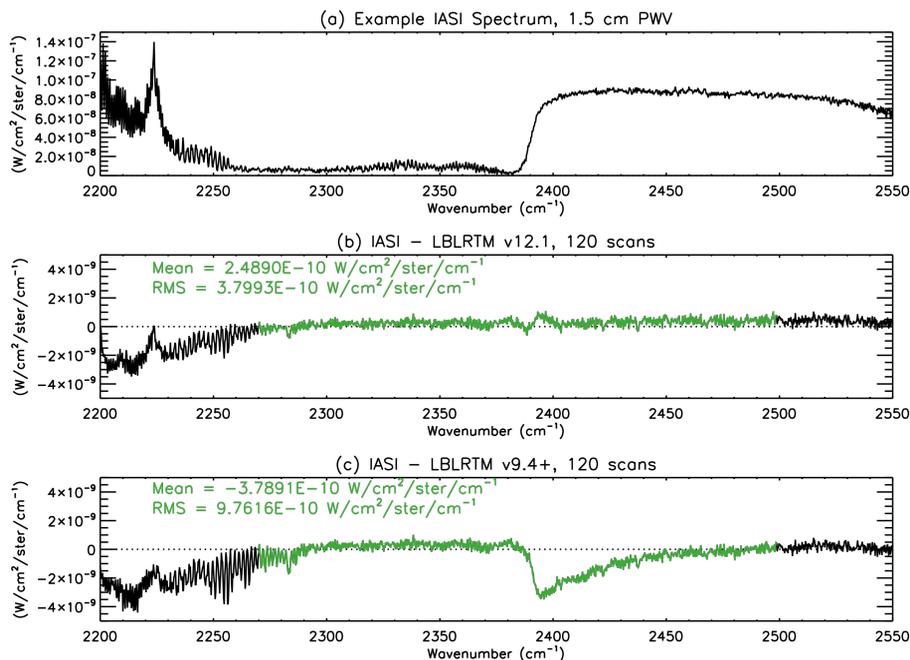
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**Fig. 7.** (a) IASI observed brightness temperature spectrum in the  $\nu_3$  band of  $\text{CO}_2$  for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The  $\text{CO}_2$   $\nu_3$  band, which was not used in the temperature retrievals, is highlighted in green.

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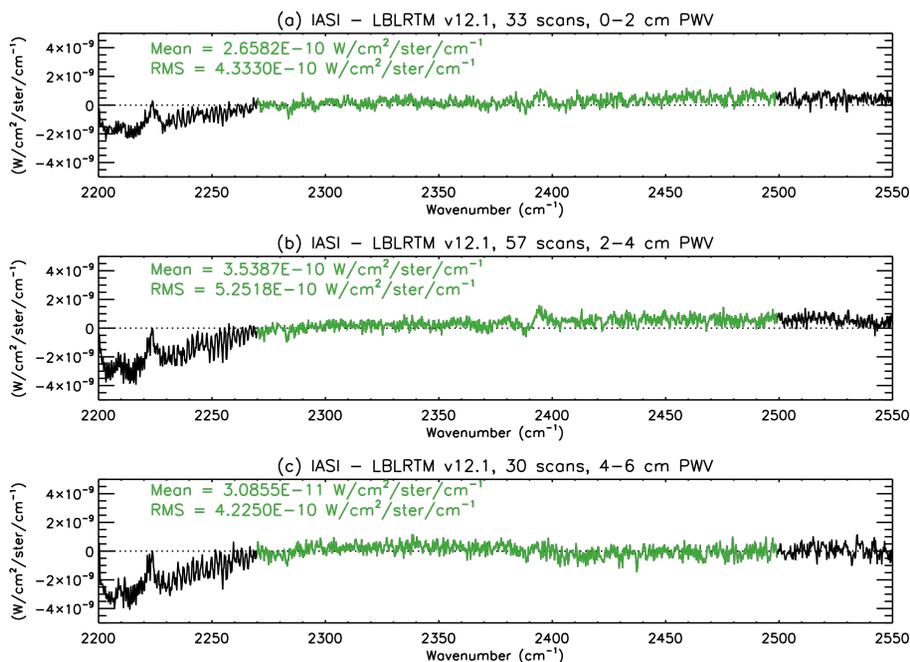
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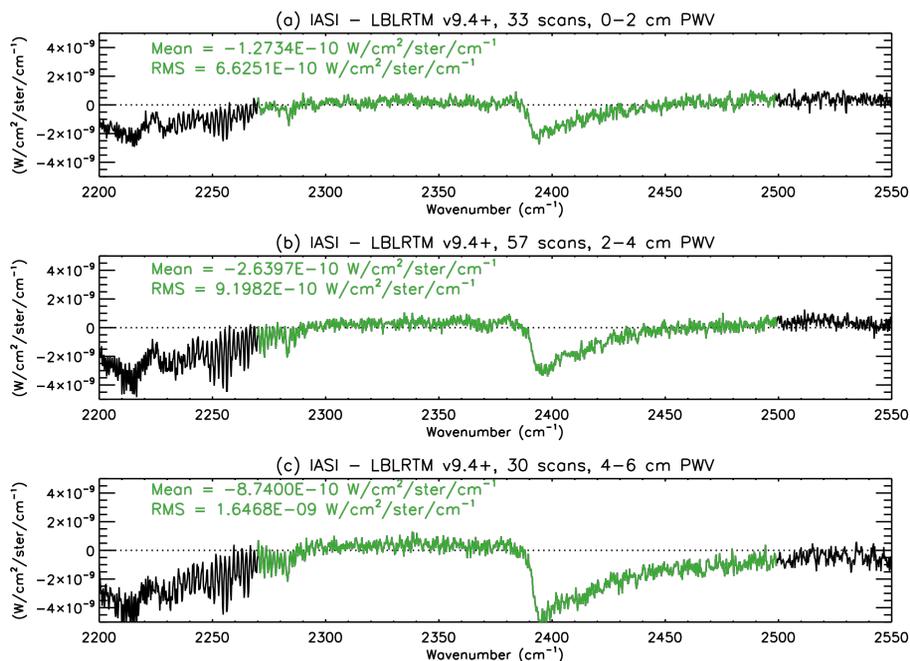


**Fig. 8.** Mean brightness temperature residuals in the CO<sub>2</sub> v<sub>3</sub> band using LBLRTM v12.1 as in Fig. 7b, but only for profiles with (a) 0 to 2 cm, (b) 2 to 4 cm, and (c) 4 to 6 cm of precipitable water vapor.

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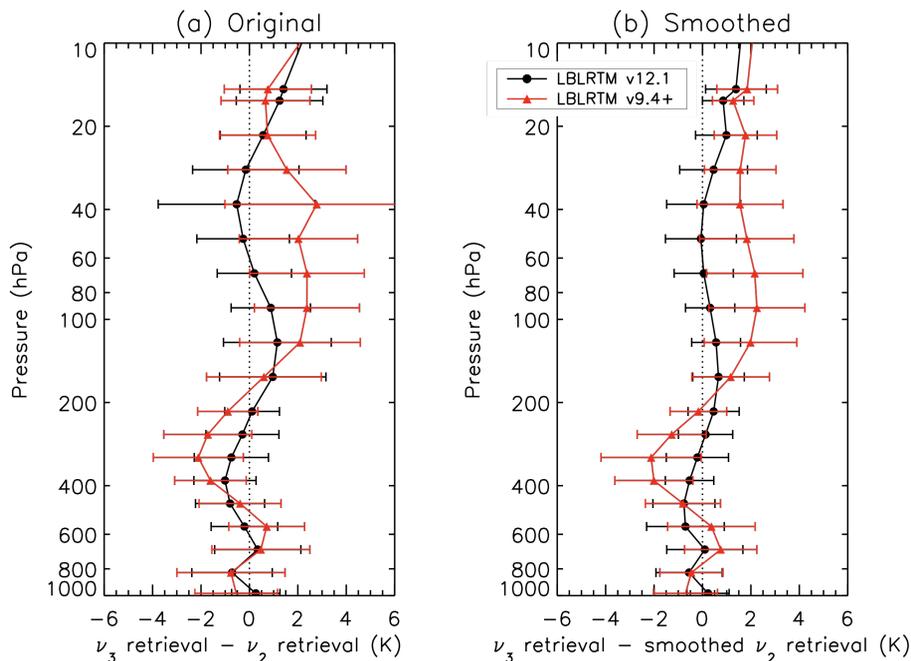


**Fig. 9.** Mean brightness temperature residuals in the CO<sub>2</sub> v<sub>3</sub> band using LBLRTM v9.4+ as in Fig. 7c, but only for profiles with (a) 0 to 2 cm, (b) 2 to 4 cm, and (c) 4 to 6 cm of precipitable water vapor.

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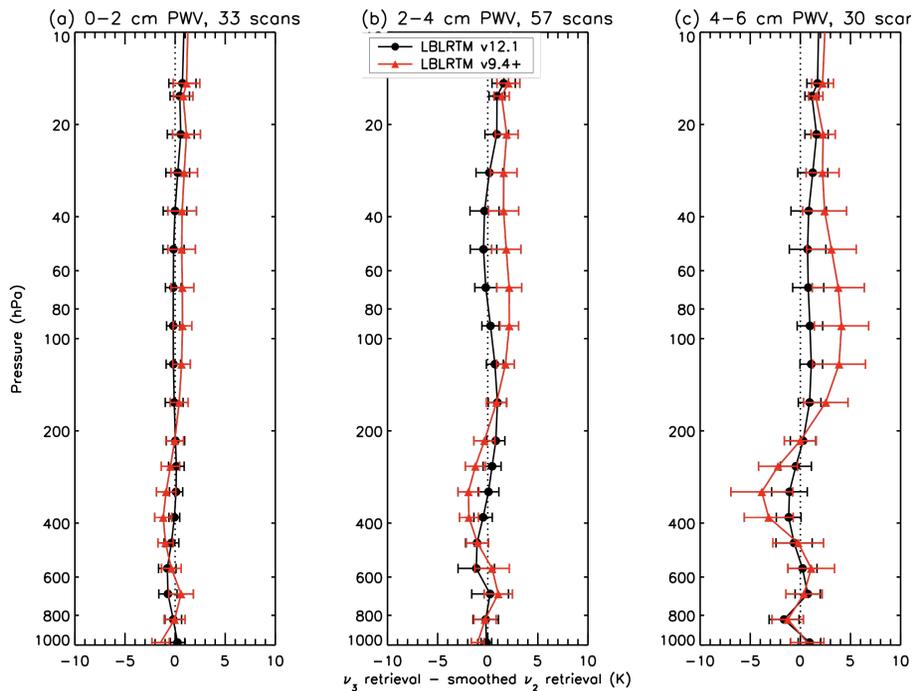


**Fig. 10.** Mean (solid line) and standard deviation (error bars) of the differences between the 120 temperature profiles retrieved with the  $\nu_3$  band of  $\text{CO}_2$  and the profile retrieved with the  $\nu_2$  retrieval. **(a)** The mean differences between the original profiles. **(b)** The mean differences between the profiles after the  $\nu_2$  retrieval was smoothed using the averaging kernel of the  $\nu_3$  retrieval, as described in the text.

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**Fig. 11.** As in Fig. 10b, but only for profiles with **(a)** 0 to 2 cm, **(b)** 2 to 4 cm, and **(c)** 4 to 6 cm of PWV.

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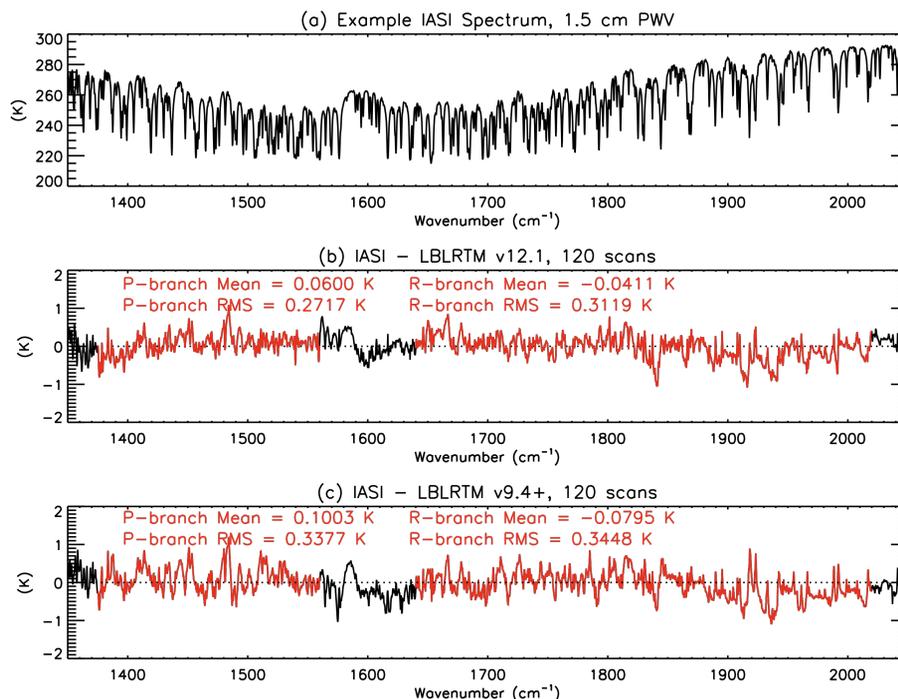
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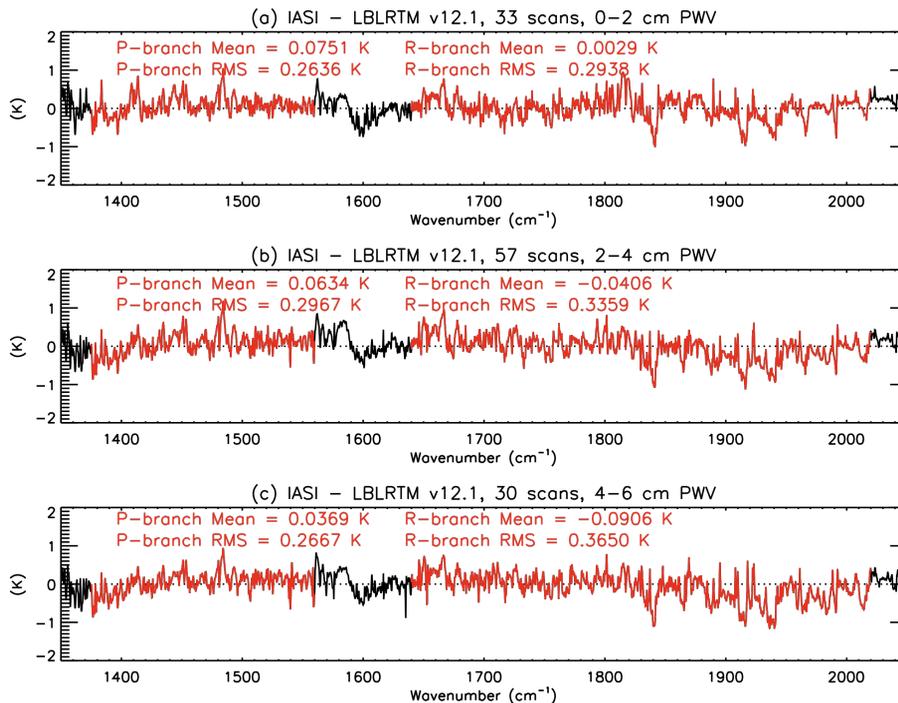


**Fig. 12.** (a) IASI observed brightness temperature spectrum in the  $\nu_2$  band of  $\text{H}_2\text{O}$  for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The P- and R-branch  $\text{H}_2\text{O}$  retrieval microwindows are highlighted in red.

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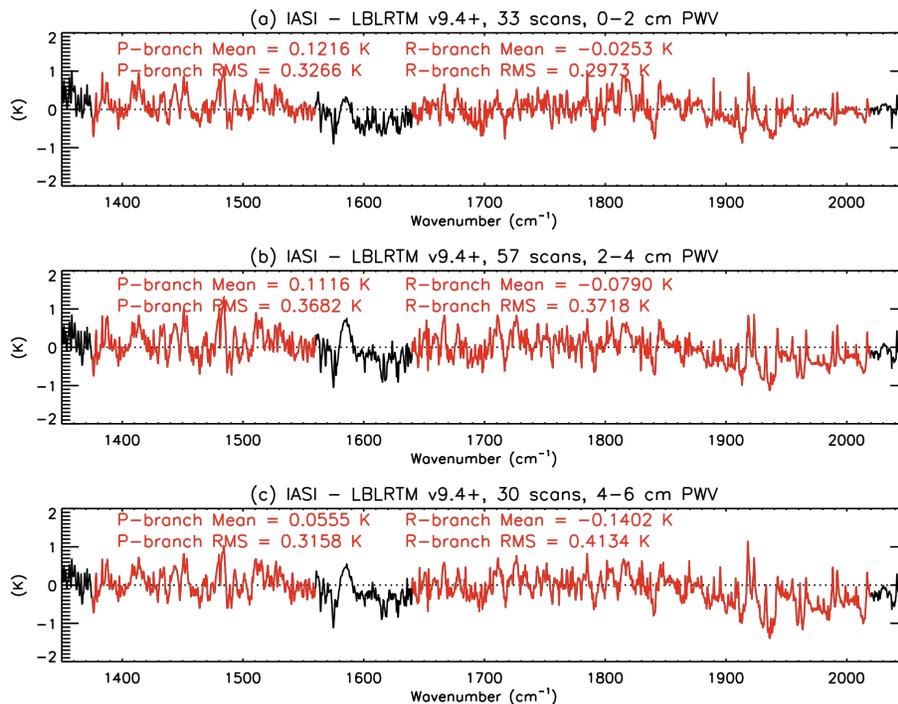


**Fig. 13.** Mean brightness temperature residuals in the  $\nu_2$  band of  $\text{H}_2\text{O}$  using LBLRTM v12.1 as in Fig. 12b, but only for profiles with (a) 0 to 2 cm, (b) 2 to 4 cm, and (c) 4 to 6 cm of precipitable water vapor.

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**Fig. 14.** Mean brightness temperature residuals in the  $\nu_2$  band of  $\text{H}_2\text{O}$  using LBLRTM v9.4+ as in Fig. 12c, but only for profiles with (a) 0 to 2 cm, (b) 2 to 4 cm, and (c) 4 to 6 cm of precipitable water vapor.

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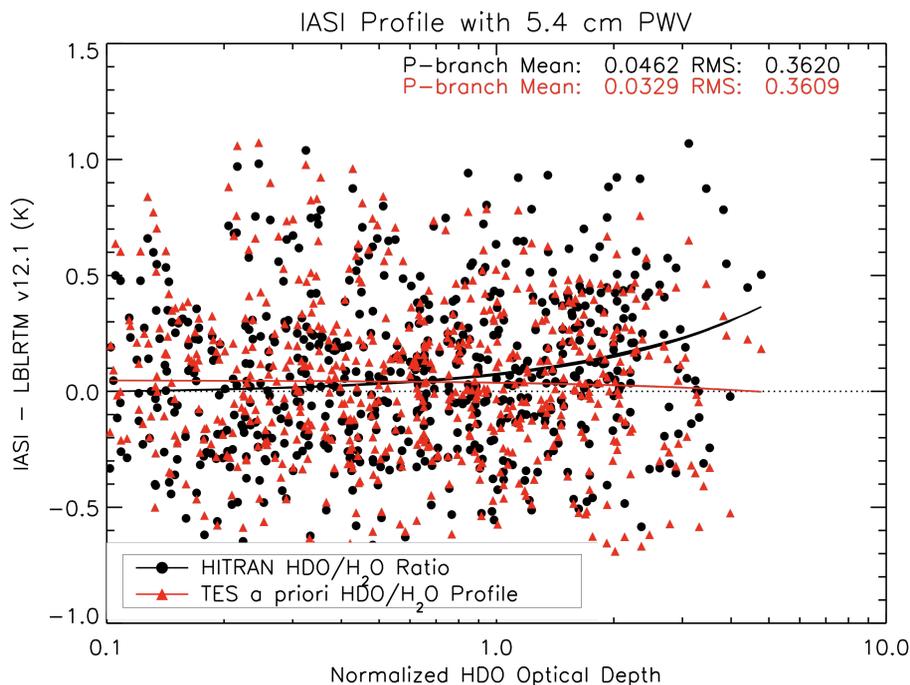
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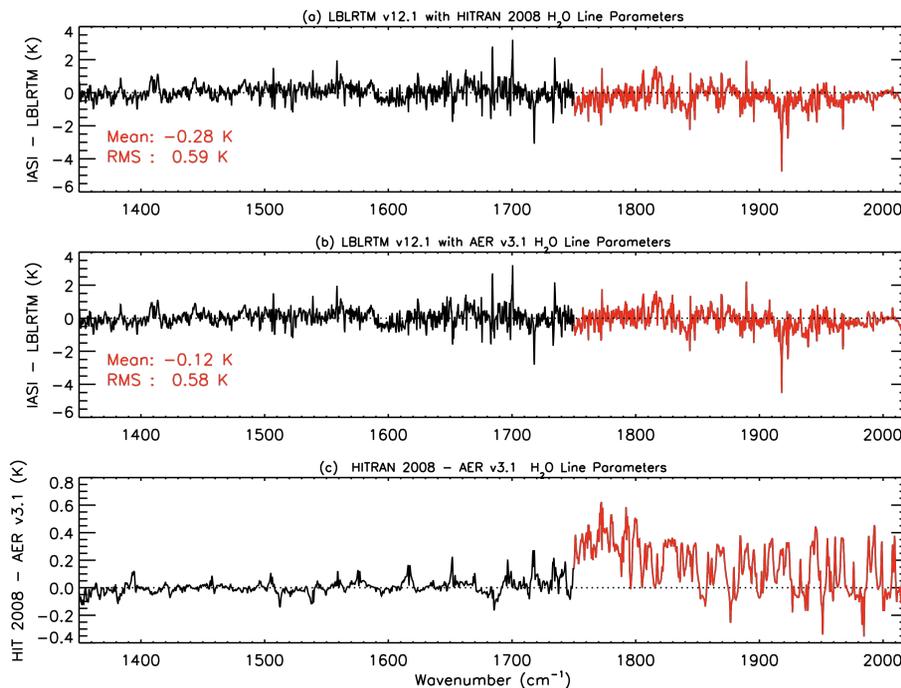


**Fig. 15.** LBLRTM v12.1 brightness temperature residuals versus HDO optical depth for a profile with 5.4 cm PWV. The black circles show the results when the constant HITRAN HDO/H<sub>2</sub>O ratio is used, while the red triangles show the results when the TES a priori profile of the HDO/H<sub>2</sub>O ratio is used. The black and red lines are the corresponding least-squares linear fits.

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**Fig. 16.** (a) Brightness temperature residuals for an example profile with 1.5 cm PWV after the  $\text{H}_2\text{O}$  P- and R-branch retrieval when the HITRAN 2008  $\text{H}_2\text{O}$  line parameters are used in LBLRTM v12.1. (b) Same as (a) but for the AER v3.1 line parameters used in LBLRTM v12.1 in the rest of this study. (c) Difference between (b) and (a), showing the discontinuity at wavenumbers greater than  $\sim 1750$   $\text{cm}^{-1}$ , where HITRAN 2008 does not use the Coudert et al. (2008) line positions and intensities included in AER v3.1.

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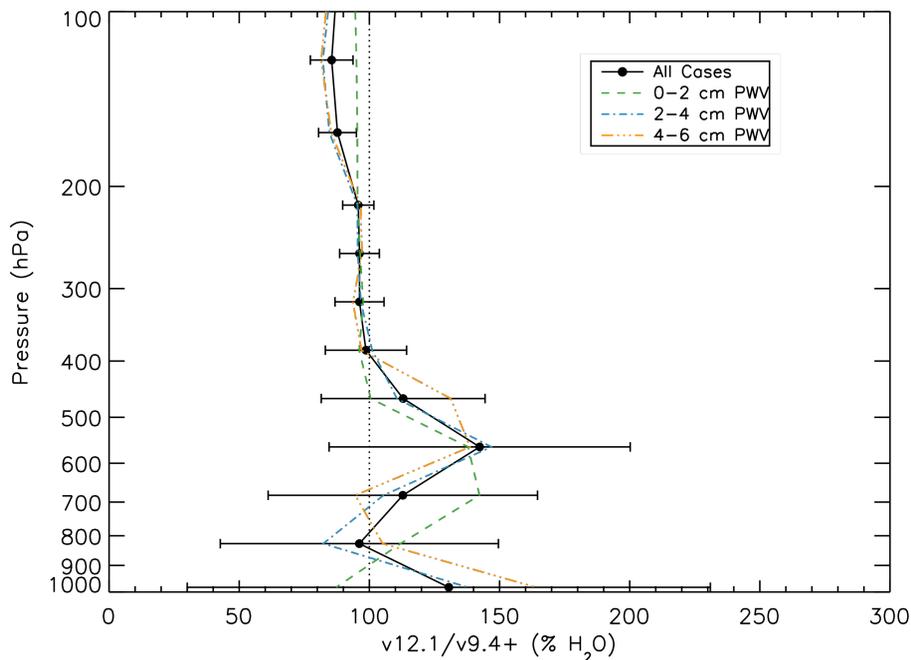
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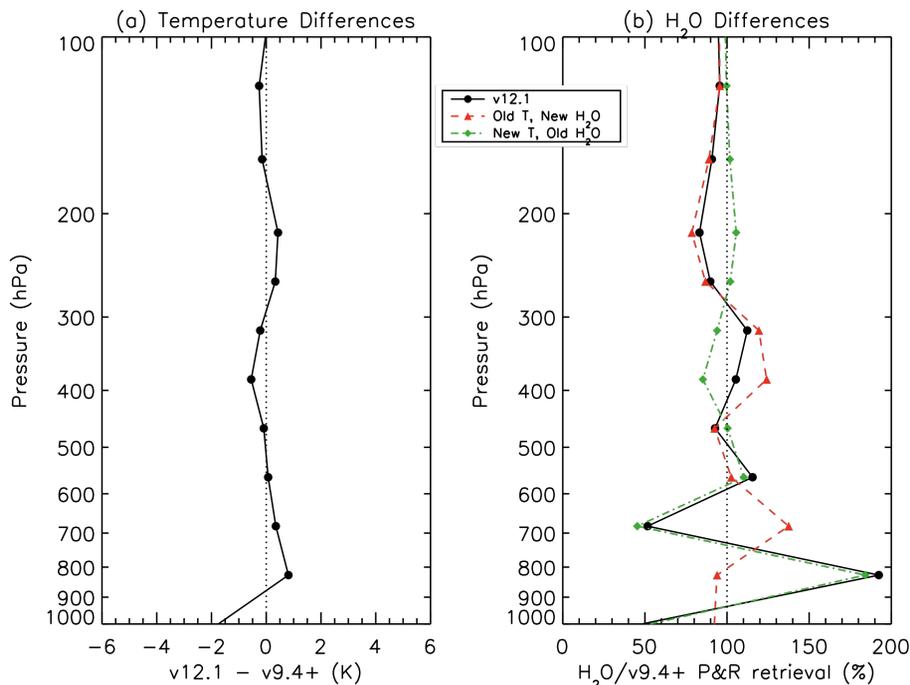


**Fig. 17.** Mean (black circles) and standard deviation (error bars) of the ratio of the 120 LBLRTM v12.1 retrieved  $\text{H}_2\text{O}$  profiles to the v9.4+ retrieved profiles, expressed as a percent. Dashed lines show the mean differences for the spectra with 0–2 cm (green), 2–4 cm (blue), and 4–6 cm (orange) PWV. Both models used the P- and R-branches in the retrievals.

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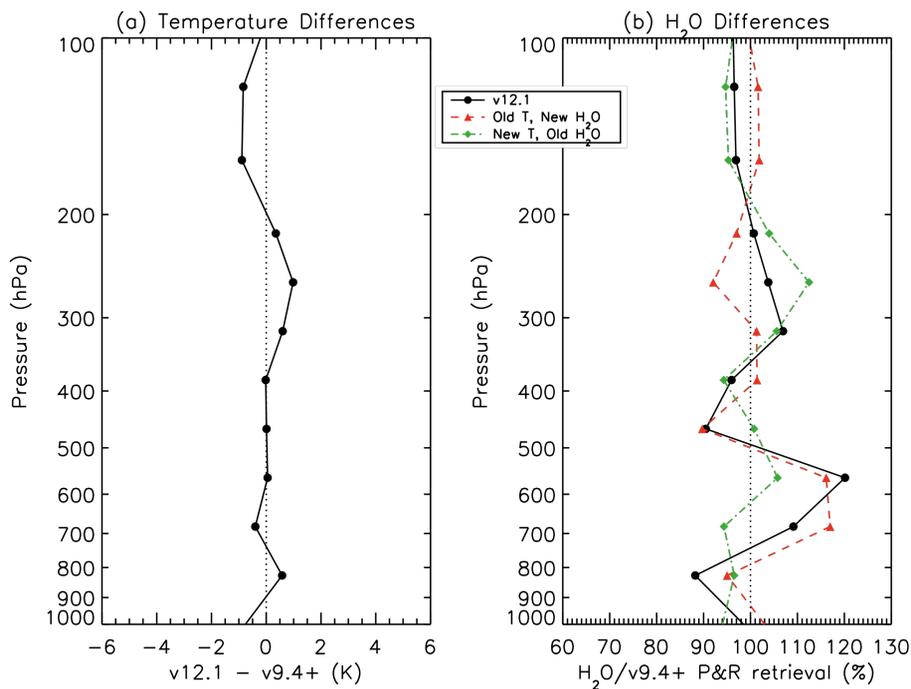


**Fig. 18.** (a) Difference of the retrieved temperature profiles between LBLRTM v12.1 and v9.4+ for an example spectrum with 1.5 cm PWV. The profiles were retrieved using the CO<sub>2</sub>  $\nu_2$  band. (b) The H<sub>2</sub>O profile retrieved using (i) LBLRTM v12.1 to retrieve both temperature and H<sub>2</sub>O (black circles), (ii) LBLRTM v9.4+ to retrieve temperature and LBLRTM v12.1 to retrieve H<sub>2</sub>O (red triangles), and (iii) LBLRTM v12.1 to retrieve temperature and LBLRTM v9.4+ to retrieve H<sub>2</sub>O (green diamonds), all normalized by the profile retrieved using LBLRTM v9.4+ to retrieve both temperature and H<sub>2</sub>O.

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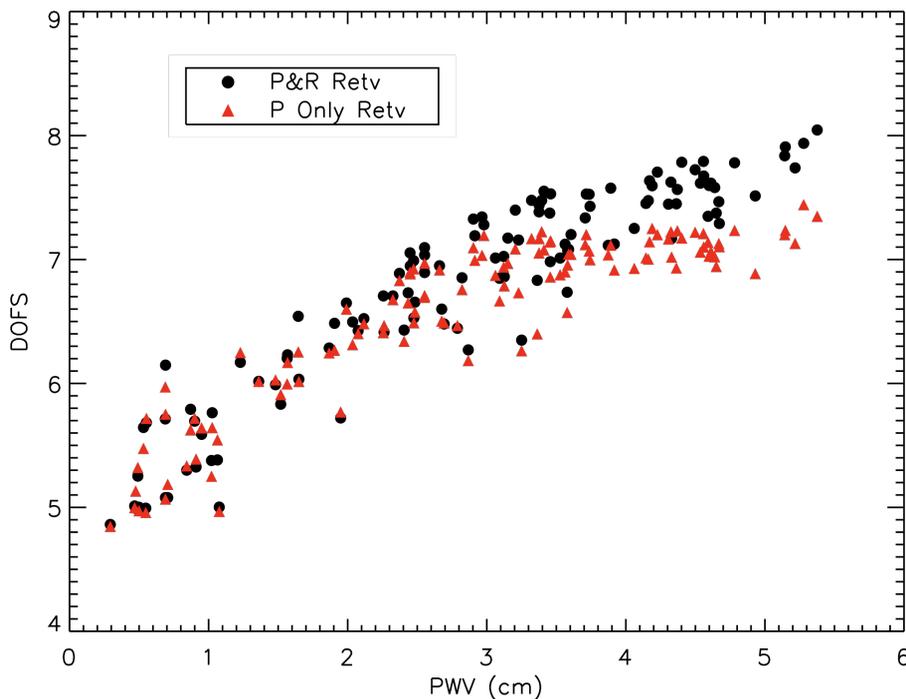


**Fig. 19.** Same as Fig. 18 but for an example spectrum with 0.3 cm PWV.

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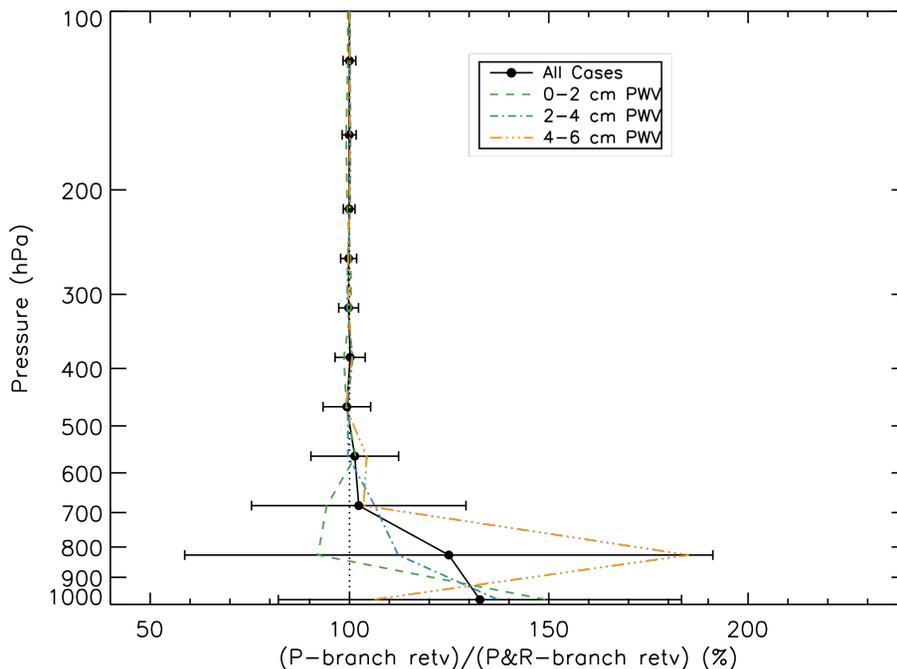


**Fig. 20.** Degrees of freedom for signal (DOFS) versus PWV for the P- and R-branch retrieval (black circles) and the P-branch only retrieval (red triangles) of  $\text{H}_2\text{O}$  using LBLRTM v12.1.

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**Fig. 21.** Mean (black circles) and standard deviation (error bars) of the ratio of the 120  $\text{H}_2\text{O}$  profiles retrieved with LBLRTM v12.1 using the P branch of the  $\nu_2$  band of  $\text{H}_2\text{O}$  to the profiles retrieved using both the P and R branches. Dashed lines show the mean differences for the spectra with 0–2 cm (green), 2–4 cm (blue), and 4–6 cm (orange) PWV.

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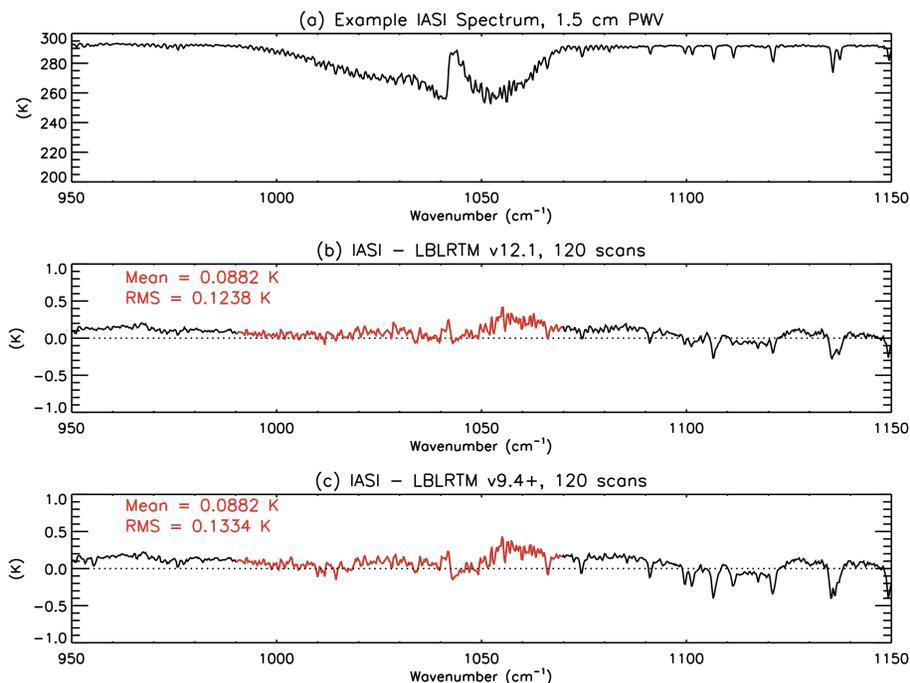
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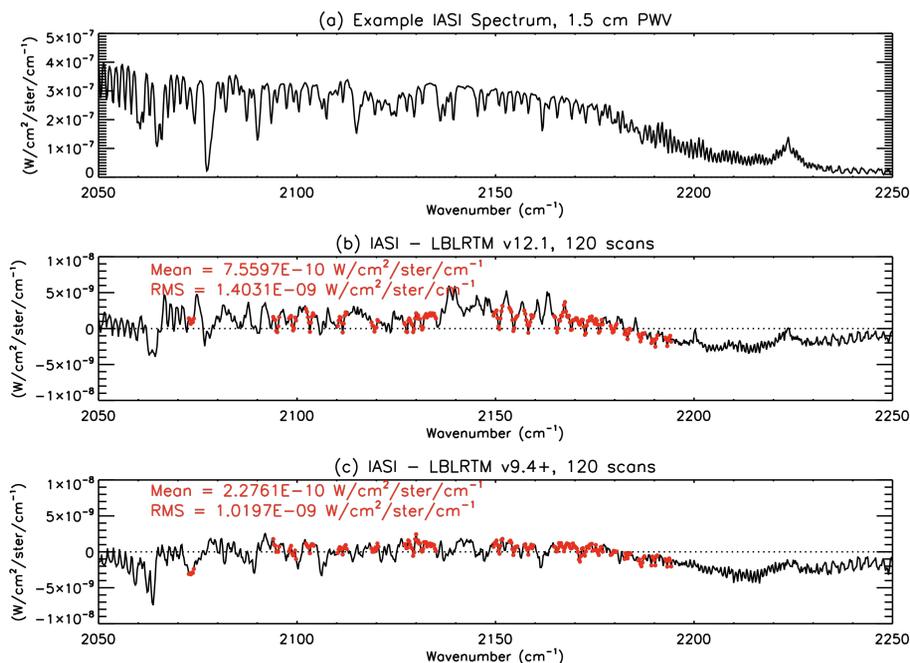


**Fig. 22.** (a) IASI observed brightness temperature spectrum in the  $\nu_3$  band of O<sub>3</sub> for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The O<sub>3</sub> retrieval microwindow is highlighted in red.

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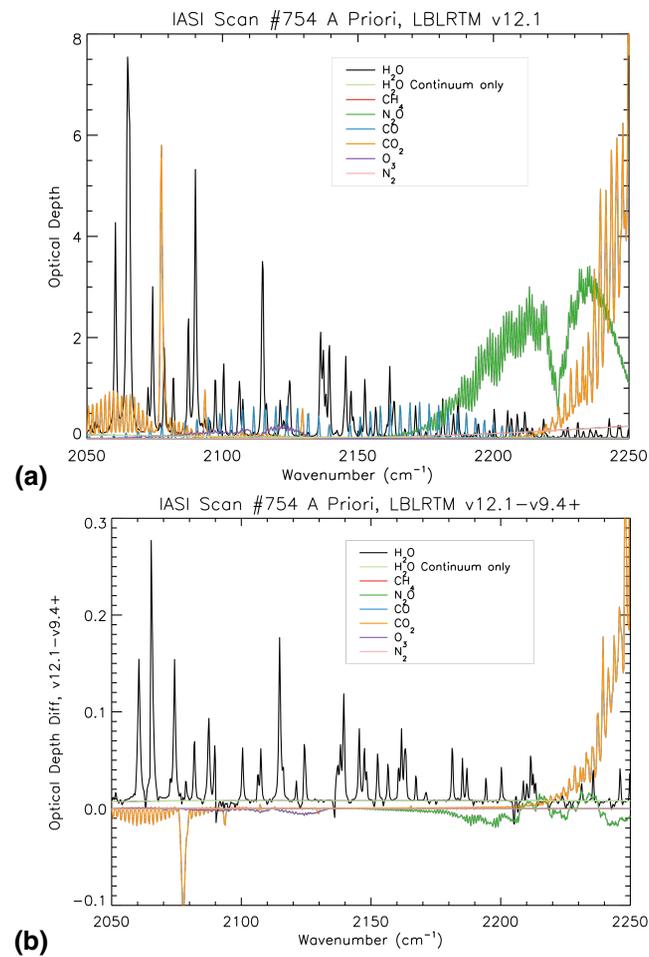


**Fig. 23.** (a) IASI observed brightness temperature spectrum in the fundamental vibrational band of CO for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The CO retrieval microwindows are highlighted in red.

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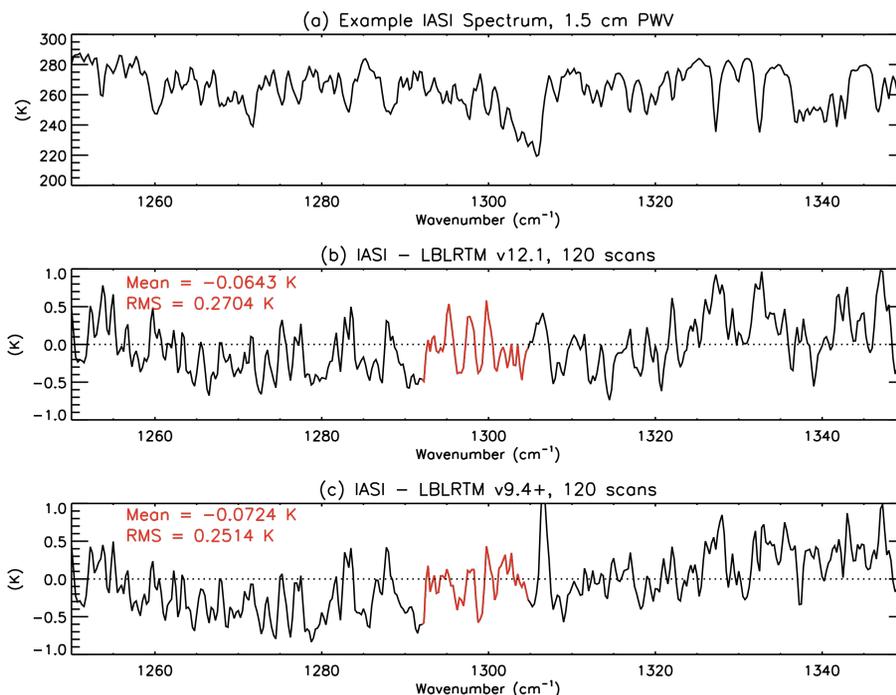


**Fig. 24. (a)** Optical depth for each species near the fundamental vibrational band of CO. **(b)** Change in optical depth between LBLRTM v12.1 and LBLRTM v9.4+

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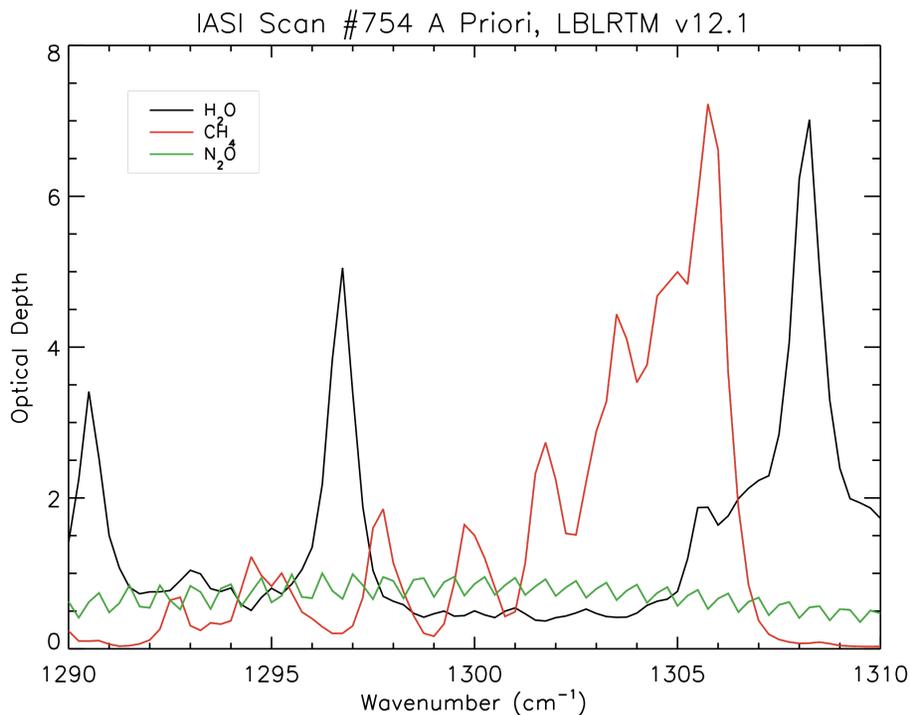


**Fig. 25.** (a) IASI observed brightness temperature spectrum in the  $\nu_4$  band of  $\text{CH}_4$  for an example profile with 1.5 cm PWV. (b) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v12.1. (c) Mean of the final brightness temperature residuals for 120 spectra using LBLRTM v9.4+. The  $\text{CH}_4$  retrieval microwindow is highlighted in red.

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**Fig. 26.** Optical depth of  $\text{H}_2\text{O}$ , (which includes HDO),  $\text{N}_2\text{O}$ , and  $\text{CH}_4$  near the  $\nu_4$  band of  $\text{CH}_4$ .

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